

Tracing hierarchical galaxy formation with stellar dynamics: From nuclei to halos

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Dynamical imprint of galaxy formation

$$t_{relax} \approx \frac{NR}{(8 \ln \Lambda)v} \gg t_{Hubble}$$

-Violent relaxation in mergers is incomplete (Lynden-Bell 1967).

-Memory of ICs encoded in integrals of motion (IOM).

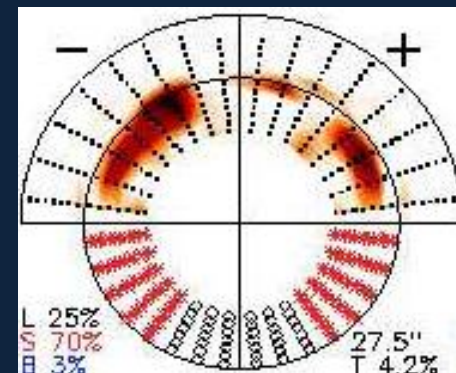
-Time-dependence/perturbations
-> diffusion in IOM phase space.

-Crossing of orbital boundaries
-> erasure of ICs (relaxation).

-To understand galactic structure and evolution, need to identify macroscopic groups of orbits that conserve qualitatively similar IOM.



Stellar orbits are the building blocks of galaxies (images: Poon & Merritt 2001).

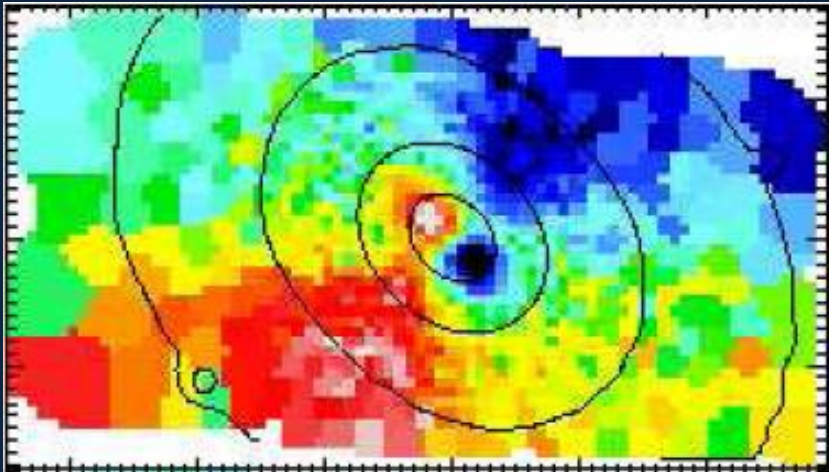


Orbital phase space need not be uniformly populated (image: van den Bosch et al. 2008).

Dynamical information in real galaxies



Image: <http://seds.org/~spider/ngc/ngc.cgi?NGC4365>



NGC 4365: SAURON image

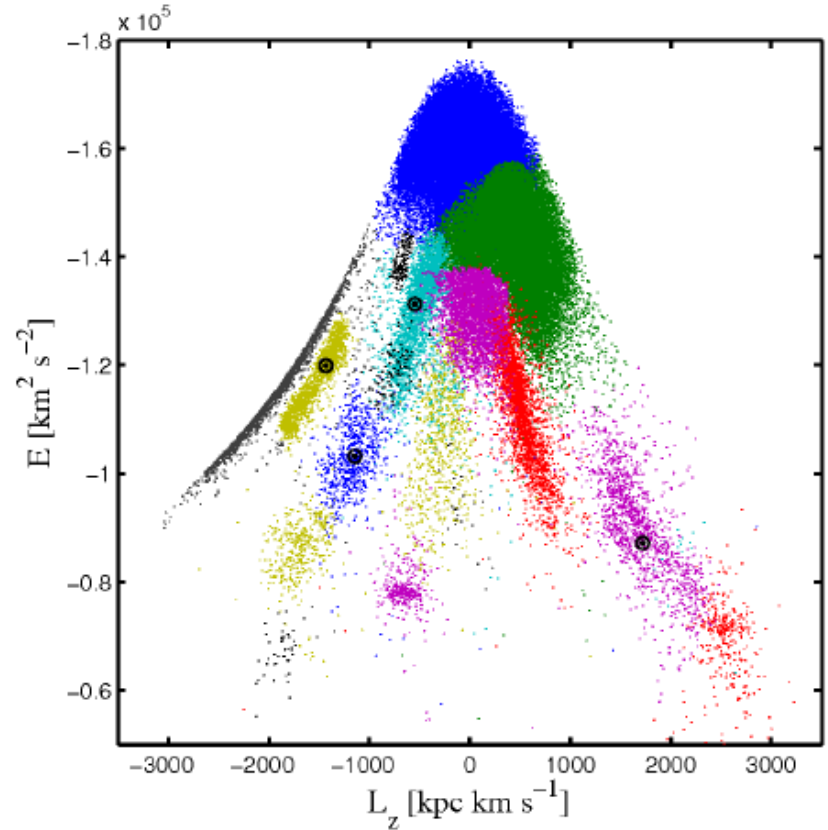
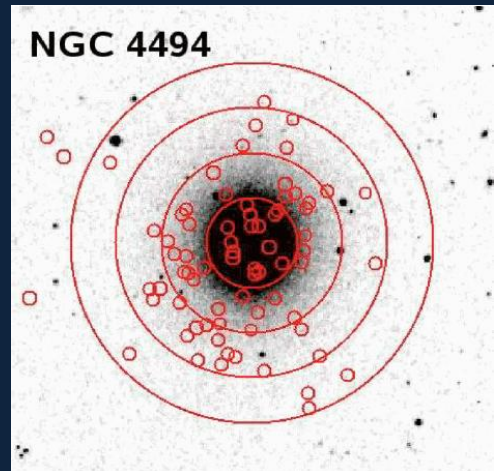
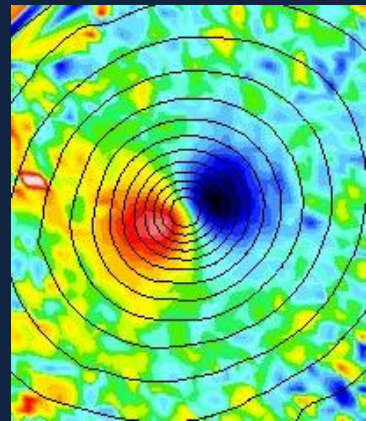
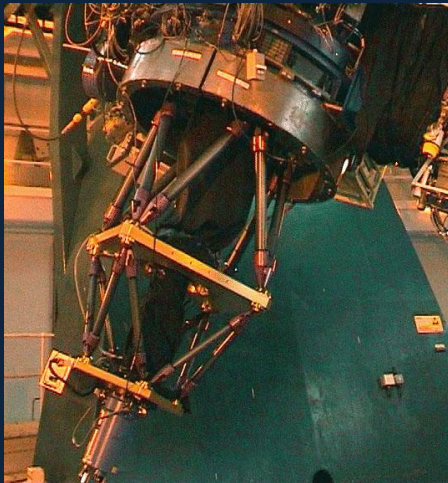


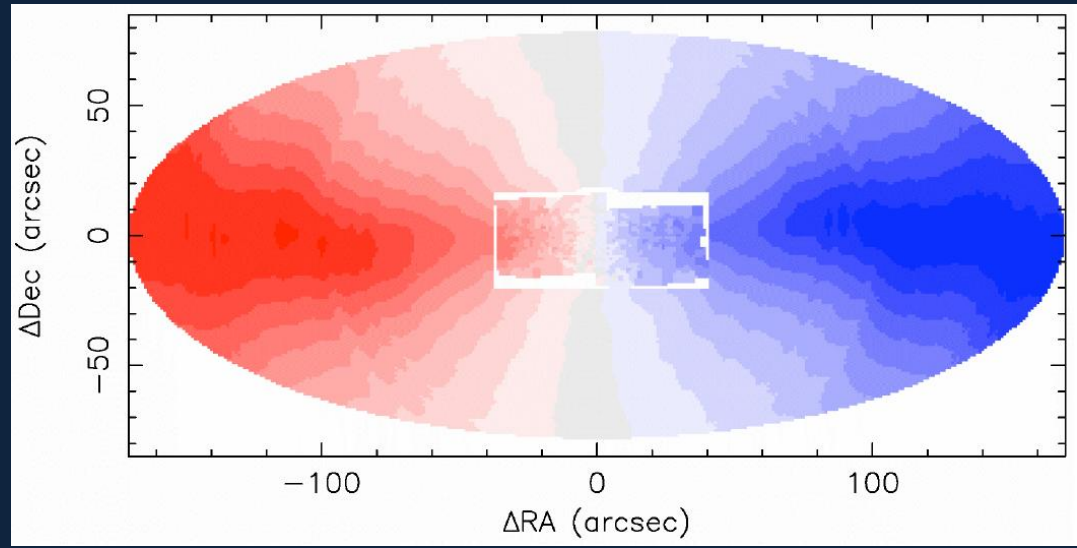
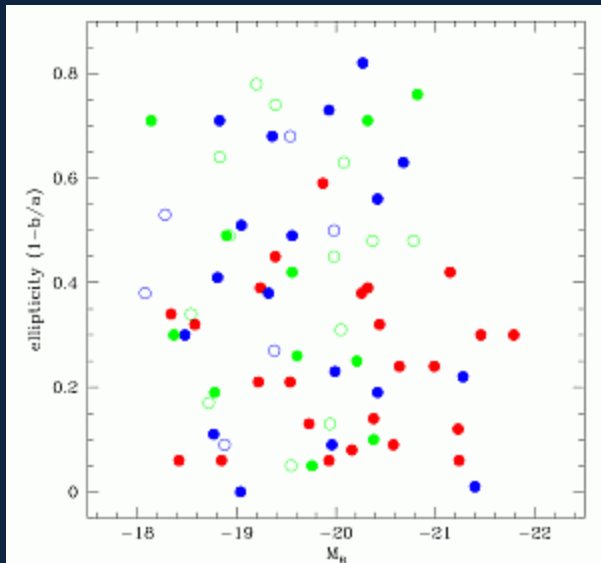
Figure 11. Distribution of stellar particles inside a sphere of 4 kpc radius at 8 kpc from the galactic centre in E vs. L_z space as would be observed by *Gaia*. The different colours show the groups identified by the Mean Shift algorithm. Black open circles denote those for which the time since accretion was successfully derived.

