

Methodology

Conditions for the formation of massive seed black holes

1. Major merger (1:3) of gas-rich late-type galaxies ($B/T < 0.2$)
2. Host halo $M_h > 10^{11} M_{\text{Sun}}$
3. No a pre-existing black hole of $M_{\text{BH}} > 10^6 M_{\text{Sun}}$

Evolution of the gas component in major merger of disk



S. Kazantzidis (2004)

GAS DENSITY

Coplanar Equal-Mass Cooling+SF

10% Gas Fraction

time = 0.00 Gyr



Multi-scale galaxy merger simulations

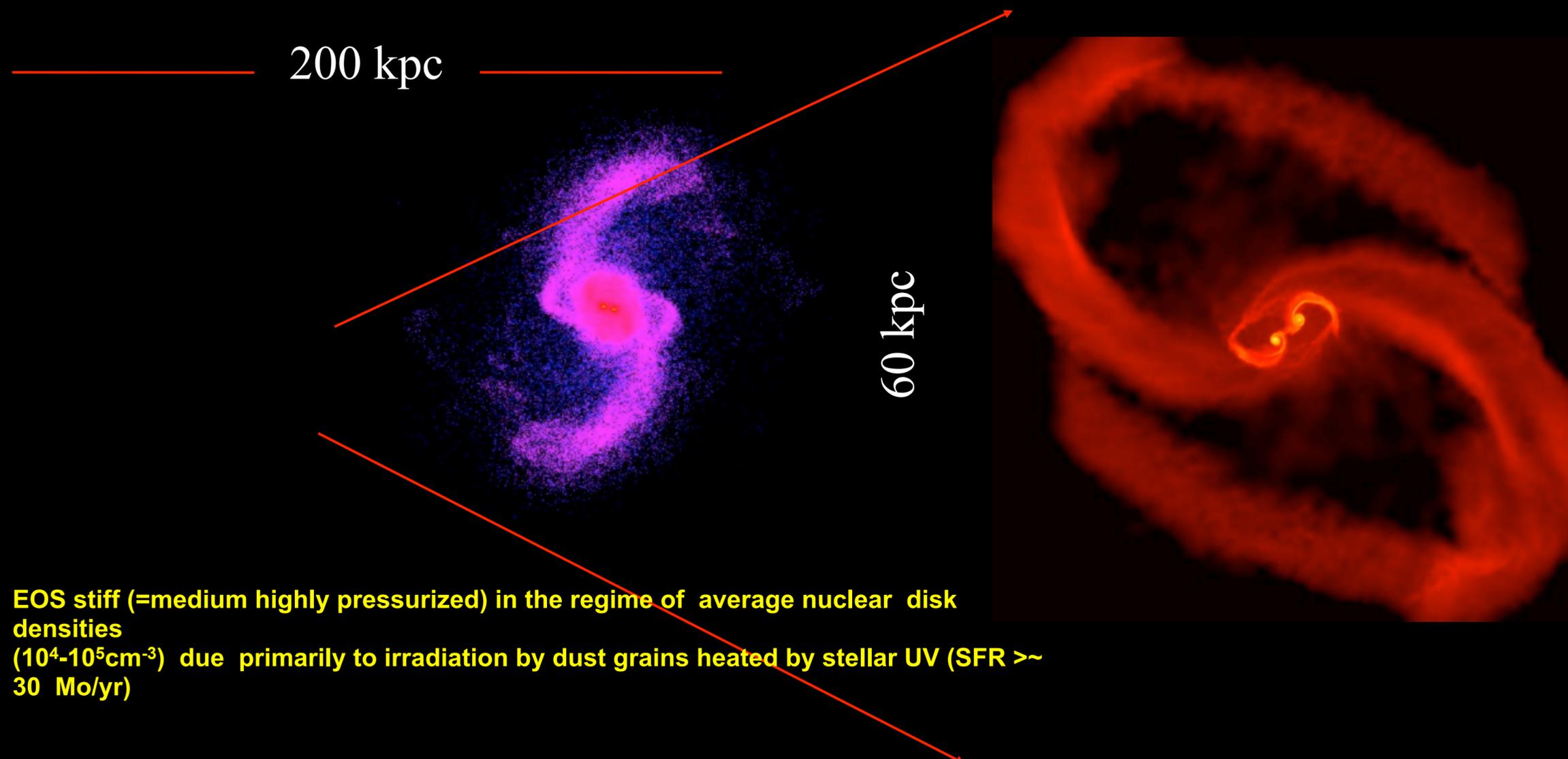
from ~ 100 kpc to 0.1 pc

Mayer et al. 2007, 2008, 2010

Using Smoothed Particle Hydrodynamics (SPH code GASOLINE) + splitting of gas particles (Kitsionas & Withworth 2002, Bromm 2004) to increase mass and spatial resolution as galaxy merger proceeds

Max. Resolution 3000 solar masses and 0.1 pc

Effective equation of state (EOS) - ideal gas, $P = (\gamma - 1) \rho \epsilon$, varying effective " γ " - to model local balance between heating and cooling in nuclear region (based on Spaans & Silk 2000; 2005 - steady-state interstellar gas model heated by starburst w/ radiative transfer)



Multi-scale galaxy merger simulations from ~100 kpc to 0.1 pc

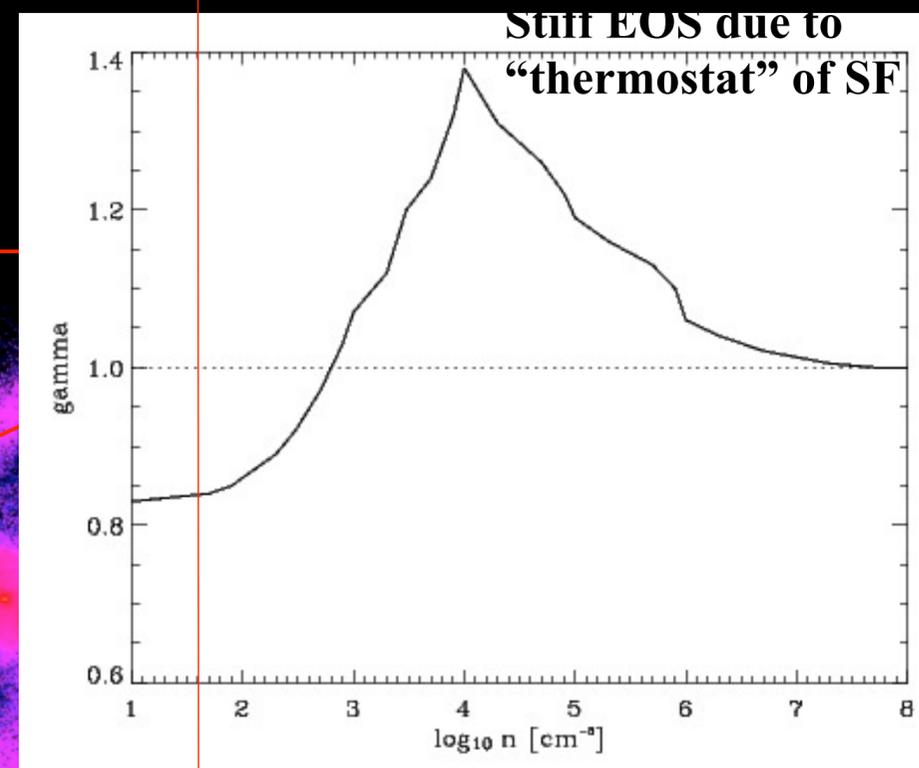
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200 kpc



EOS stiff (=medium highly pressurized) in the regime of average nuclear disk densities (10^4 - 10^5cm^{-3}) due primarily to irradiation by dust grains heated by stellar UV (SFR $> \sim 30 \text{ Mo/yr}$)

SELF-GRAVITATING GAS DISKS: STABILITY and INFLOWS

THREE REGIMES:

Toomre parameter $Q = \kappa v_s / \pi G \Sigma$

(from linear local
perturbative analysis
of self-gravitating
rotating fluid in
infinitesimally thin disk)

$Q < 1$ locally unstable to collapse \rightarrow fragmentation on
dynamical timescale (t_{dyn}) \rightarrow gas clumps make stars

$1 < Q < 2$ locally stable, globally unstable
to non-axisymmetric modes
(spiral modes, bar modes)

\rightarrow angular momentum transport (on a few t_{dyn})
via spiral density waves (Lynden Bell & Pringle 1979;
Lin & Pringle 1987; Laughlin & Adams 2000)
 \rightarrow gas inflow towards state of minimum energy

$Q > 2$ locally and globally stable \rightarrow dynamically
uninteresting

-Sweet spot ($1 < Q < 2$): a non-fragmenting globally unstable disk to sustain central
gas inflow

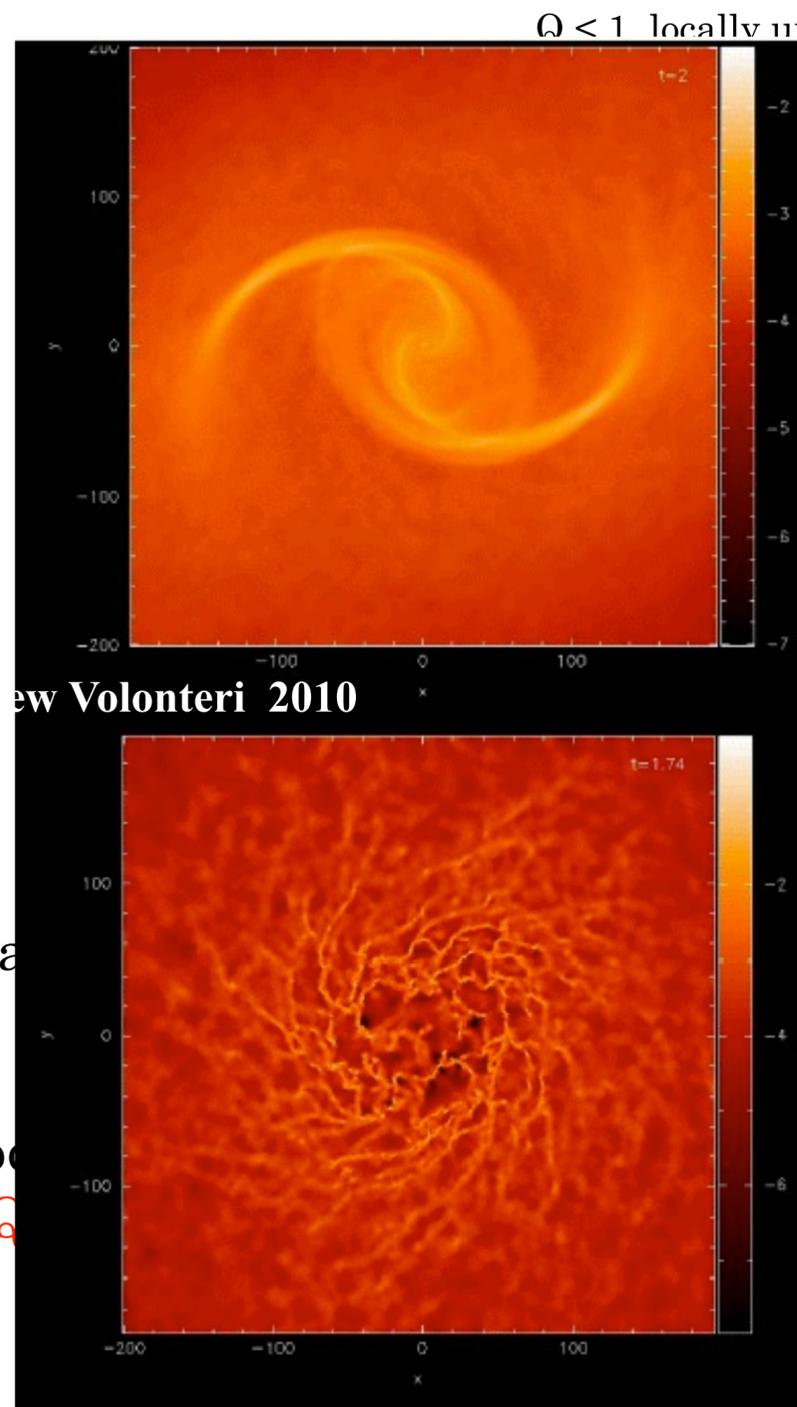
- The dissipation rate in the system is crucial – if cooling efficient amplitude of
non-axisymmetric modes increases \rightarrow inflow increases but $Q \lesssim 1$ approached
($T_{\text{cool}} < T_{\text{dyn}}$ drives Q below 1, while with $T_{\text{cool}} > T_{\text{dyn}}$ self-regulation to $Q \gtrsim 1$)

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new Volonteri 2010

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globally unstable
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accretion. The maximum amplitude of
perturbations is limited by the
fact that as $Q \lesssim 1$ is approached
(Lynden Bell self-regulation to $Q \gtrsim 1$)

- Sweet spot ($1 < Q < 2$): a
gas inflow
- The dissipation rate
non-axisymmetric modes
($T_{cool} < T_{dyn}$ drives Q)