Gomez et al. 2010

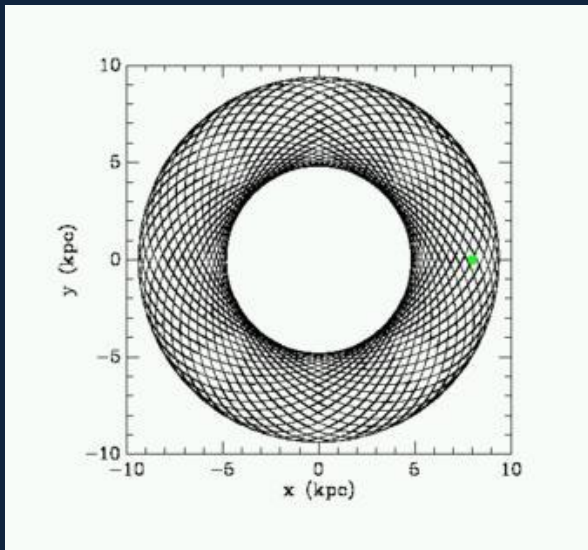
The stellar population of the NGC 4365 KDC is indistinguishable from that of the outer galaxy.



Proctor
et al. 2009



Spherical potentials

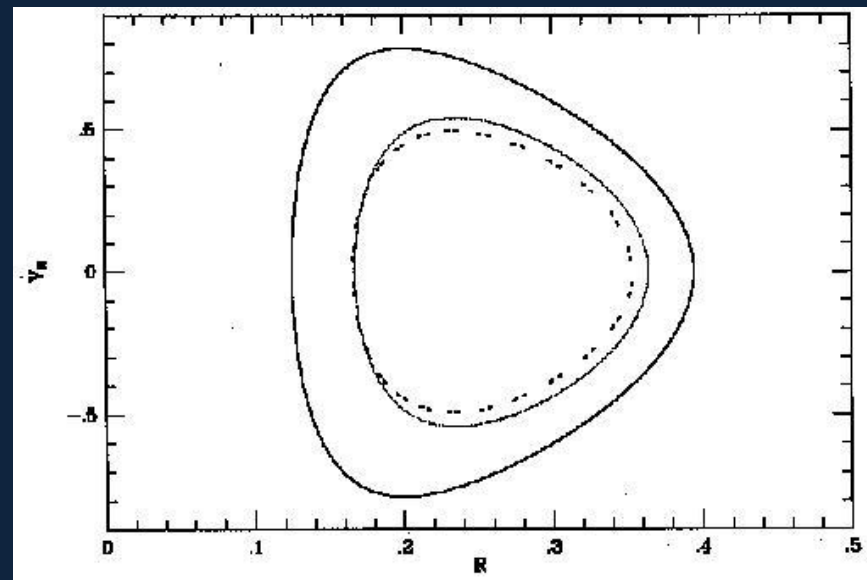
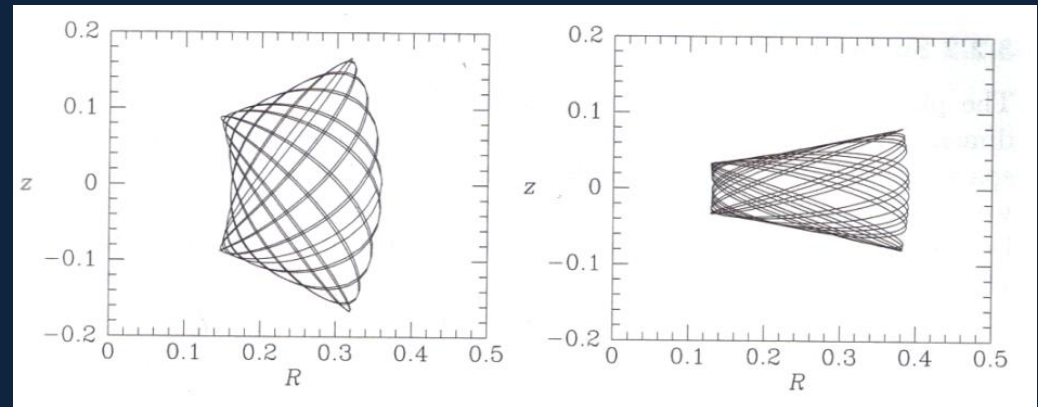


-Energy, angular momentum conserved \rightarrow orbit confined to plane

-Rosettes (precessing ellipses) uniformly fill annulus $r_0 < r < r_1$ over long times

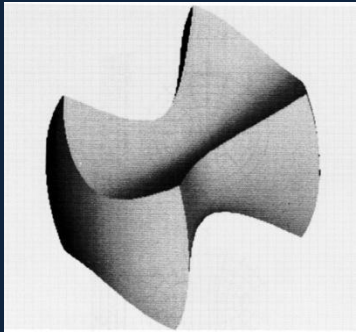
-Centrifugal barrier shields star from center of potential

Axisymmetric potentials

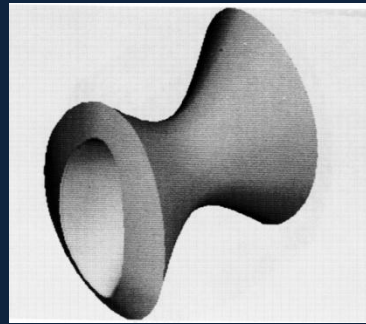


Binney & Tremaine 2008

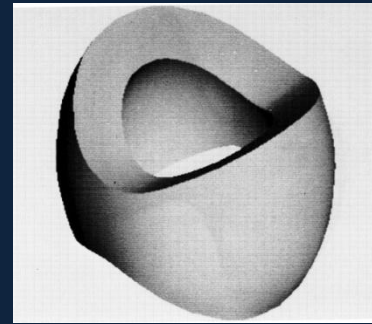
3D triaxial potentials



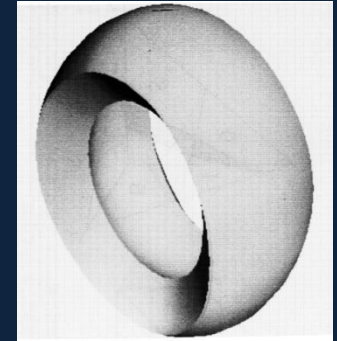
Box



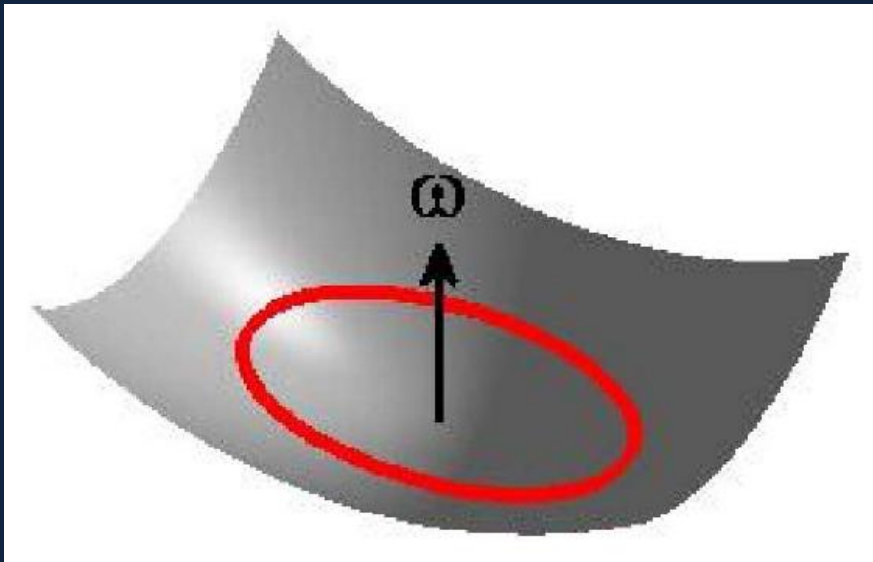
Inner long-axis tube



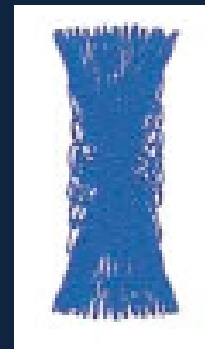
Short-axis tube



Outer long-axis tube

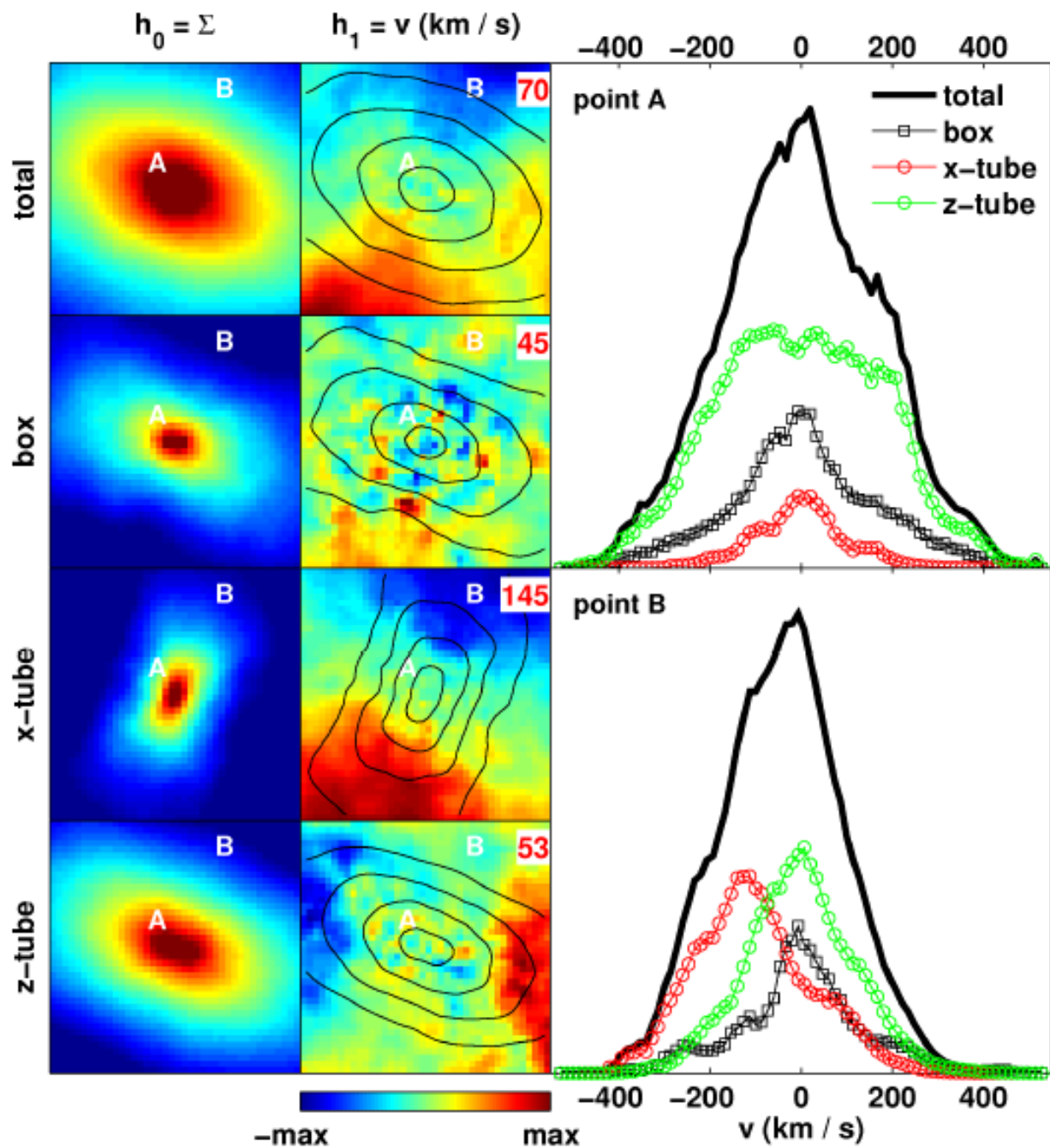


Binney and Spergel 1982;
figure from Hoffman et al. 2010

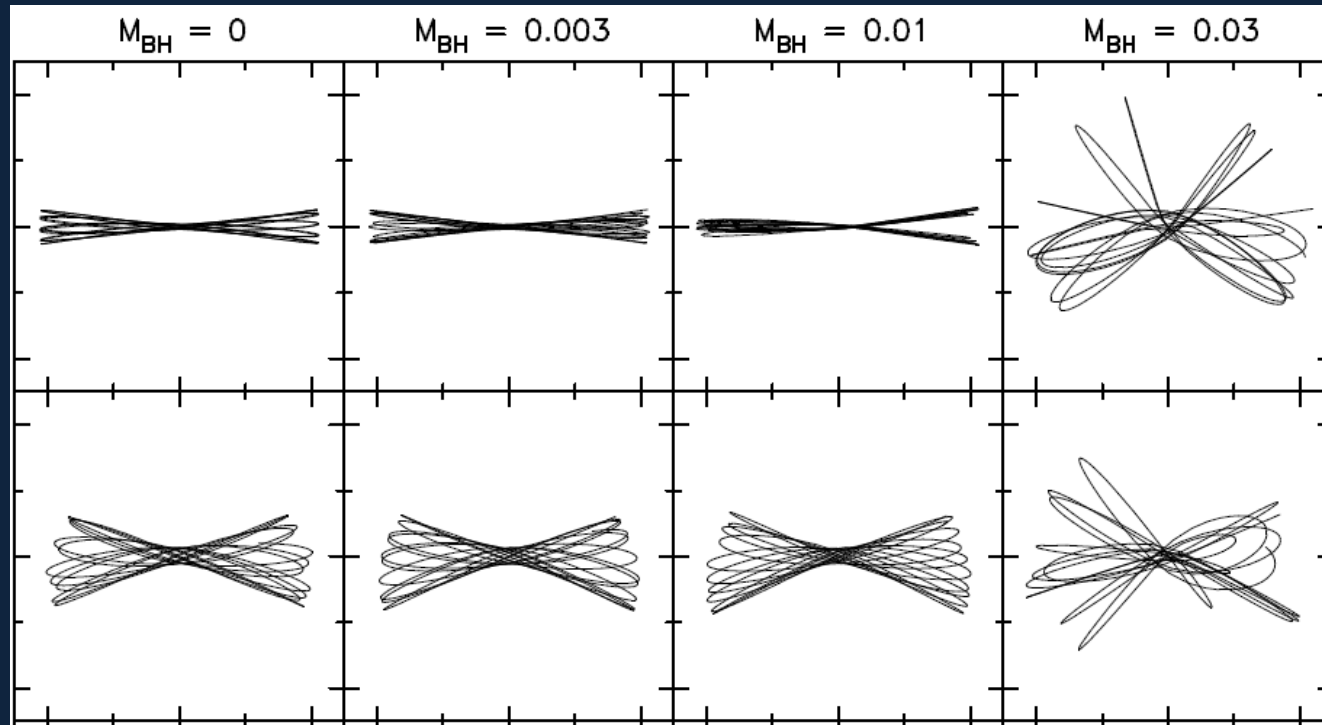


Poon & Merritt 2001

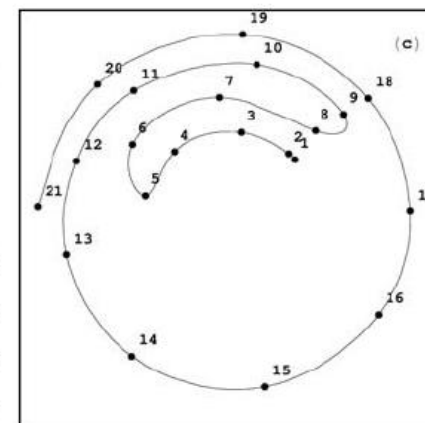
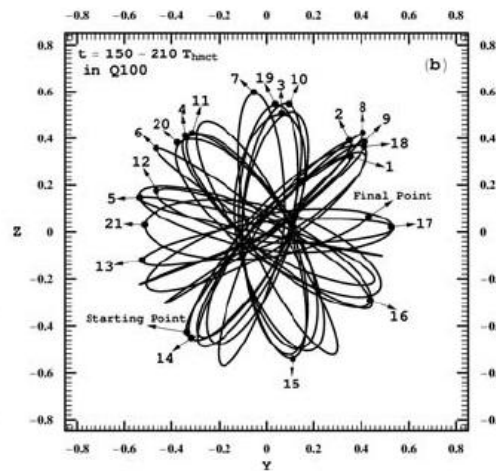
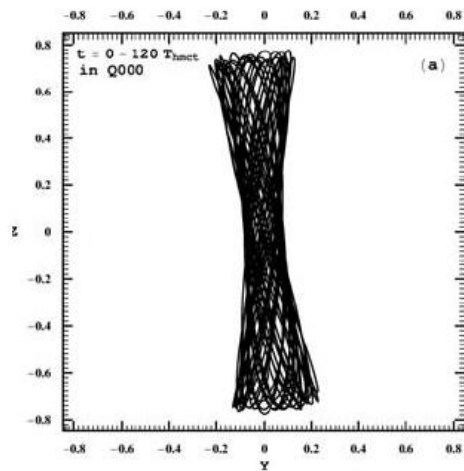
Imprint of orbital structure on kinematics



Origin of the orbital structure - effect of a CMC

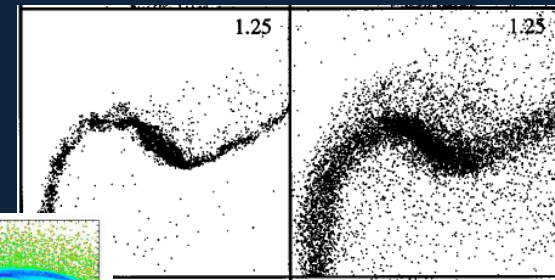
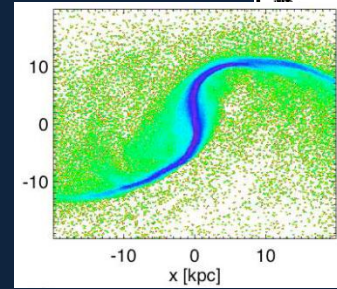
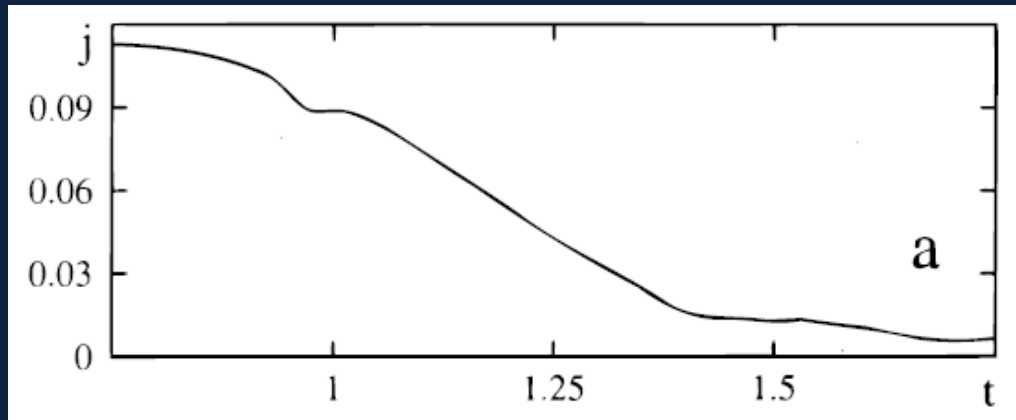