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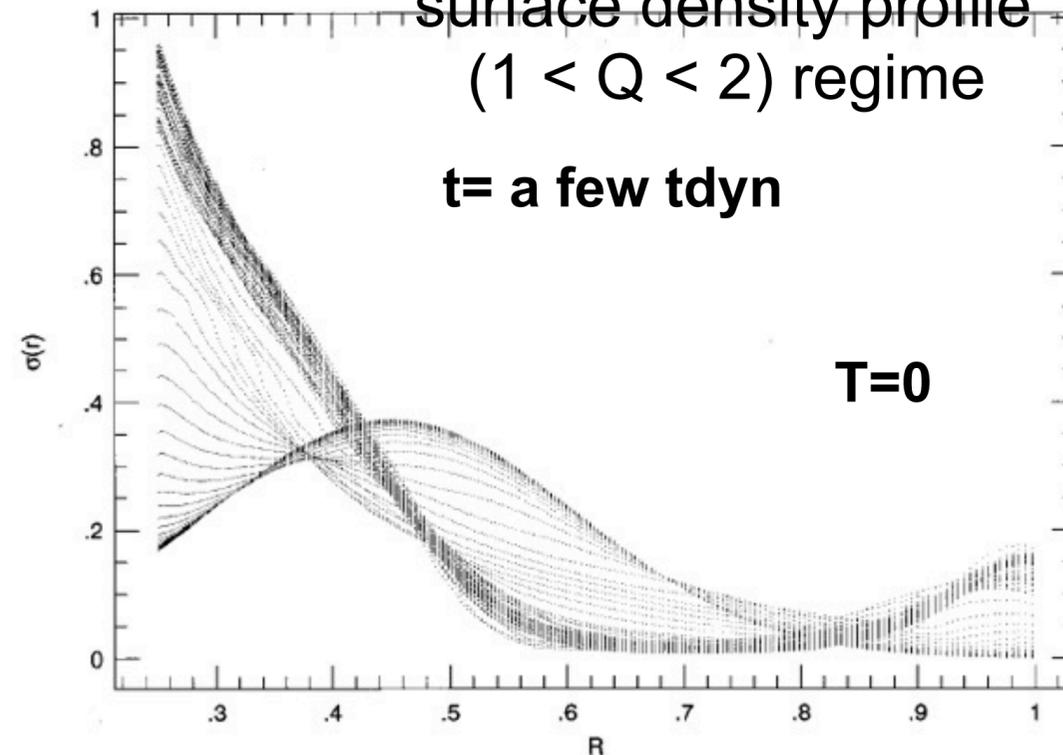
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Evolution of disk gas
surface density profile
($1 < Q < 2$) regime

$t = \text{a few } t_{\text{dyn}}$



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INFLOW BOTTLENECK: COOLING AND FRAGMENTATION

In system that cools rapidly ($t_{\text{cool}} < t_{\text{dyn}}$) and accumulates gas via inflow eventually Q drops to < 1 and fragmentation/star formation takes over

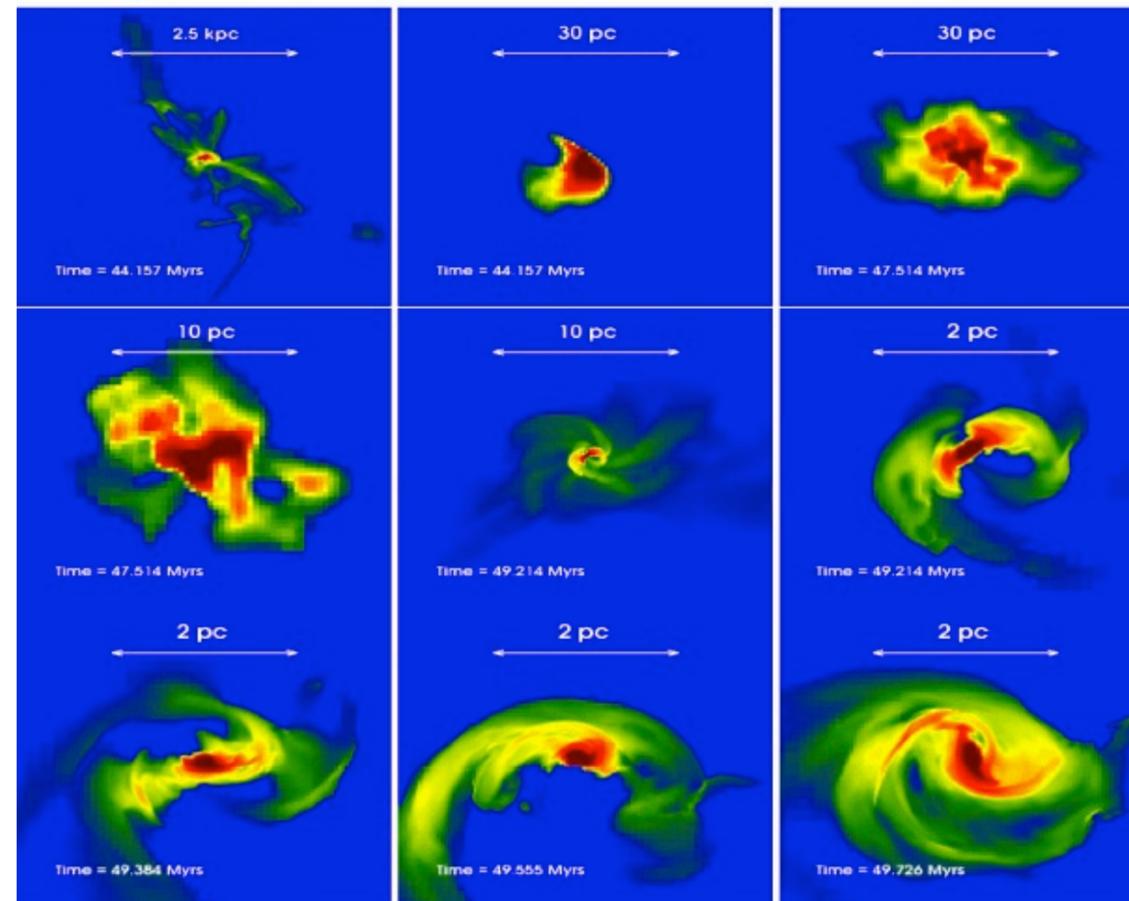
CONVENTIONAL WAY-OUT:
**SUPPRESS FRAGMENTATION
BY SUPPRESSING COOLING**
(keep $T > 10^4$ K) \rightarrow NEED
METAL-FREE GAS + H₂
dissociation by Lyman-Werner
UV bg above mean cosmic value
at $z > 2$

BUT METAL-FREE GAS
UNREALISTIC CONDITION!

(a) Metallicity $> 10^{-5}$ solar
reached at $z > 10 \rightarrow$ sufficient
to trigger rapid cooling esp. in
presence of dust (Omukai et
al. 2008).

(b) Weak inflow rates < 1 Mo/yr
(Wise et al. 2008; Regan & Haehnelt
2009, 2010)

Not enough to assemble supermassive clouds/SMS
Indeed no self-gravitating compact object forms



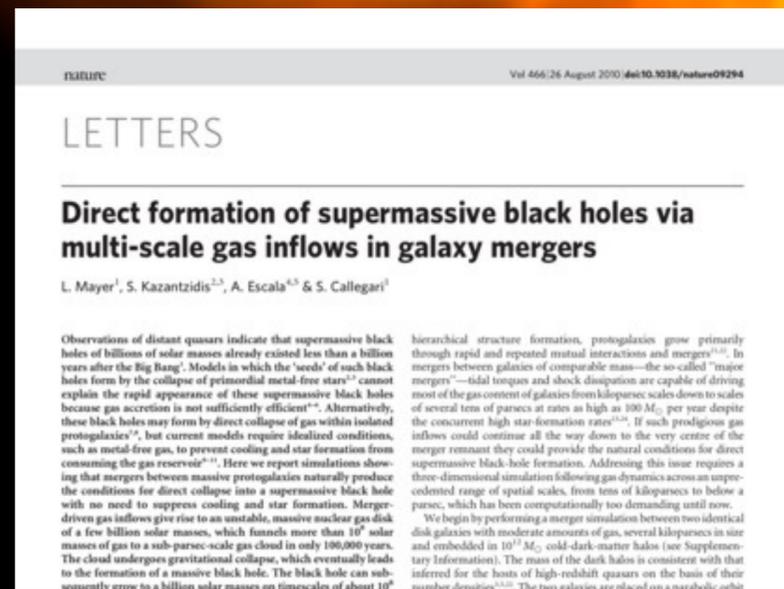
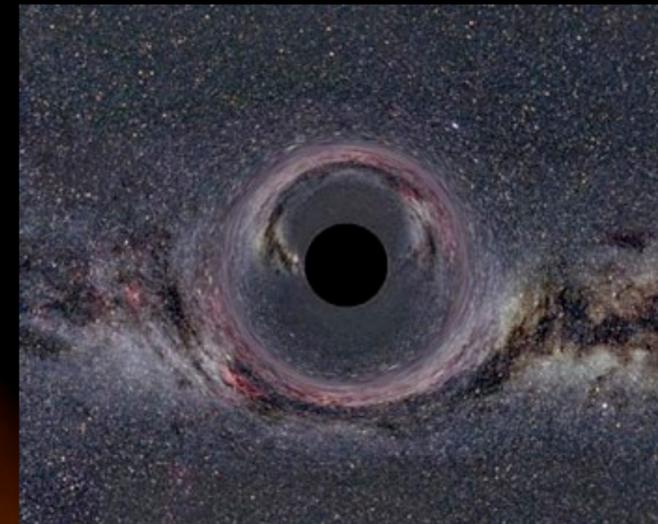
**Metal-free protogalaxy
simulation**
Regan & Haehnelt 2009

Formation of supermassive black holes by direct gas collapse in galaxy mergers



Lucio Mayer

University of Zurich



Collaborators:

Stelios Kazantzidis (CCAPP Ohio State Univ.)
Simone Callegari (Univ. of Zurich)
Andres Escala (KIPAC Stanford/UCHile)
Silvia Bonoli (Univ. Zurich)

Direct gas collapse model: brief intro

Rapid formation of massive BH seed --- mass $M_{\text{BH}} \sim 10^5 - 10^9 M_{\odot}$

If happens early ($z > \sim 8-10$) can explain high- z QSOs ($M_{\text{BH}} > 10^9 M_{\odot}$) without requiring the continuous Eddington accretion needed for $< \sim 100 M_{\odot}$ Pop III (Volonteri & Rees 2006)

Simulations show Pop III seeds accrete well below Eddington, eg Johnson & Bromm 2006; Wise et al 2008; Milosavljevic et al. 2010) due low density gas plus their own radiative feedback

I - Gas inflow in galaxy from kpc to $\ll 1$ pc scales to form supermassive gas cloud ($M > 10^6 M_{\odot}$) - need efficient loss of angular momentum in galactic disk gas across many spatial scales (eg Lodato & Natarayan 2006)

II - Depending on mass and internal rotation of supercloud (T/W) two pathways:

(a) supermassive cloud collapses dynamically and globally into massive black hole with $M_{\text{BH}} \sim M_{\text{cloud}}$ due to radial GR radial instability (Fowler & Hoyle 1966; Zeldovitch & Novikov 1972; Baumgart & Shapiro 1999; Shibata & Shapiro 2002; Saijo & Hawke 2009) ---> direct formation of SMBH

(b) forms a short-lived ($> \sim \text{Myr}$) supermassive star collapsing into BH at the center due to catastrophic neutrino cooling (Begelman et al. 2006; Begelman 2008; Begelman & Volonteri 2010). Even if BH initially only 10-100 M_{\odot} it accretes super-Eddington from a pressure-supported convective envelope powered by BH accretion energy ("Quasi-star") reaching $> 10^{4-5} M_{\odot}$ before cloud dispersal in a few Myr ---> formation of massive BH seed

This talk: how can step (I) be achieved?

TIMESCALE FOR SUPERMASSIVE CLOUD ASSEMBLY: REQUIRED GAS INFLOW RATE

Simple argument: a supermassive star ($M_{\text{star}} \gtrsim 10^6 M_{\odot}$) has short lifetime ($t_{\text{life}} \sim 10^6 \text{ yr}$) must be assembled on $t_{\text{form}} < t_{\text{life}}$

----→ Characteristic gas inflow rate to feed the cloud $dM_{\text{g}}/dt > M_{\text{star}}/t_{\text{life}} > 1 M_{\odot}/\text{yr}$ for $M_{\text{star}} \gtrsim 10^6 M_{\odot}$
(Begelman 2008)

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HOW DO WE GET SUCH HIGH GAS INFLOW RATES
AT $< \text{pc}$ scales?

---->Gravitational torques in self-gravitating, marginally unstable protogalactic disk

(bars-in-bars, spiral modes, see eg [Begelman et al. 2006](#), [Lodato & Natarayan 2006](#); [Levine et al. 2009](#))

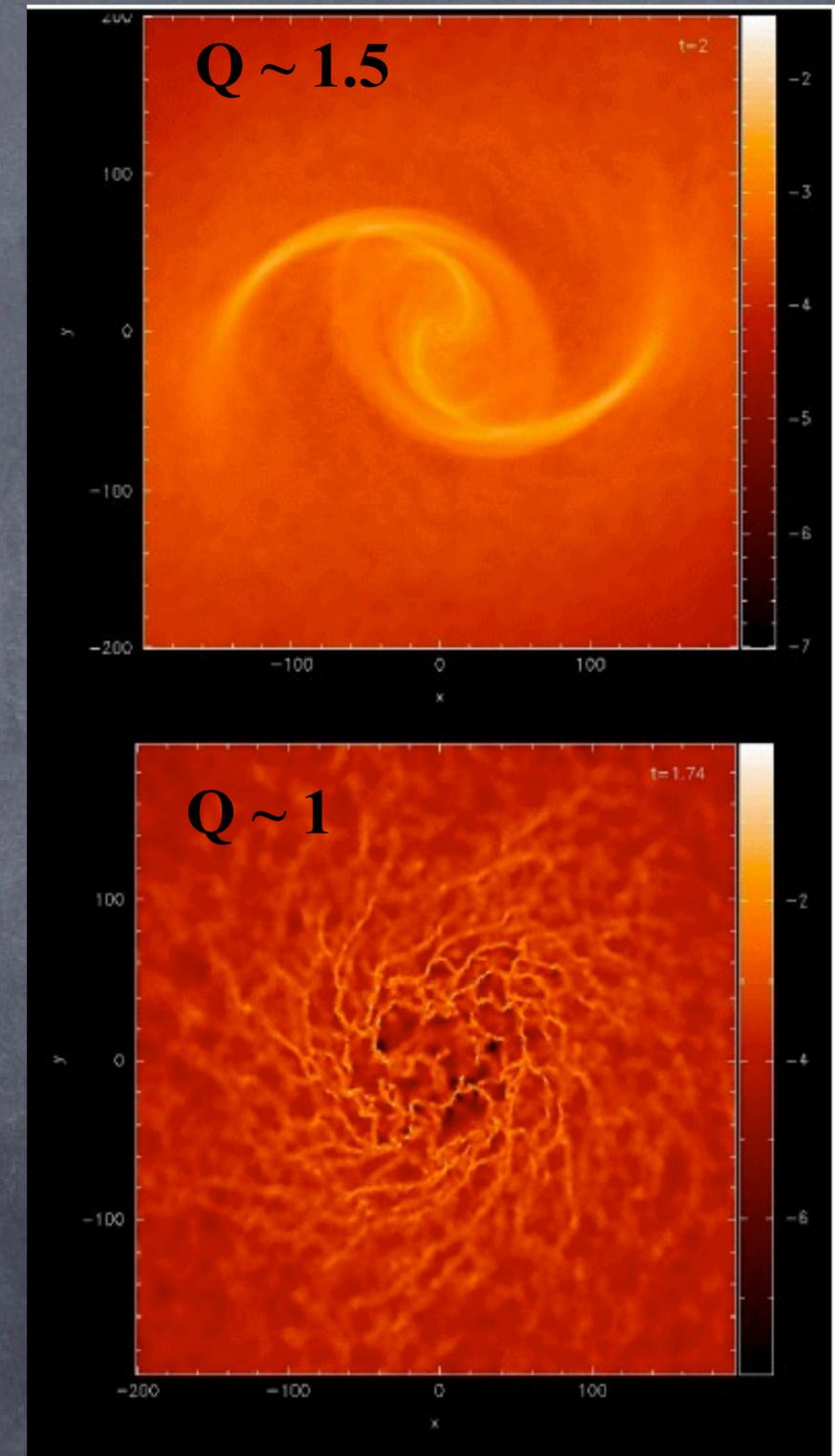
Needs massive but "warm" disk (Toomre $Q \sim 1.5-2$)

How do we keep the disk "warm" and stable?

Standard way: suppress molecular cooling and metal cooling below 10^4 K to keep $Q > 1$, avoiding fragmentation and star formation (otherwise gas makes stars rather than BH seed --- star formation bottleneck)

-- potentially can work at very high redshift ($z > 15$) with very low metallicity gas, perhaps requires proximity with massive star forming galaxies shining with high LW flux dissociating H_2 ([Dijkstra et al. 2009](#); [Agarwal et al. 2012](#))

-- characteristic host protogalaxy mass small ($\sim < 10^8 M_\odot$), a potential problem since inflow rate $dM_{\text{gas}}/dt \sim V_{\text{halo}}^3/G \sim M_{\text{halo}}/G < \sim 1 M_\odot/\text{yr}$ neglecting residual angular momentum (roughly consistent with simulations of [Wise et al. 2008](#), [Regan & Haehnelt 2009](#))



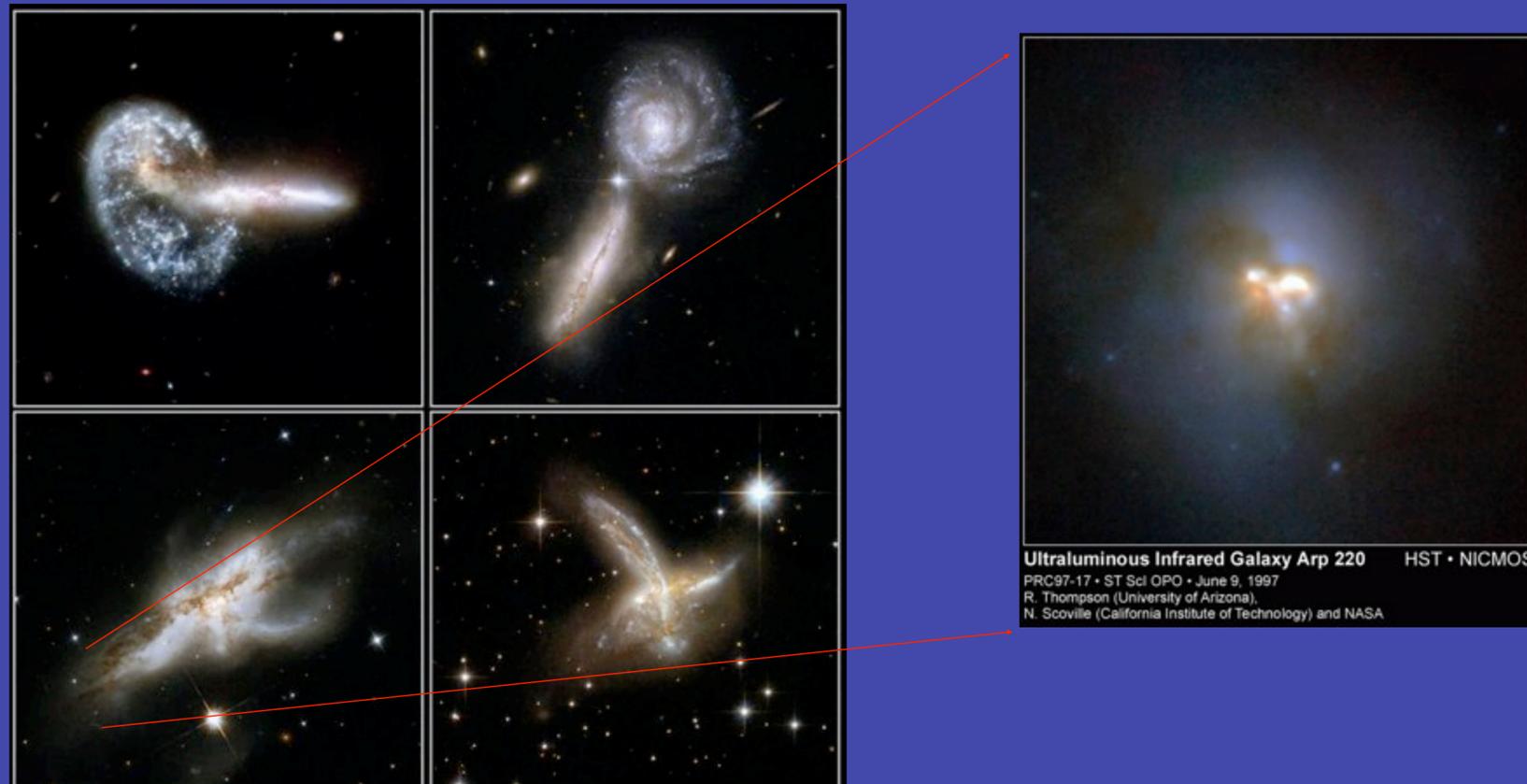
TOWARDS NEW BH FORMATION SCENARIO: MASSIVE MULTI-SCALE GAS INFLOWS IN GALAXY MERGERS

-Galaxy mergers are known to trigger the strongest gas inflows in galaxies at 100 pc- 1 kpc scales (due to tidal torques and shocks extracting angular momentum)
----> simulations show $dM/dt > 100 M_{\odot}/yr$ (eg Kazantzidis et al. 2005; Li et al. 2006), can sustain high SF rates in ULIRGs and sub-mm galaxies (eg Hopkins et al. 2008)

In mergers gas inflows effective, still most of the gas does not turn into stars!
from observations SF rate $\sim \epsilon_{sf} M_{gas}/t_{dyn}$, $\epsilon_{SF} = 0.01-0.1$, highest efficiencies occurring in high z merging systems (see eg Genzel et al. 2010, Tacconi et al. 2012)

-> slow gas consumption timescale compared to inflow timescale

$$t_{dyn}/\epsilon_{SF} \gg t_{dyn} \sim t_{inflow}$$



Bottom line: in mergers there is no “star formation bottleneck”, at least down to 100 pc scales, and there is a lot of low angular momentum gas...

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Can the merger-driven inflow continue all the way from 100 pc to \ll pc scales and form the precursor of a massive BH?

Gas-rich major mergers of massive proto-disk galaxies ($M_{\text{disk}} \sim 6 \times 10^{10} M_{\odot}$, $6 \times 10^9 M_{\odot}$ of gas at merger time) in $10^{12} M_{\odot}$ halos at $z \sim 8$

Resolution 0.1 pc in ~ 30 kpc volume using SPH particle splitting with EOS appropriate for nuclear starburst (Spaans & Silk 2000, 2005)

Galaxy halo mass consistent with abundance of high- z SDSS QSOs (Fan et al. 2006, Morlock et al. 2010) i.e. rare $3-4\sigma$ peaks at $z > 6$ (Volonteri & Rees 2006; Li et al. 2007)



Shown box size =
200 pc on a side
(galaxy cores a few Myr
before final collision)

60% of total gas
mass accumulated
within 200 pc due
to tidal torques and
shocks

Gas thermodynamics with effective equation of state (EOS):
polytropic with effective adiabatic index $\sim 1.1-1.4$

EOS based on model by Spaans & Silk 2005 (also Klessen et al. 2007) calibrated with radiative transfer calculation

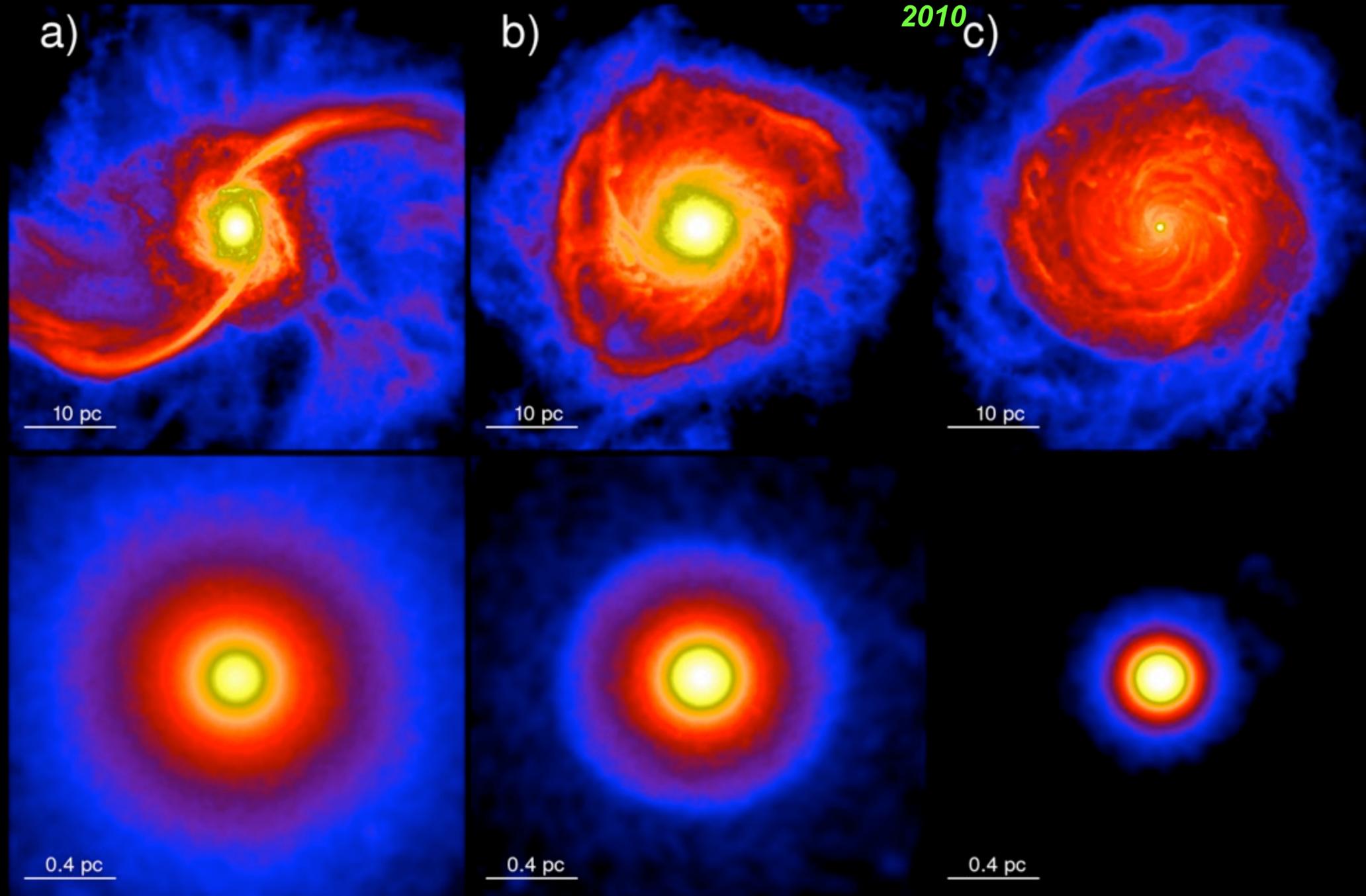
Accounts for thermal equilibrium between radiative cooling and heating (UV, IR from dust, cosmic rays) for density range 0.1 to 10^7 atoms/cc in dusty starburst with metal enriched gas (metallicity solar).