Merritt &
Valluri 1997

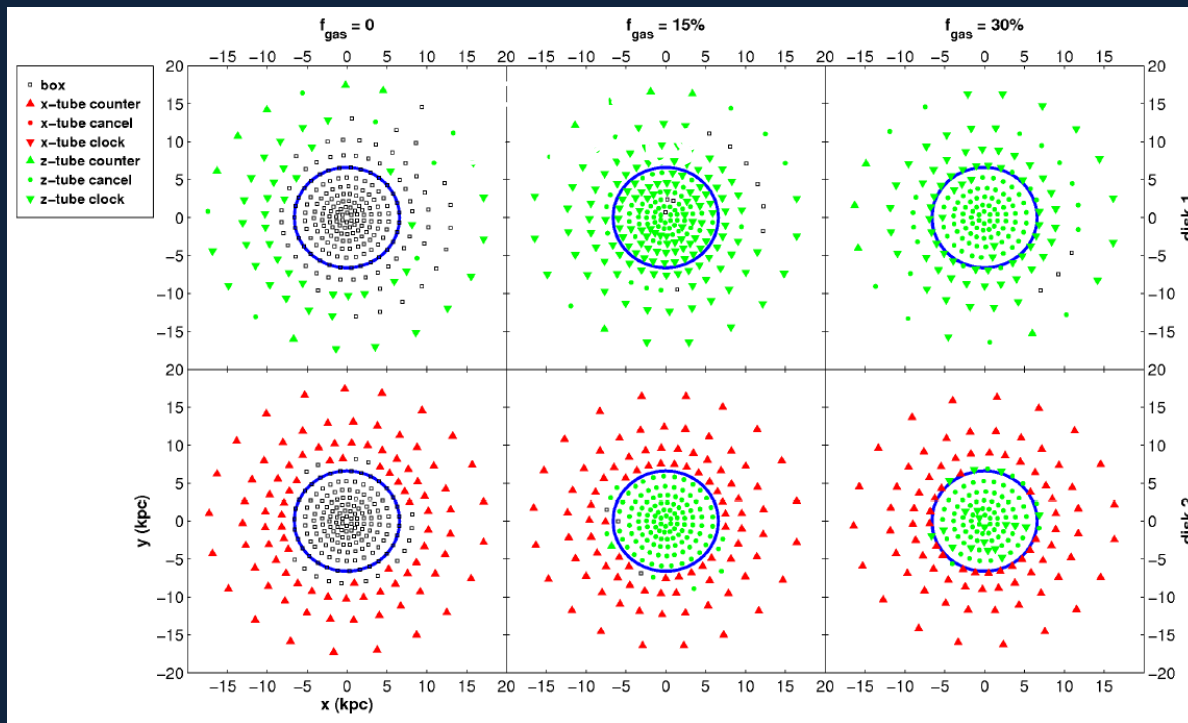


Kalapotharakos
et al. 2004

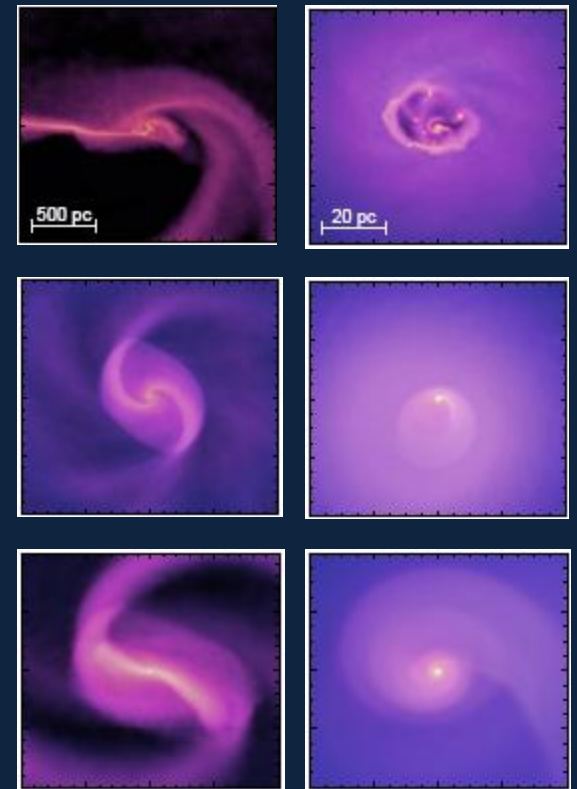
Origin of the orbital structure



Barnes & Hernquist 1996,
Hopkins et al. 2009

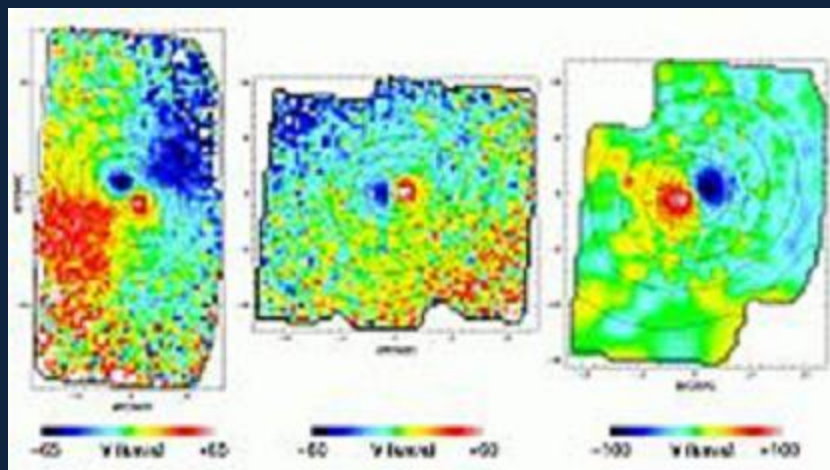