Multi-stage gas inflow down to sub-pc in gravitationally unstable circumnuclear gas disk forming in major merger

---> rapid formation of supermassive ($> 10^8 \text{ Mo}$) sub-pc scale gas cloud in only $\sim 10^5$ years after merger (SMBH precursor)

Below logarithmic density map spanning 10^5 yr after merger

Mayer, Kazantzidis, Escala & Callegari, Nature, 2010

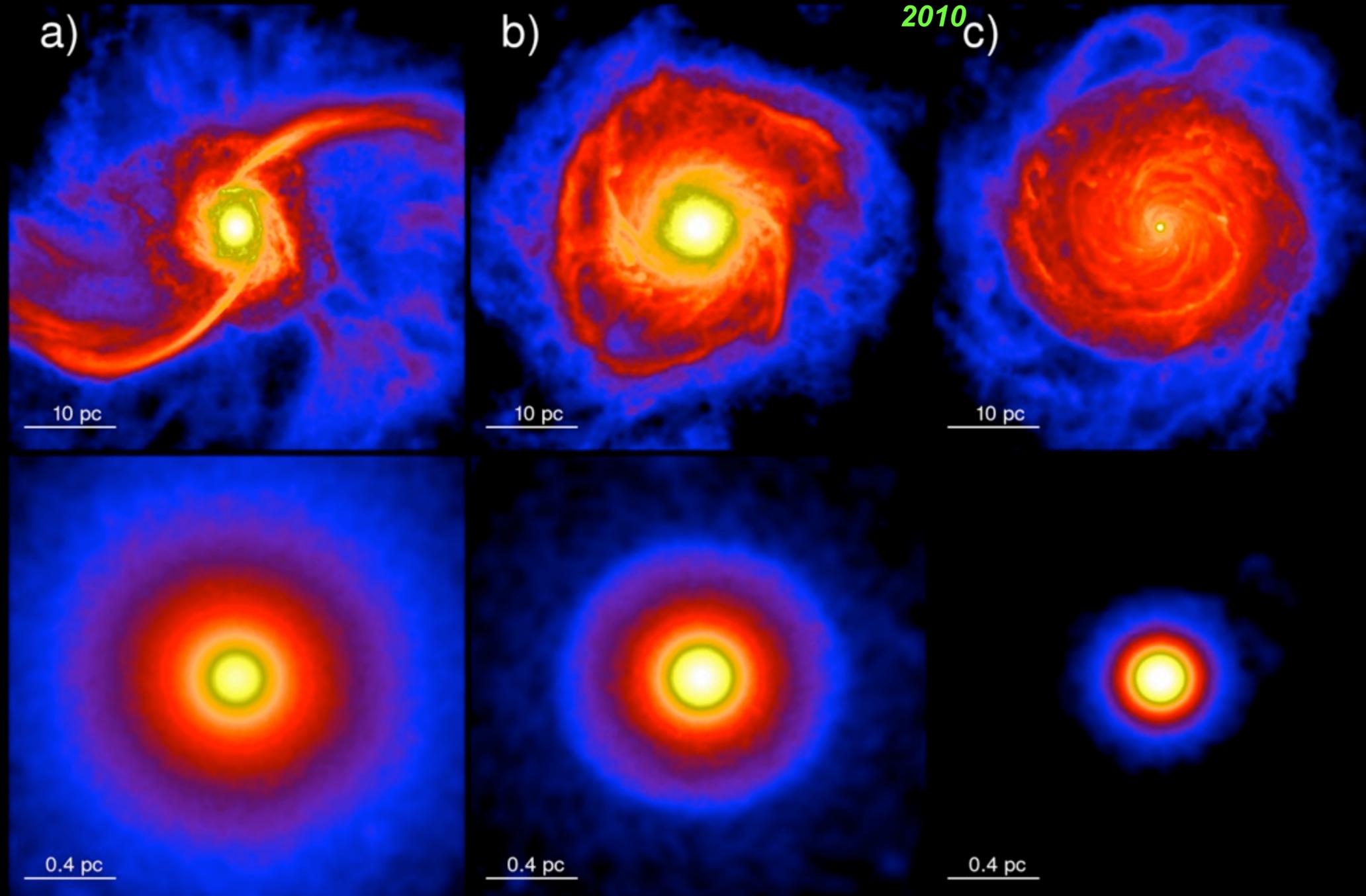


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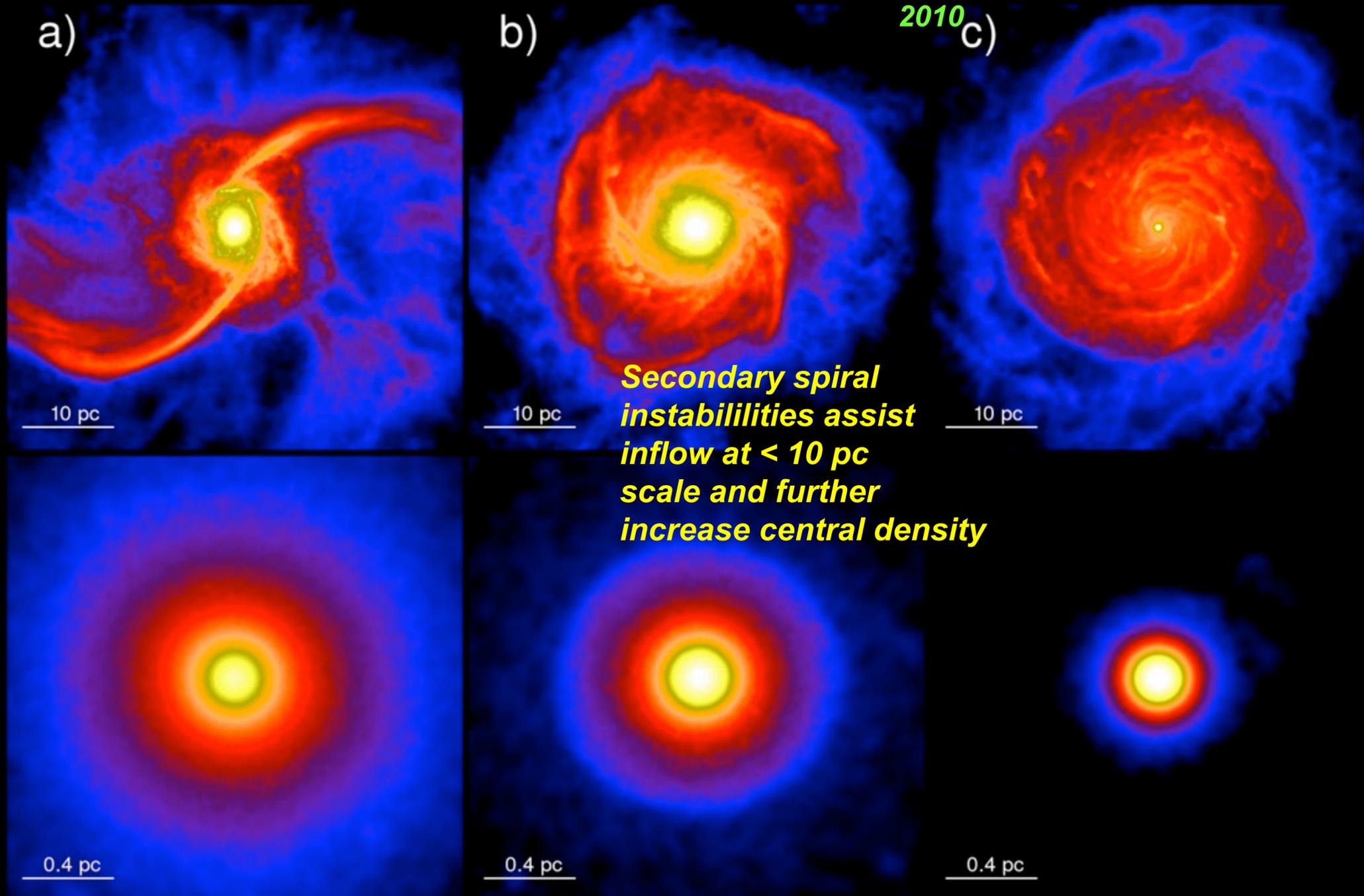
Large scale $m=2$ mode imprinted by galaxy collision starts inflow in nuclear disk

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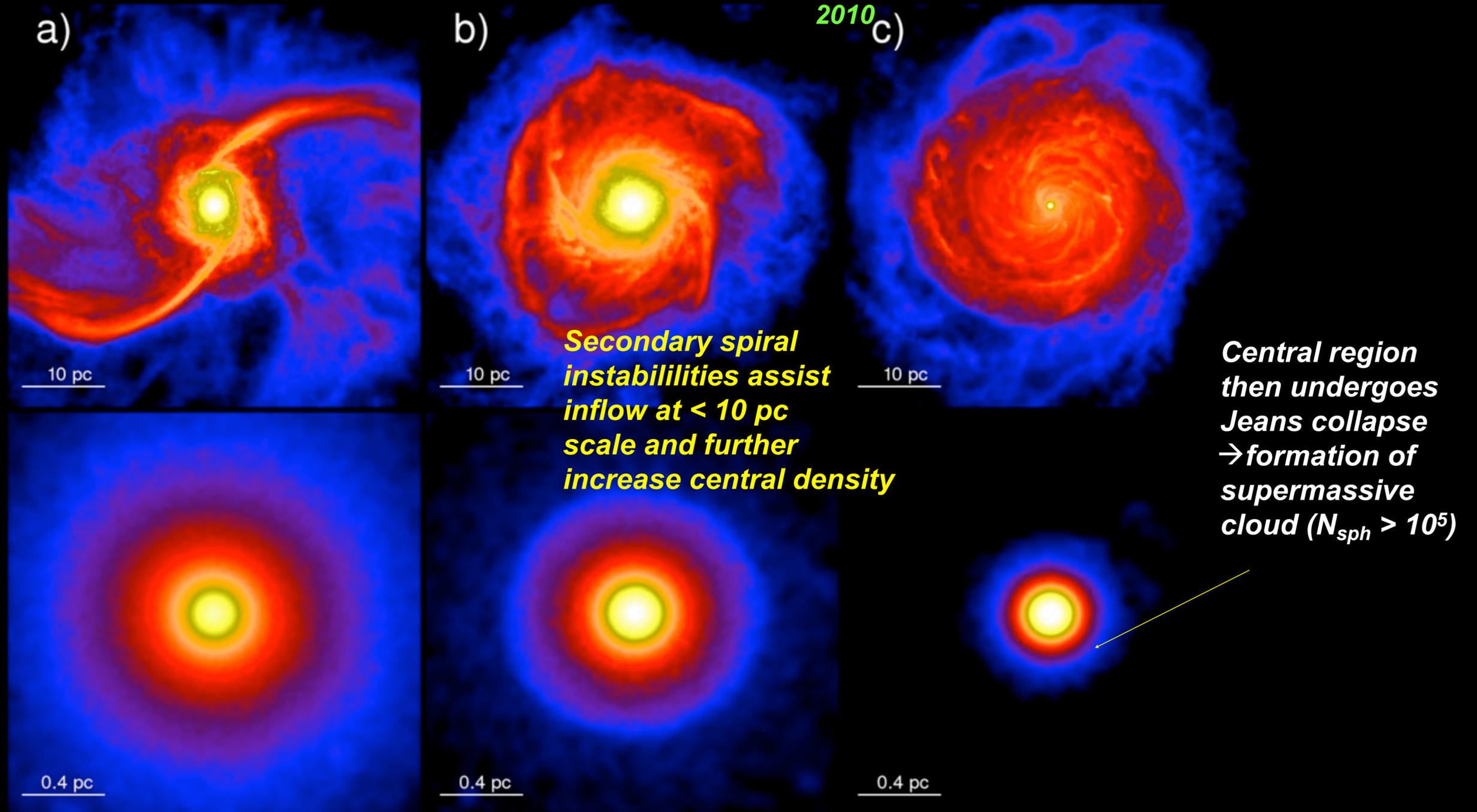
Secondary spiral instabilities assist inflow at < 10 pc scale and further increase central density

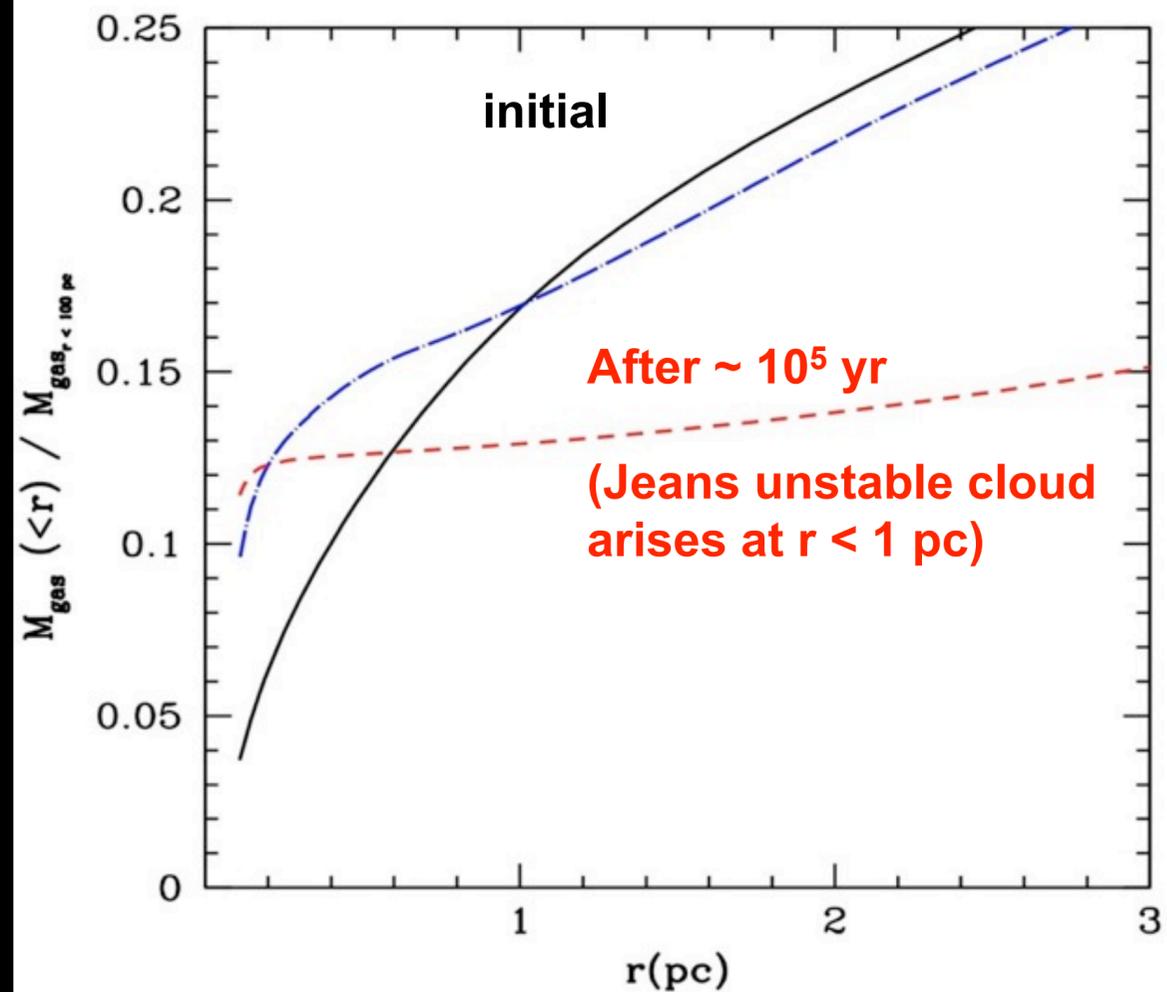
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In first 10^5 yr after merger:

Mass inflow rate

$\sim 10^4 - 10^5 \text{ Mo/yr}$

Star formation rate

$(\sim 0.1 \times M_g / T_{\text{orb}})$

$\sim 10^3 \text{ Mo/yr}$

\rightarrow gas inflow up to 2 orders of magnitude higher than star formation rate

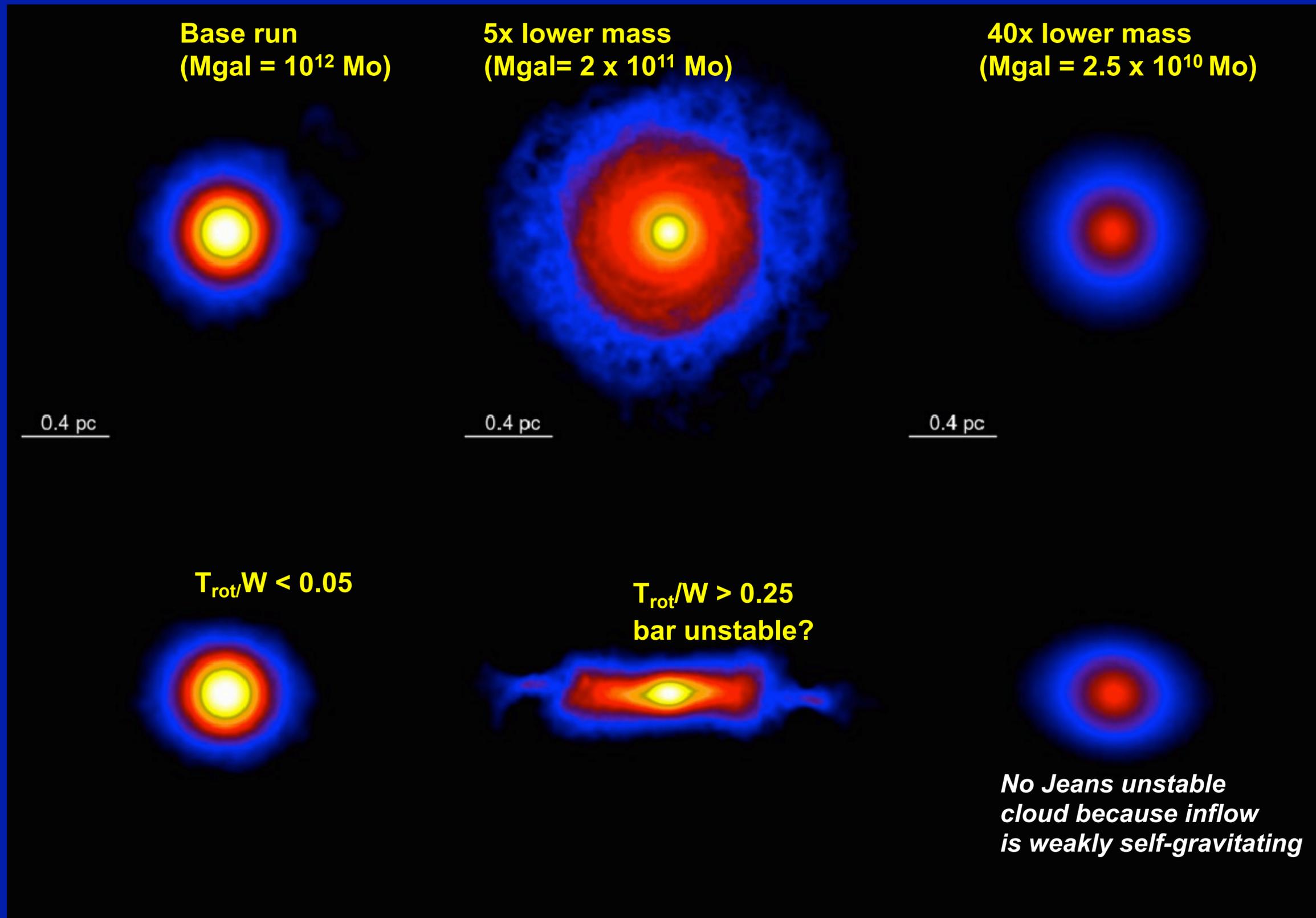
▪ **Supercloud Jeans unstable to resolution limit** – further collapse (a) into supermassive star or (b) directly into $> 10^8 \text{ Mo}$ SMBH via post-newtonian instability (route (b) requires $R \sim 640 GM/c^2 \sim 0.02 \text{ pc}$ for $M \sim 10^8 \text{ Mo}$ from numerical GR simulations results (Shibata et al. 2002; Saijo et al. 2009), for us $R_{\text{cloud}} \sim 0.5 \text{ pc}$)

▪ Assuming route (a) and, conservatively, that $> \sim 10^5 \text{ Mo}$ BH forms from ultimate collapse of SMS ($< 0.1 \%$ super-cloud mass!):

If initial black hole forms at $z \sim 9$ then can grow at $> \sim 0.7 \times$ Eddington rate to 10^9 Mo in $< 3 \times 10^8 \text{ yr}$, i.e before $z \sim 7$

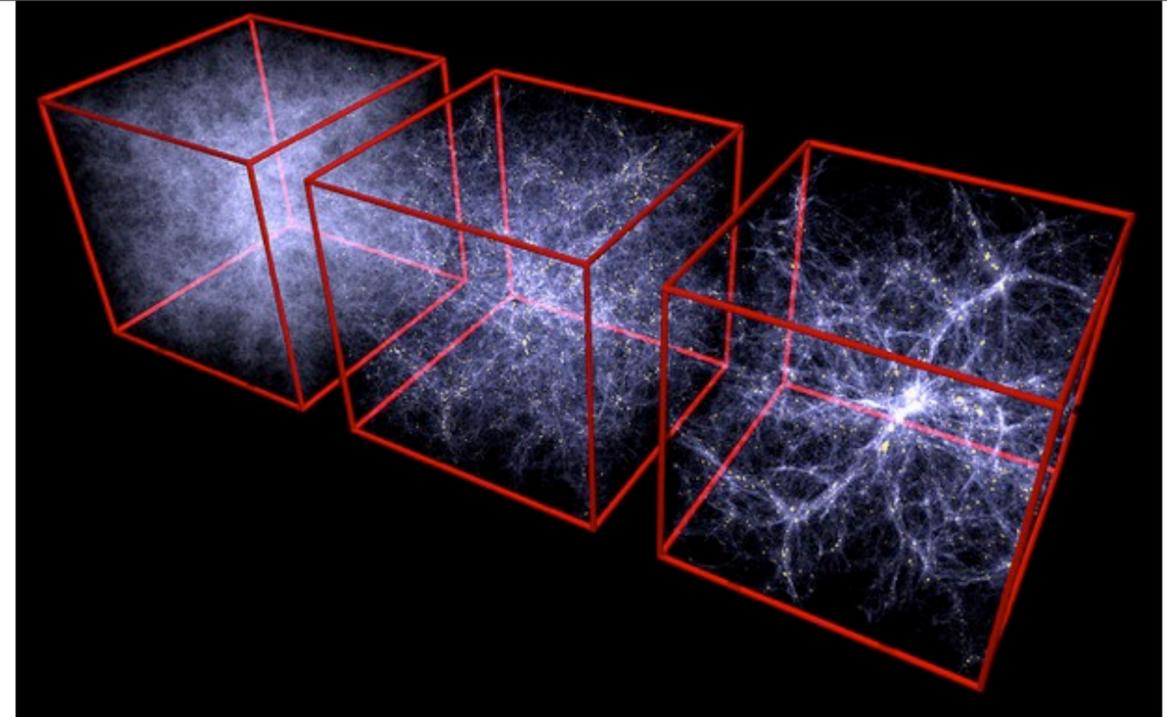
In low mass galaxies ($\sim 10^{10}$ Mo) no SMBH precursor forms

1:1 mergers between galaxies with a range of masses



Shown on left:
Logarithmic
gas density
maps

Embedding our formation scenario in the LCDM galaxy formation paradigm (Bonoli, Mayer & Callegari 2012)



We use the semi-analytical Munich model of galaxy formation (Croton et al. 2006; Bonoli et al. 2009), applied to the outputs of the Millennium Simulation



We have a full population of galaxies evolving in a cosmological framework that allows us to seek the BH seed formation conditions from hydro simulations and statistically test our scenario

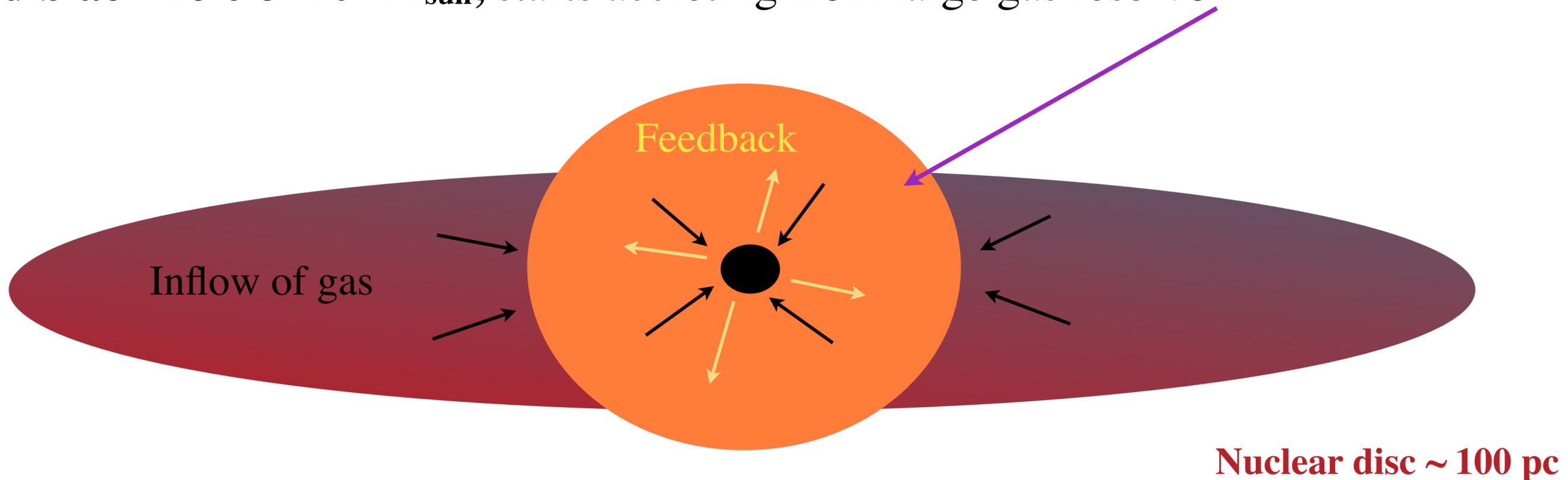
We follow the cosmological evolution of galaxies and their black holes:

- **PopIII seeds ($M = 1000 M_{\odot}$)** populate ALL newly formed galaxies
- **Direct collapse seeds ($M = 10^5 M_{\odot}$)** are formed during major mergers (and replace PopIII black holes), IF certain conditions implied by our simulations are satisfied

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- ✓ 3. No pre-existing black hole of $M_{\text{BH}} > 10^6 M_{\text{Sun}}$

Seed black hole of $10^5 M_{\text{Sun}}$, starts accreting from large gas reservoir

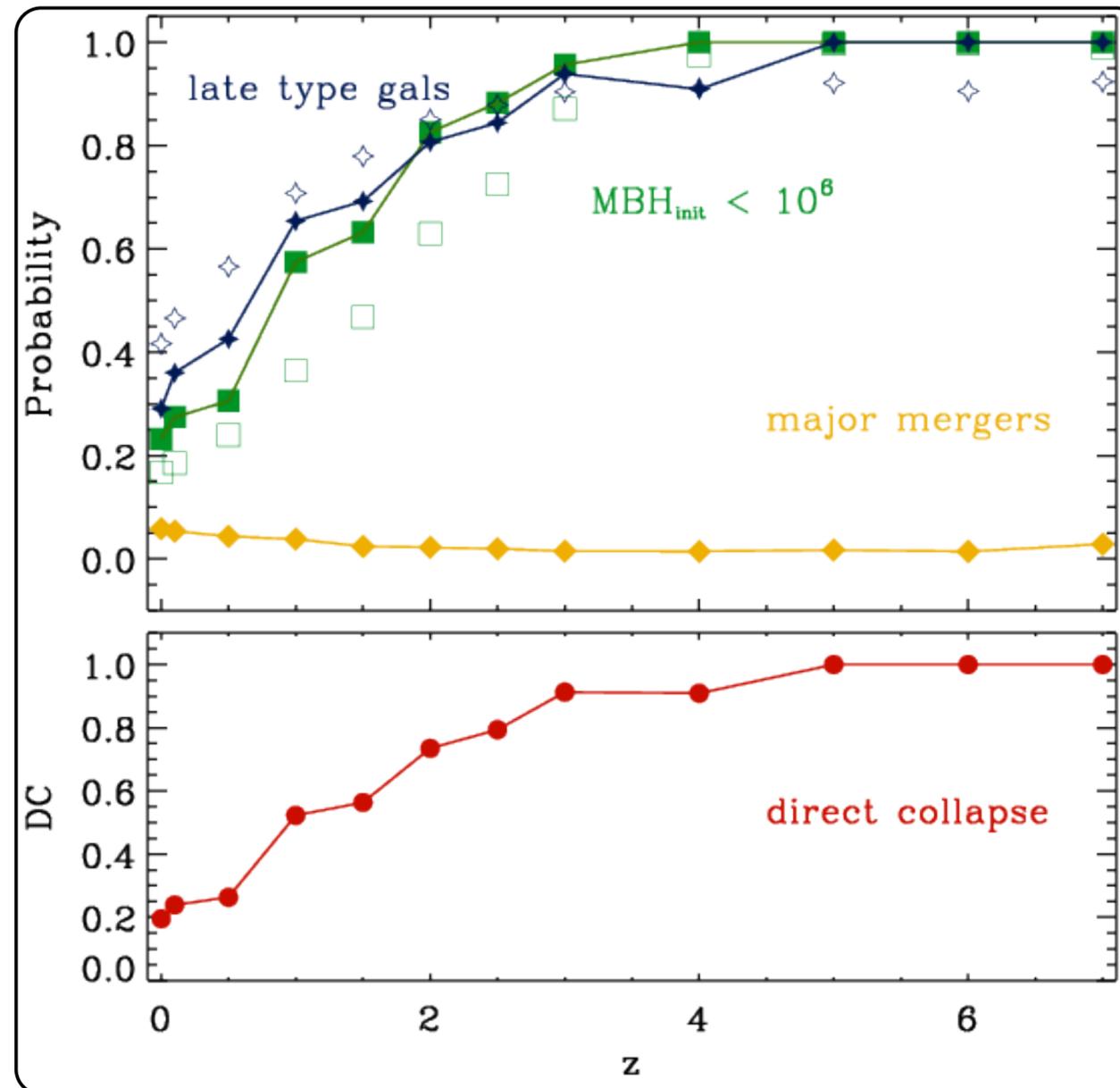


Self regulation: accretion stops once the feedback energy released by the black hole unbinds the reservoir (assumed isotropic thermal feedback with 0.05 coupling efficiency). BH will continue grow Eddington limited during subsequent mergers in the same way as Pop III seeds (a la [Croton et al. 2006](#))

Radius of the reservoir is a free parameter (0.1-1pc), determines its binding energy

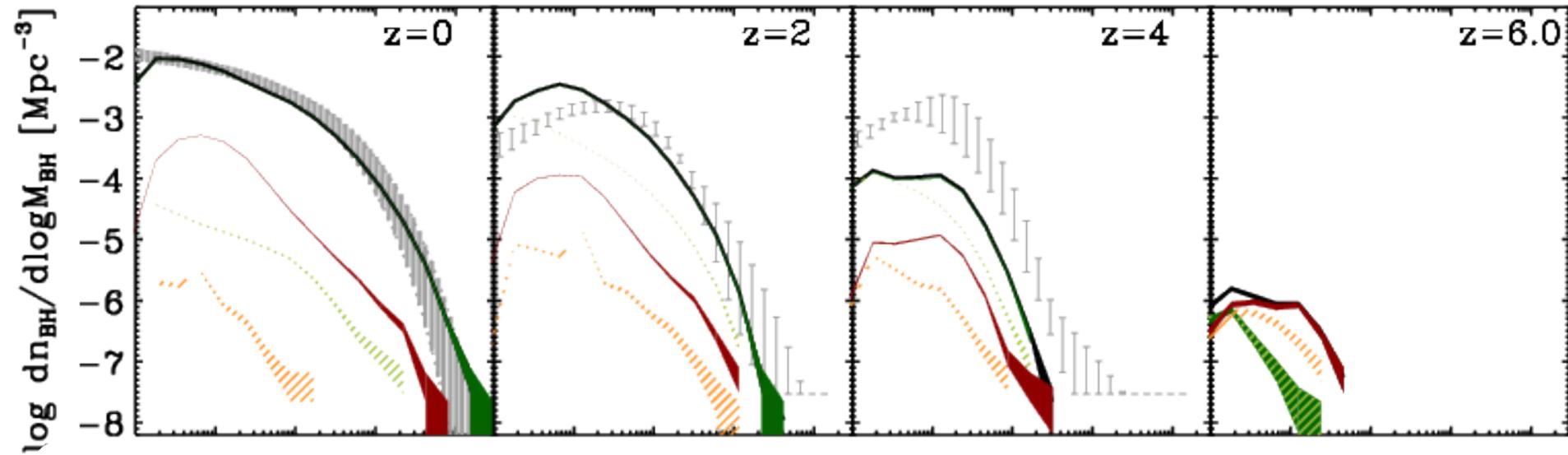
How frequent is our direct collapse route as a function of redshift?

- Above $z \sim 4$ all major mergers could lead to direct collapse
- Major merger events giving rise to direct collapse MBH seeds can happen even at low z (though large majority at $z > 3$)



Properties of the mass function (data from Merloni & Heinz 2005)

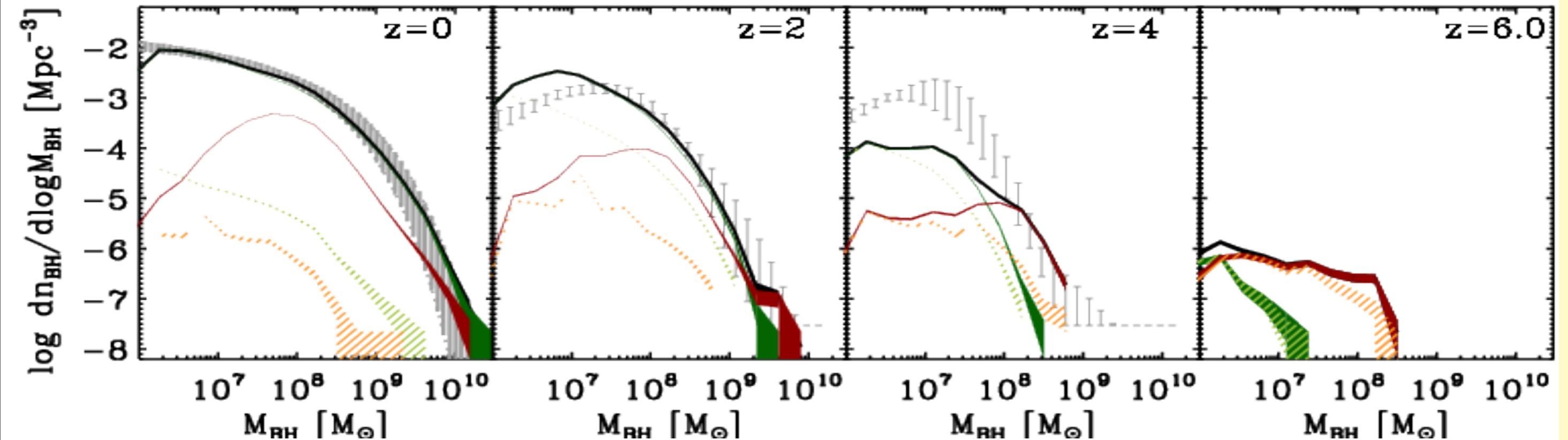
1 pc reservoir



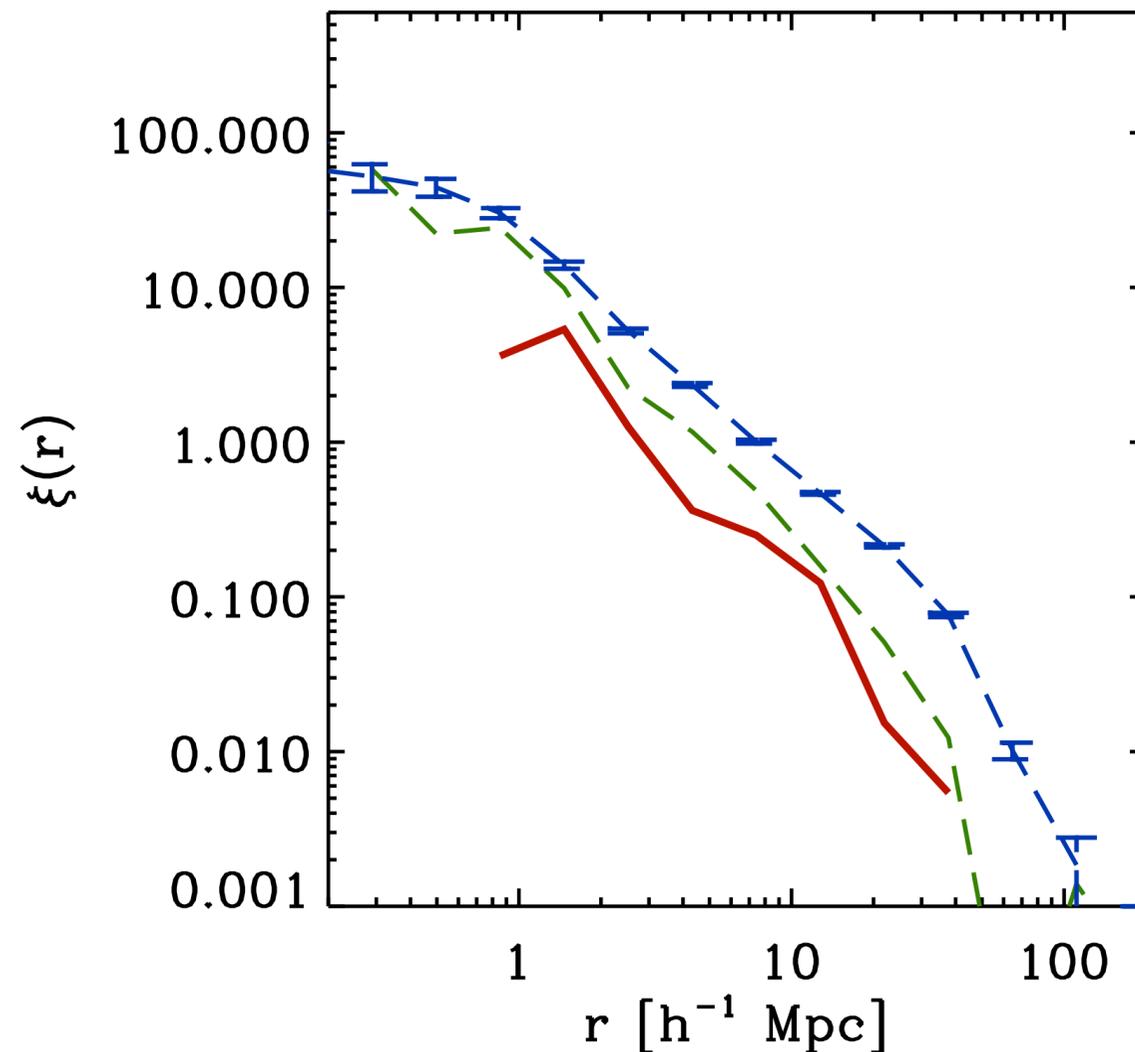
POP III

DC

0.1 pc reservoir



Clustering of galaxies forming direct collapse BHs at $z < 0.1$: two-point correlation function