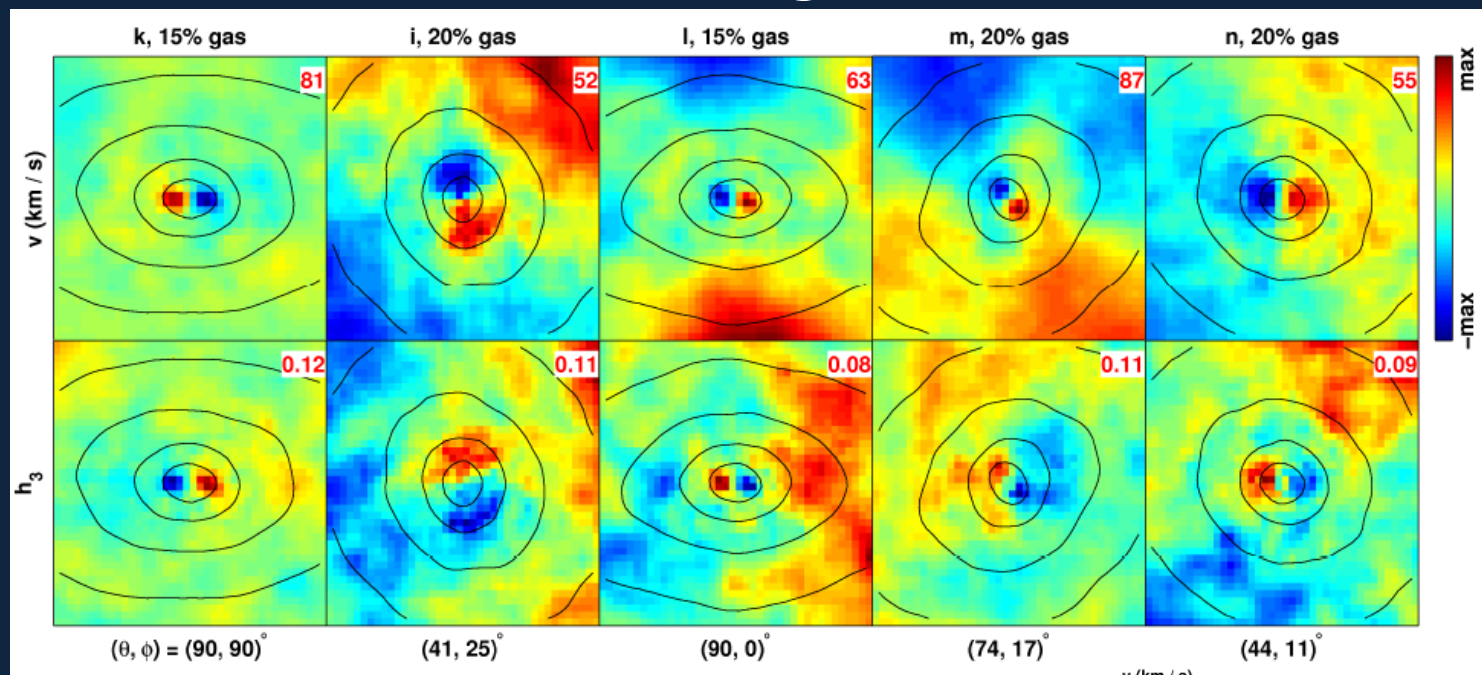


Hoffman et al. 2010

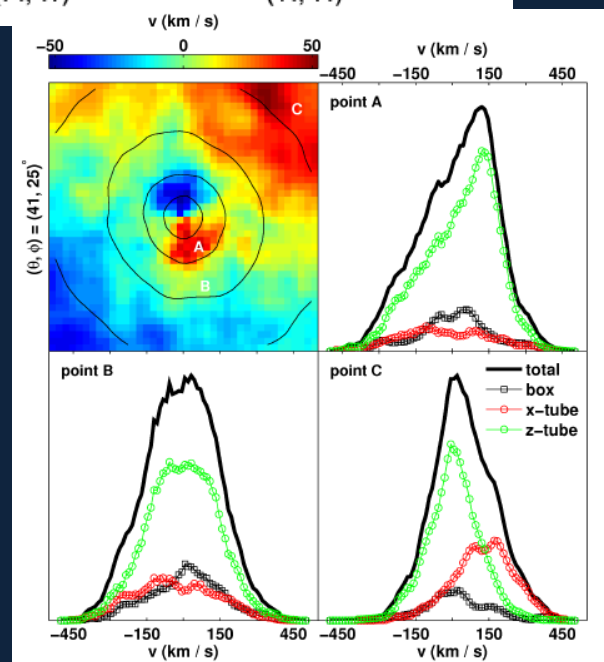


Hopkins et al. 2009

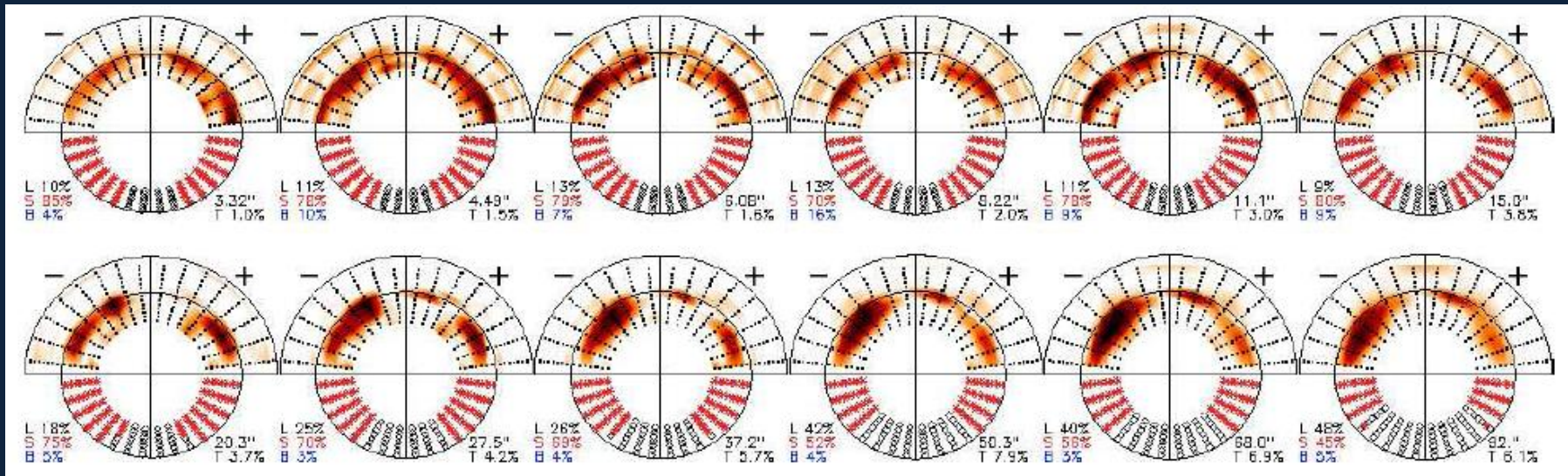
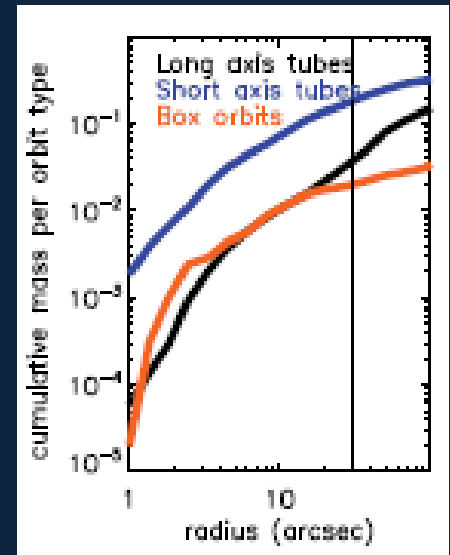
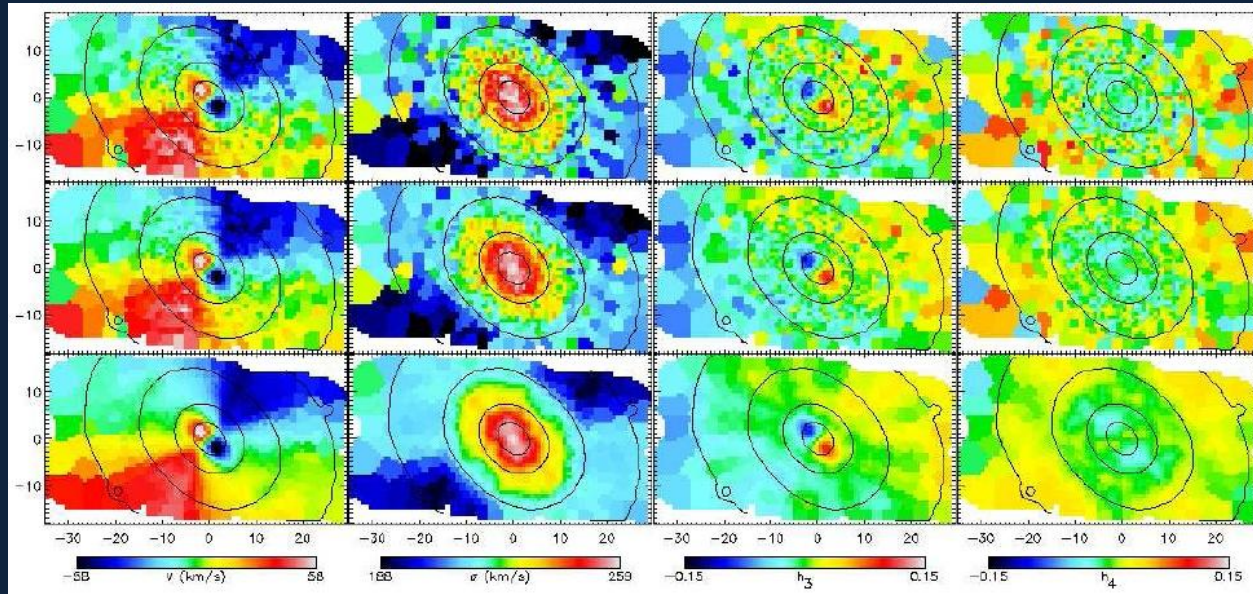
KDCs in 15–20% gas remnants