**Red : Hosts of Direct Collapse
BH seeds formed at $z < 0.1$**

**Green: Recent major mergers
which do not form BH seeds
by direct collapse (but have Pop III
seeds) and have same galaxy stellar
mass distribution**

**Blue: Random Sample with
same stellar mass distribution
of host galaxies**

Low clustering amplitude relative to global BH population

because host galaxies had few or no mergers for nearly an Hubble time (otherwise Pop III seed grows and prevents direct collapse), i.e. fairly isolated objects

Qualitatively similar to low clustering of blue galaxies vs. global galaxy population (eg [Li et al. 2006](#))

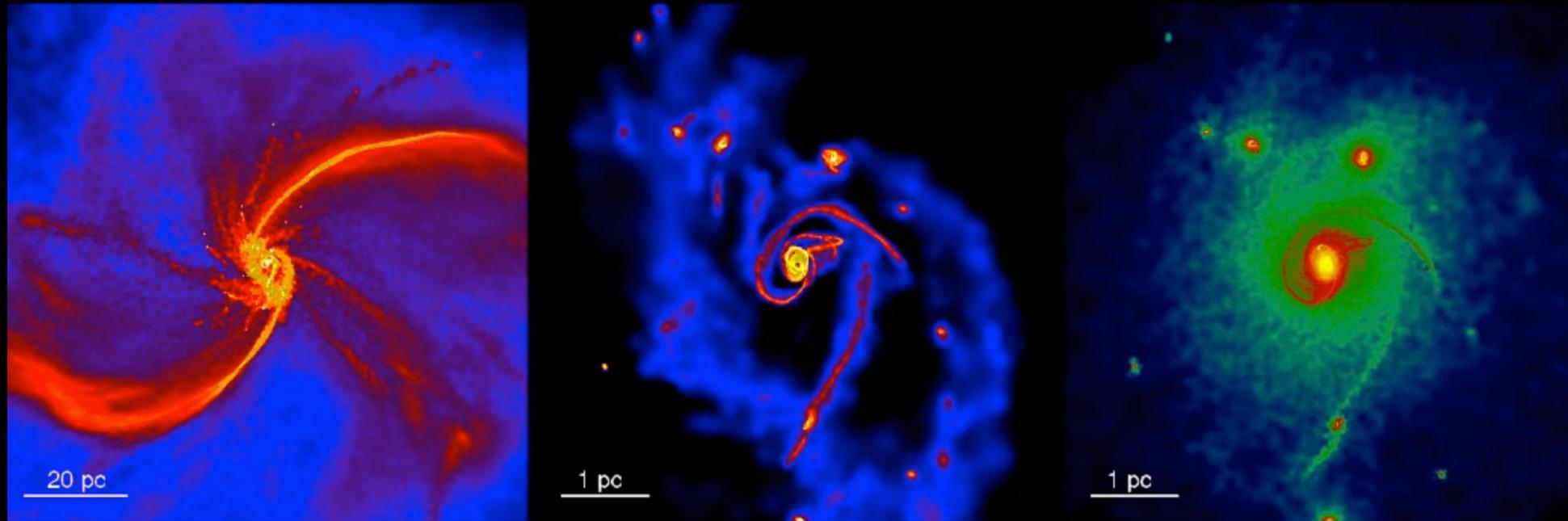
Open issues and implications

- Does direct collapse into a SMBH really occur after formation of supermassive cloud and *Which path does it take? Global post-newtonian instability? Supermassive star + quasi-star?*
-- Modeling of cloud collapse at even higher resolution w/post-newtonian effects and then interface with full General Relativistic simulations
-- Better characterization of cloud physical state – beyond EOS w/radiative transfer, neutrino diffusion in collapsing hot core etc..
- How does direct SMBH formation scenario depend on the structure/initial angular momentum content of merging galaxies? What is the role of gas turbulence?
Disks at high z clumpier and more turbulent than our ICs! Gravitoturbulence should aid collapse by extracing angular momentum further
- Does it stop working at low galaxy mass as our models with effective EOS suggest?
Likely yes --- in galaxies with $M < \sim 10^{10} M_{\odot}$ supernovae driven outflows should prevail over inflows, remove 2/3 of baryons (Governato, Brook, Mayer et al. 2010; Brook et al. 2011)

Predictions (simulation combined with SAM):

- BH formed by merger-driven collapse also at low z, and should have low clustering amplitude (those that form at high z are instead highly clustered as expected for high-sigma peaks)
- At $z > 2$ large deviations from the local $M_{\text{bulge}}-M_{\text{BH}}$: SMBH already in place while galaxy/ bulge has nearly an Hubble time left to grow
- If quasi-star phase precedes BH seed formation could be observable with JWST (blackbody emission at a few microns), although only very few per JWST field expected (see also [Volonteri & Begelman 2010](#)). At low z such events about an order of magnitude less frequent but gamma ray and radio emission could be detected if jets develop in quasi-stars, perhaps explaining unidentified sources in gamma-ray catalogs ([Czerny et al. 2012](#))

From EOS model to model with explicit radiative cooling and star formation

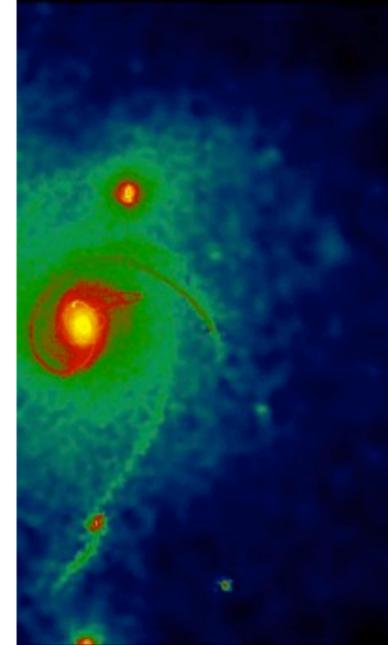
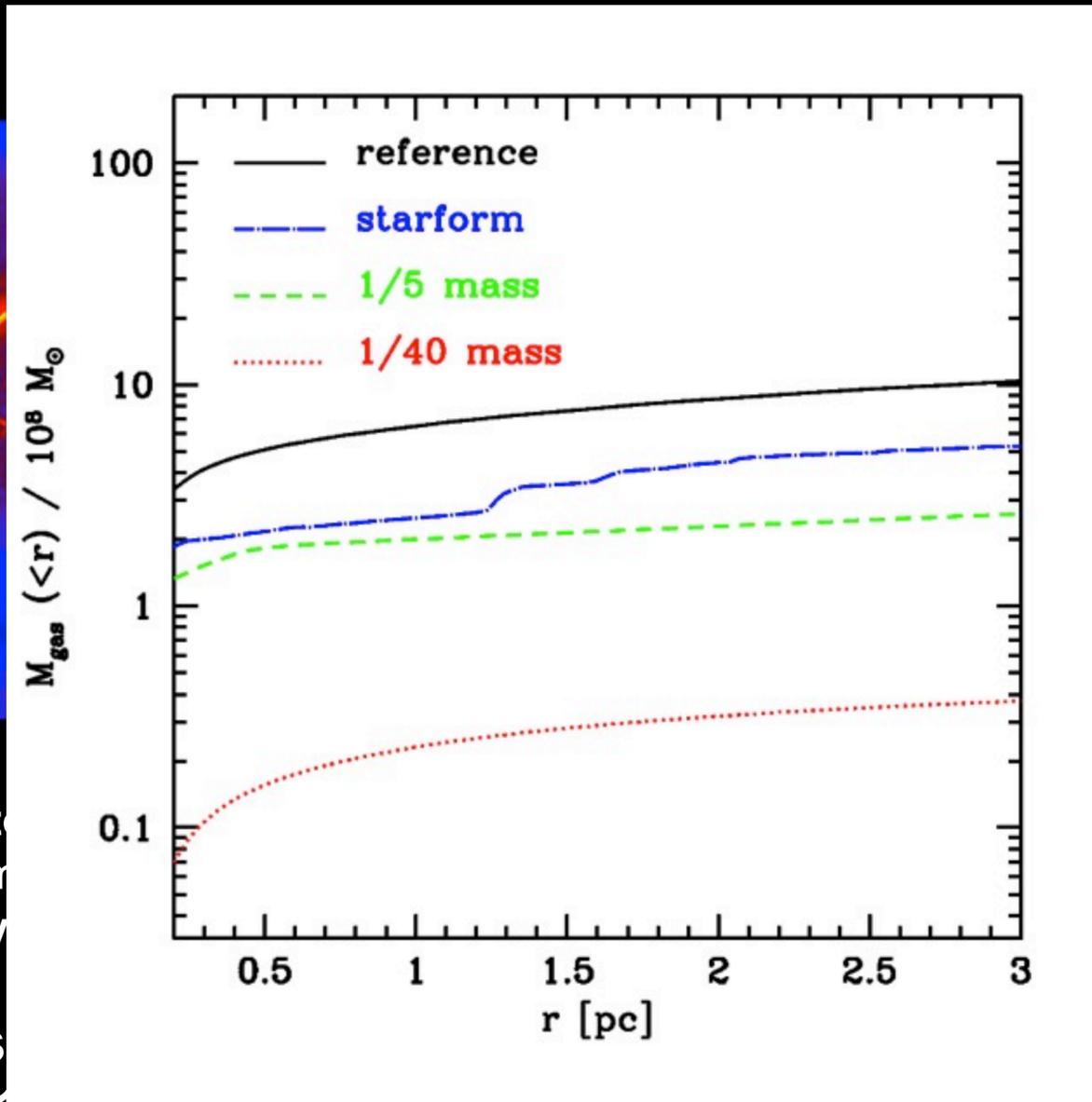
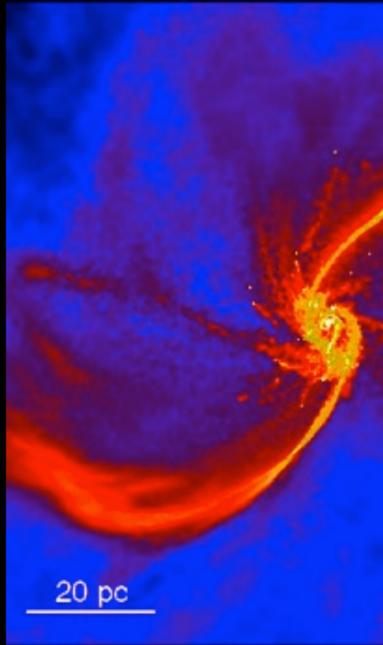


Gas cools radiatively and turns into stars above a density of 10^4 cm^{-3} + pressurization of medium to avoid spurious fragmentation below local Jeans length (no radiative transfer or heating by stellar/supernovae feedback, so max. fragmentation)

< 10^5 yr after the merger star formation has turned 30% of the nuclear gas disk into stars but $> 10^8 \text{ Mo}$ of gas still concentrates at $< 0.5 \text{ pc}$ in supermassive flattened cloud

-> **even stronger inflow than with EOS model** (gravitoturbulent regime, see **also Begelman & Shlosman 2010**)

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n)
 clear

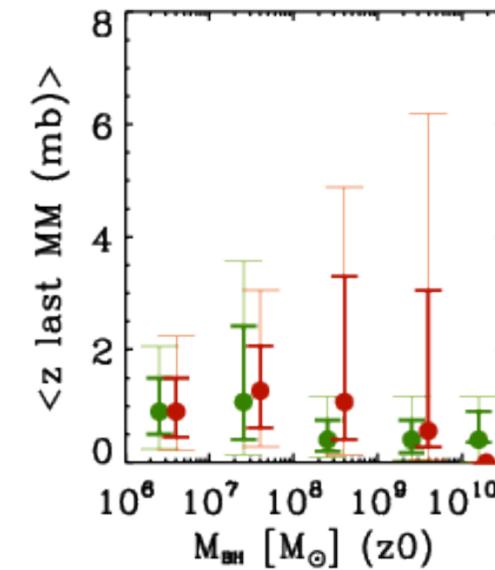
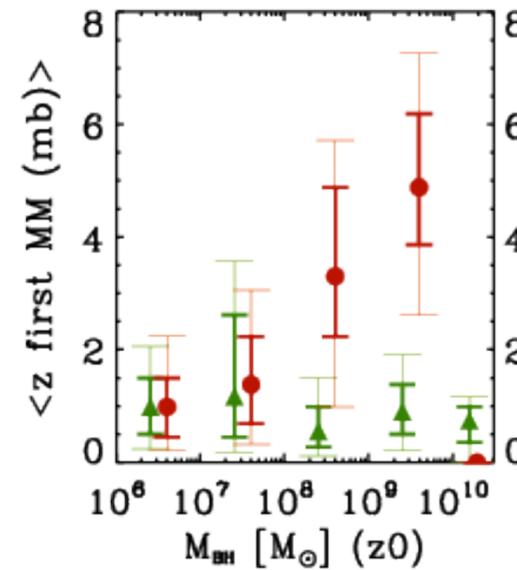
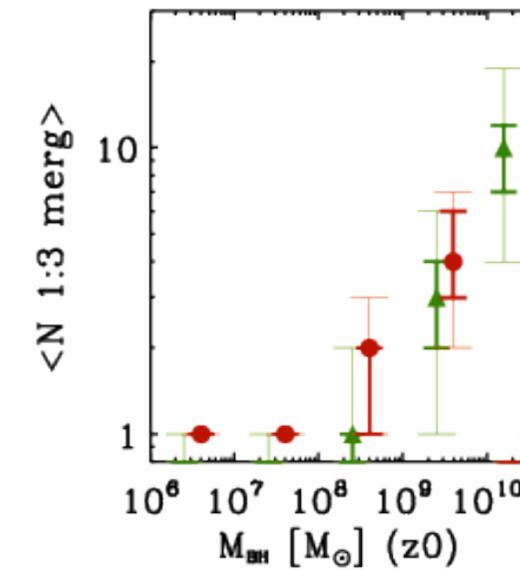
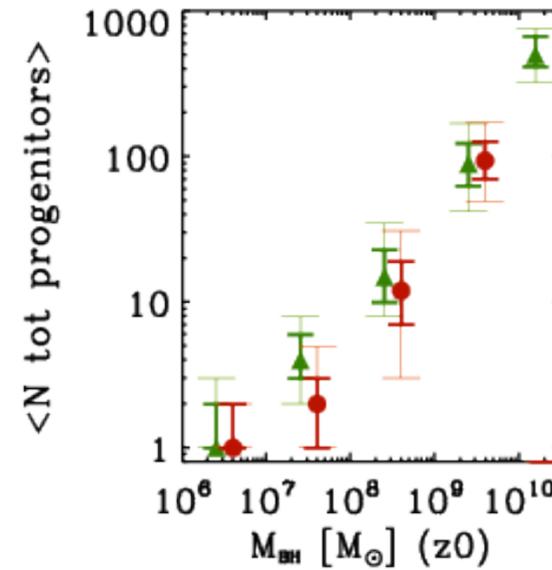
Features of merger histories for galaxy hosts of different BH seeds