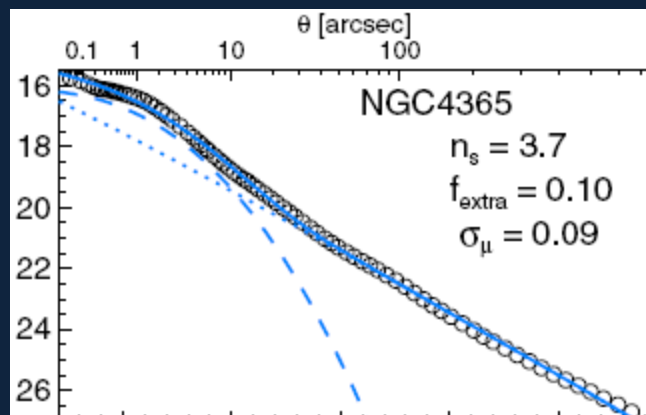
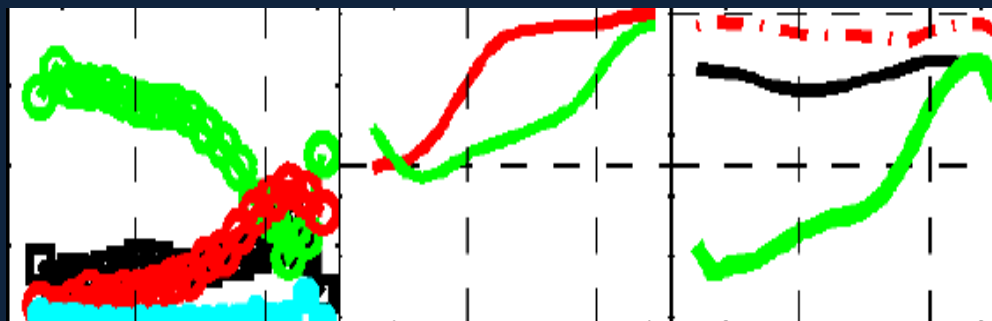
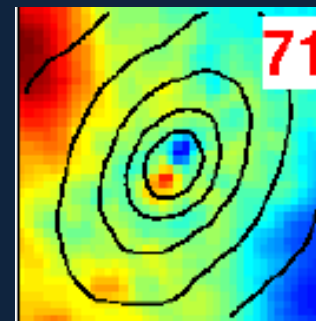
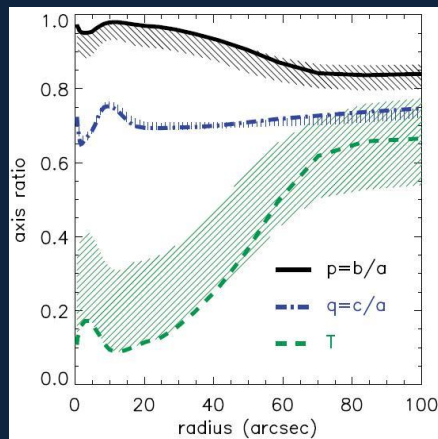
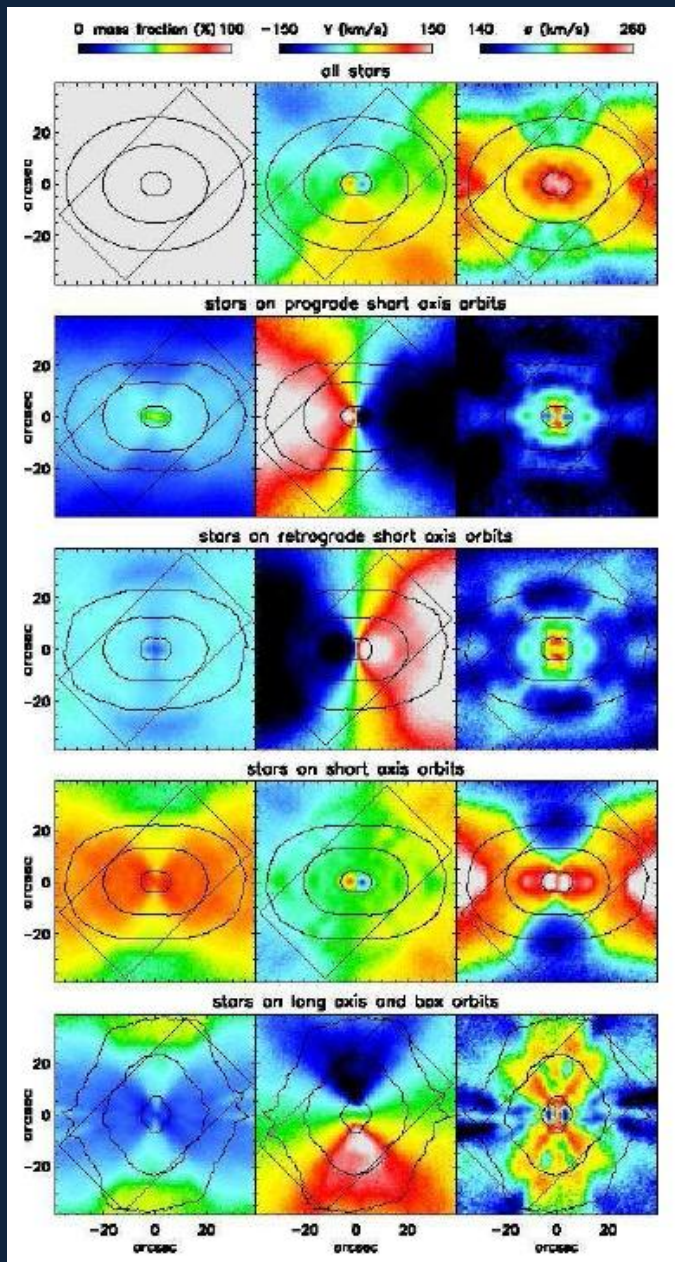
SAURON data