POP III (small seeds)

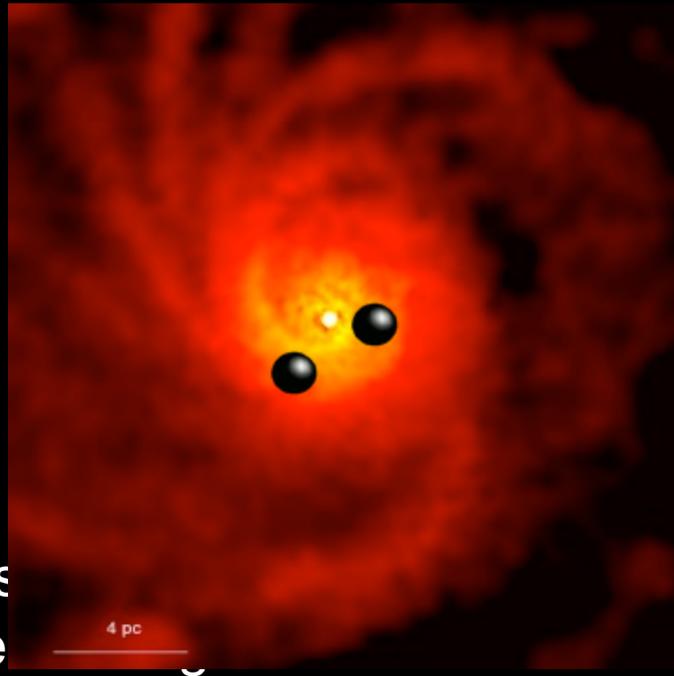
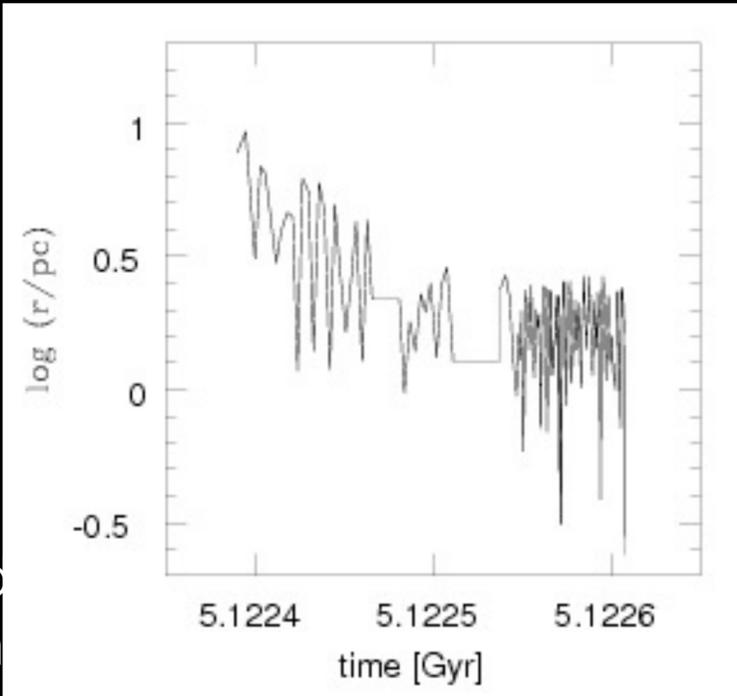
DC

other colors for other percentiles

Only clear distinction between seeding scenarios: hosts of direct collapse seeds have first major merger earliest



A binary SMBH stalls at ~ 0.5 pc in a nuclear disk capable of forming a central massive supercloud

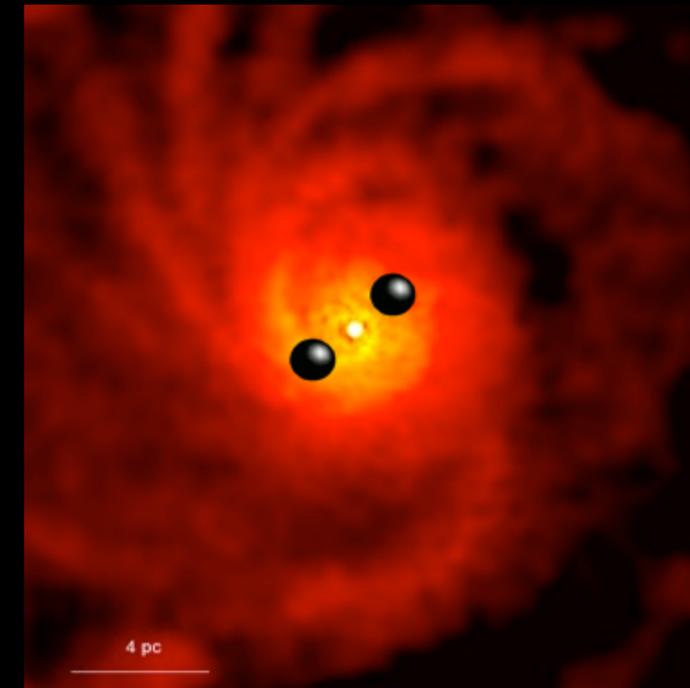
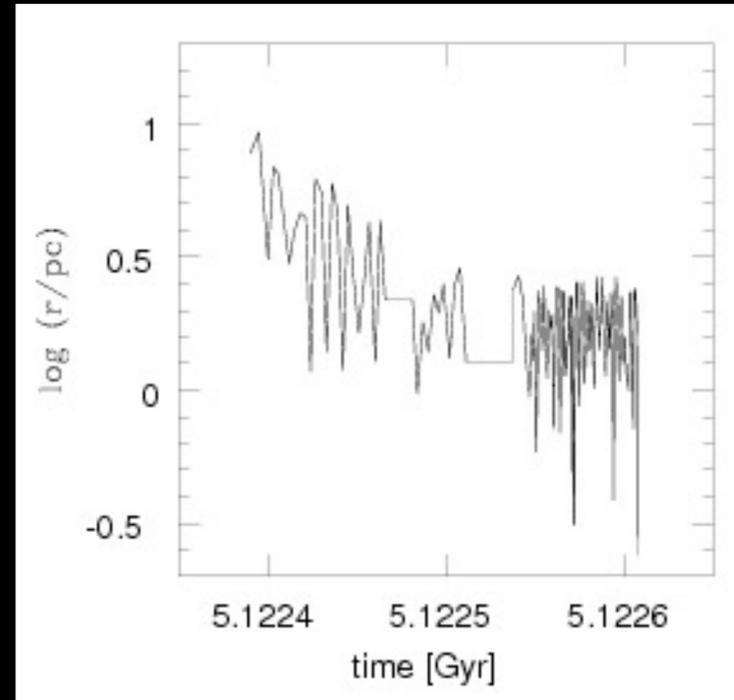


Dynamical friction from background gas is the local gas density and v_{bh} the relative velocity. Both SPH (Mayer et al. 2007) and AMR simulations (Chapon, Mayer & Teyssier 2011) show that binary of SMBH hardens down to about ~ 1 pc separation in ~ 10^6 yr.

But in less than 10^5 yr (1) density increases by x10 (the supercloud) at scales < 1 pc over ~ 10^5 years but decreases x10 just outside 0.5 pc +(2) v_{bh} increases because larger mass in the center (=supercloud)

*---> **df slows inefficient because $t_{inflow} < t_{df}$***

A binary SMBH stalls at ~ 0.5 pc in a nuclear disk capable of forming a central supercloud



But is a nuclear gas disk with a central supercloud a realistic configuration when two SMBHs are already present?

The answer is: probably not

Attractive scenario (to be investigated);

- *when no pre-existing black hole is present disk is violently unstable, drives a strong inflow and central supercloud collapse --→ massive SMBH seed formation in massive merging protogalaxies at $z > 5$*

(ii) when one or two massive black holes are already in place in the nuclear disk ($M_{BH} \geq 10^6 M_{\odot}$) they accrete gas and heat the disk via radiative feedback, stabilizing it against spiral instabilities and thus suppressing the central collapse

($Q > 2$ from Eddington limit accretion and 10% of accretion energy released as thermal/turbulent kinetic energy over about 10^8 yr)→

the disk profile does not become so steep and the binary can sink down to separations < 0.1 pc.

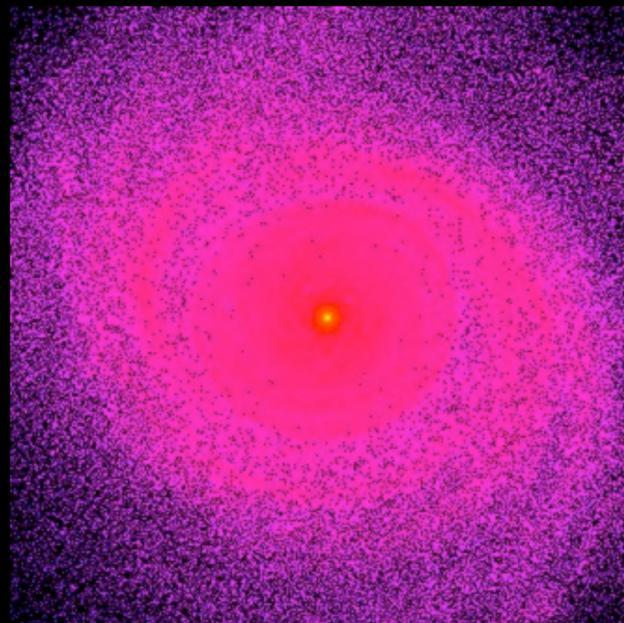
→ Binary SMBH coalescence successful + no formation of new SMBH seed

Which fate of the “supercloud”?

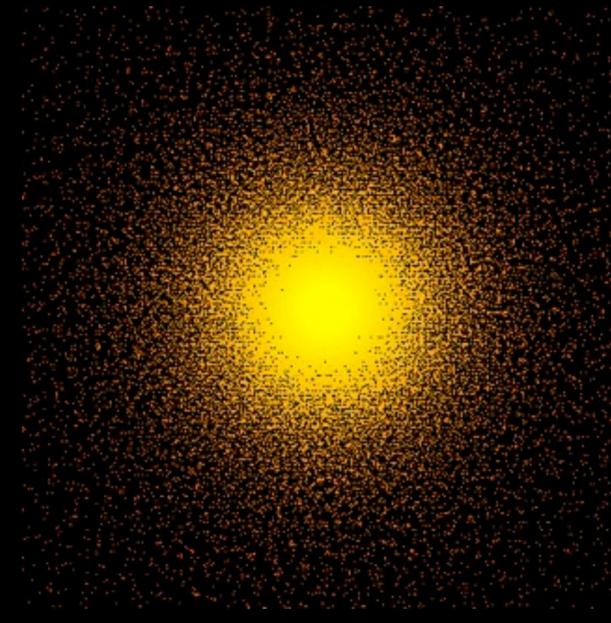
New ongoing simulation campaign to study supercloud collapse to post-newtonian regime (**Mayer, in prep.**)

First step; verification that cloud collapses continues below 0.1 pc in the newtonian case, including superclouds with highest angular momentum, by repeating simulations with 0.02 pc resolution (at even higher res PN corrections necessary)

3 pc box



0.05 pc box



Cloud evolved with $\gamma = 1.1$ and $\gamma = 4/3$ (likely more realistic, should be optically thick to its own radiation \rightarrow radiation pressure supported cloud)

After 2 free-fall times $\gamma=4/3$ cloud in sim with highest T_{rot}/W ($> \sim 0.25$) has turned into a core-disk envelope structure (no bar instability occurs)

Core contains $\sim 7 \times 10^7$ Mo, is ~ 0.04 pc in size and is still Jeans unstable at $t=2t_{ff}$ (end of sim)