NGC4365 - van den Bosch et al. 2008

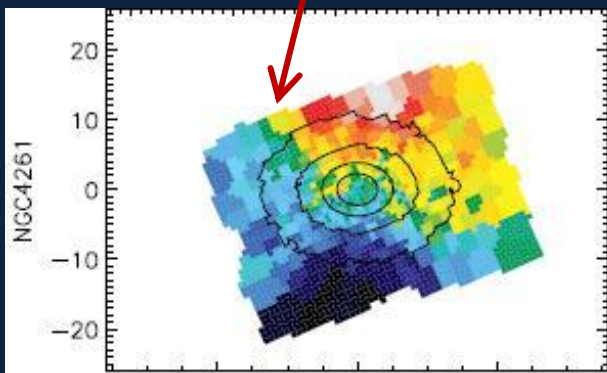
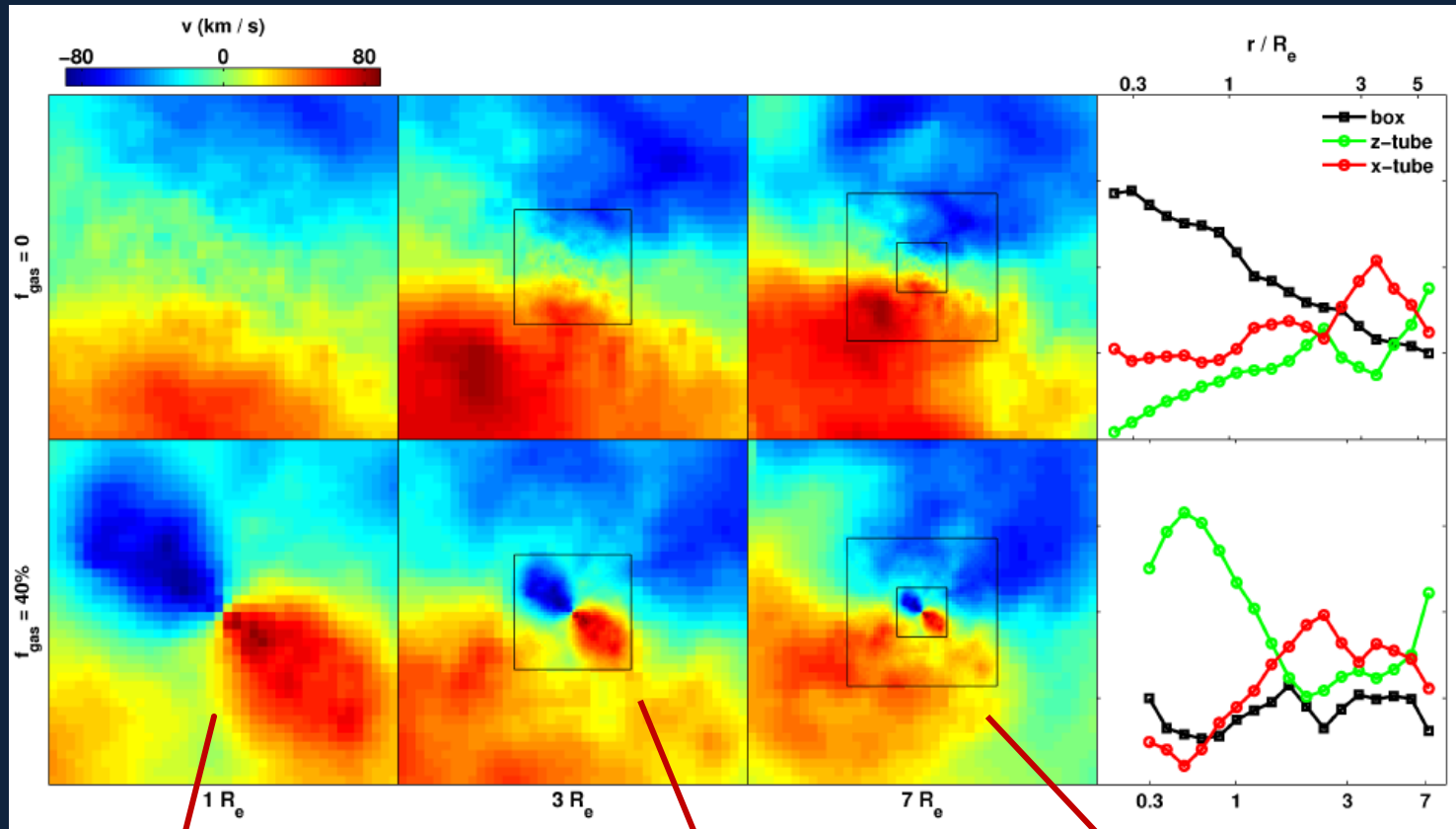


NGC 4365 - comparison of 3D structure with simulations

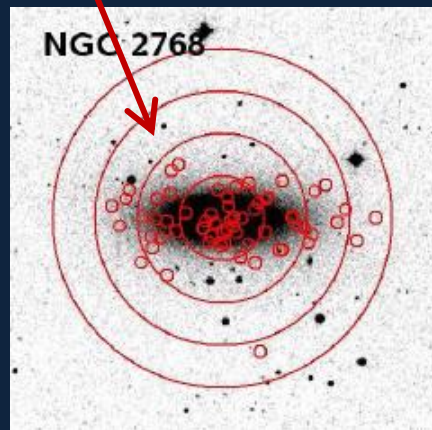


Hopkins et al.
 2009:
 10% "extra
 light"
 component
 in NGC 4365

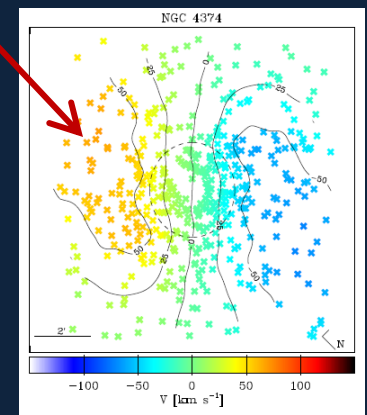
Constraints from outer kinematics



Emsellem et al. 2004

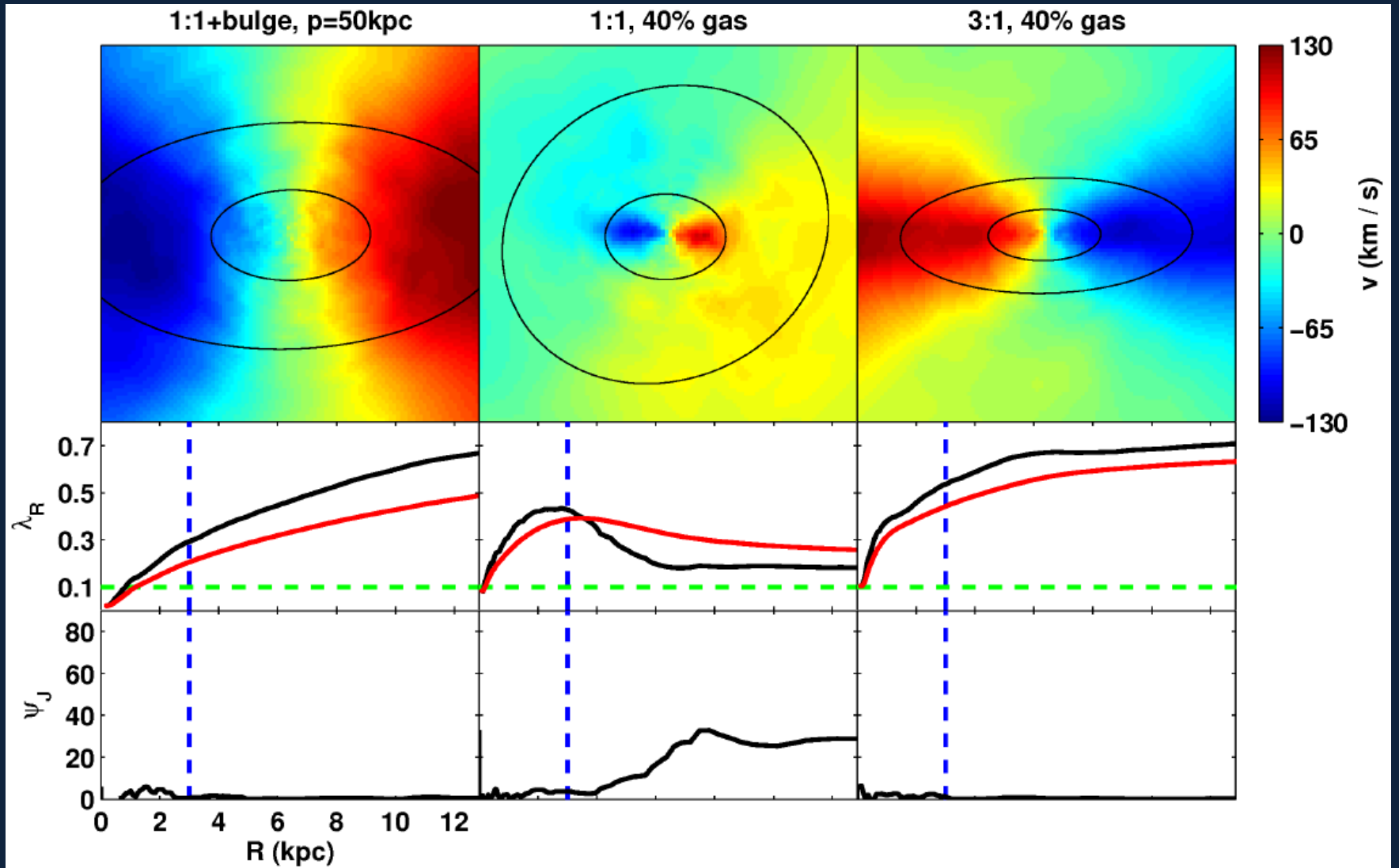


Proctor et al. 2009

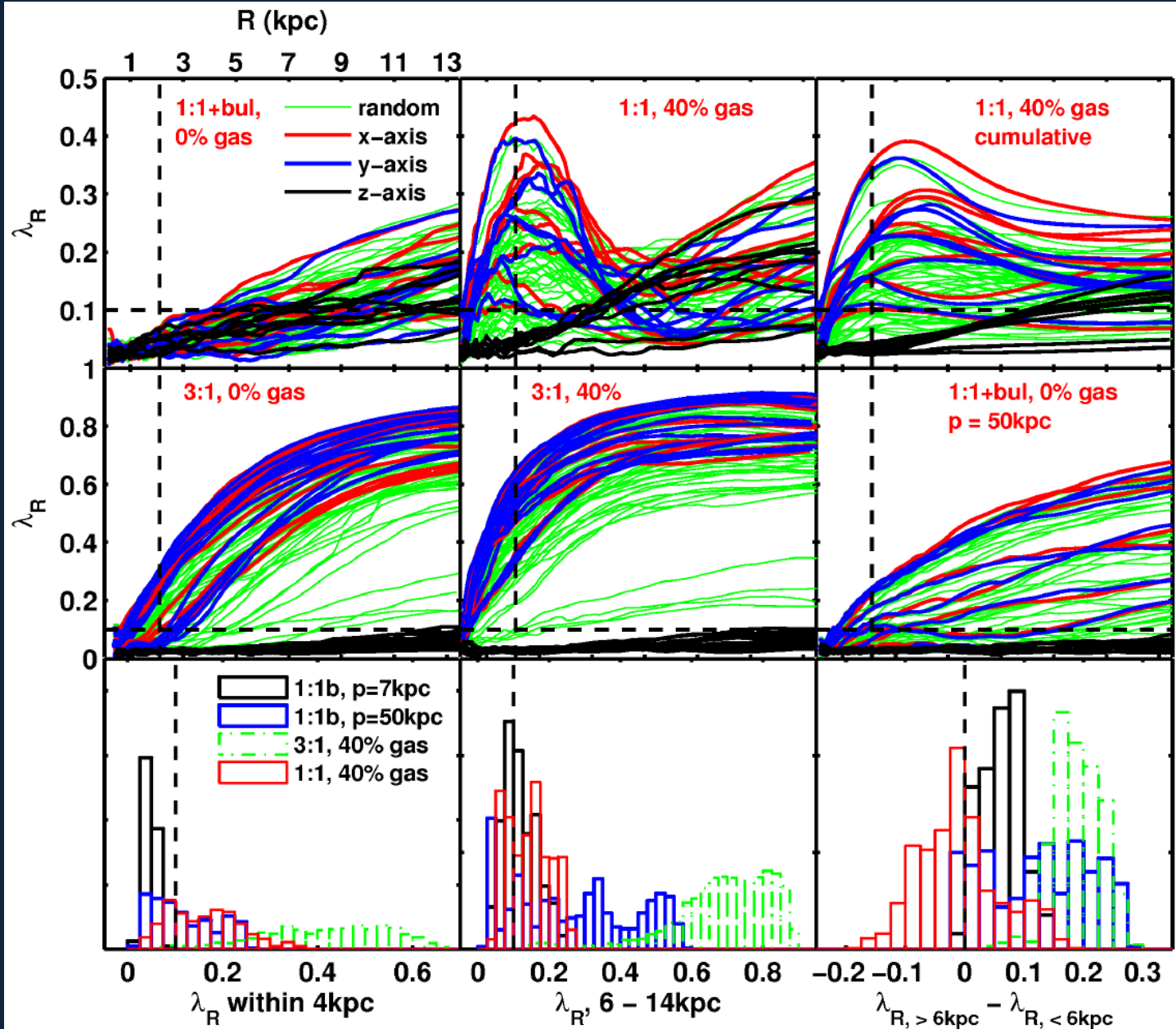


Cocato et al. 2009

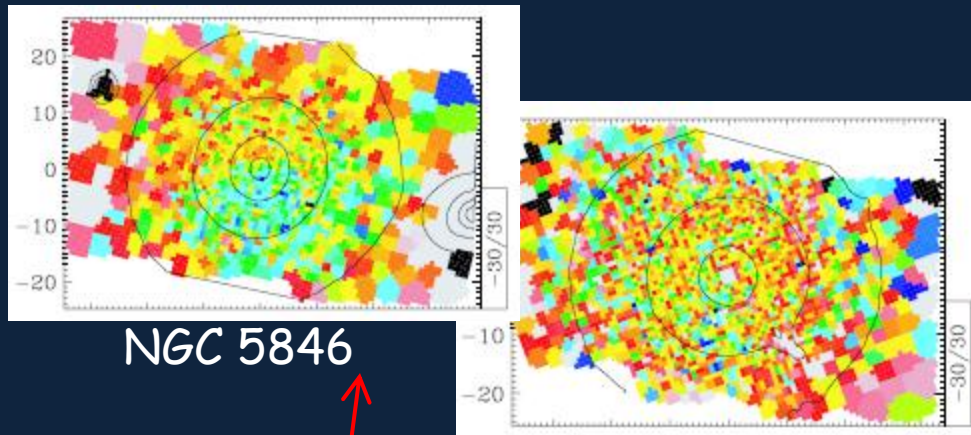
Outer kinematics II



Outer kinematics III

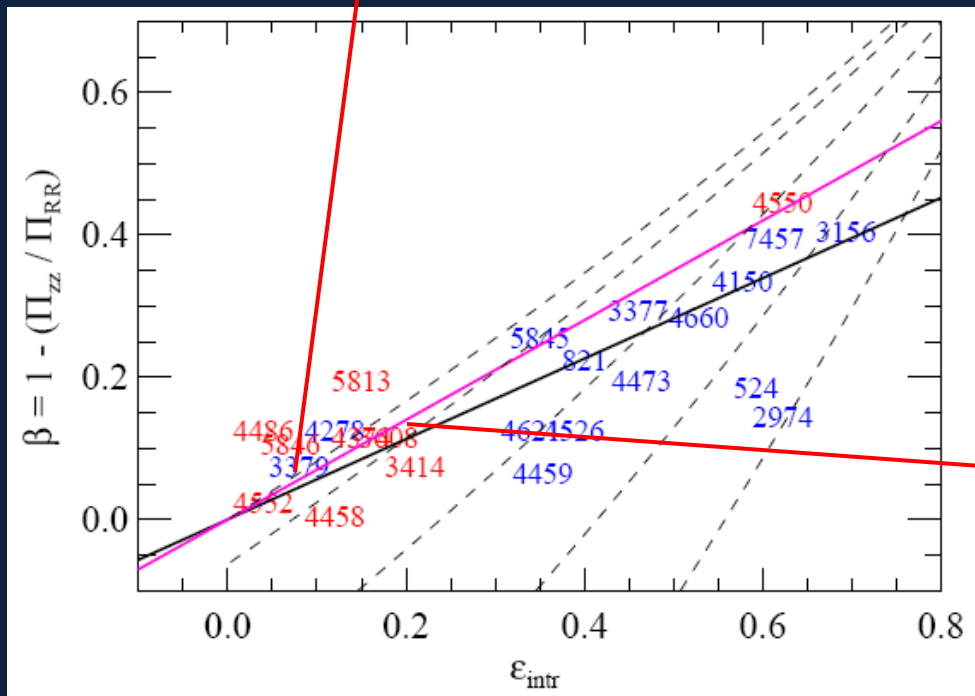


The SAURON challenge



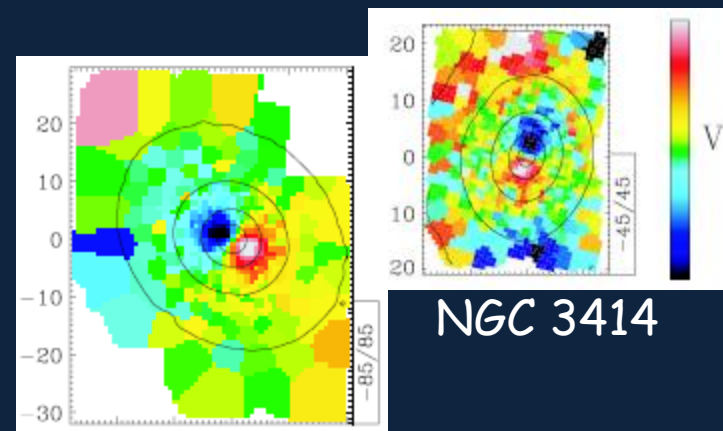
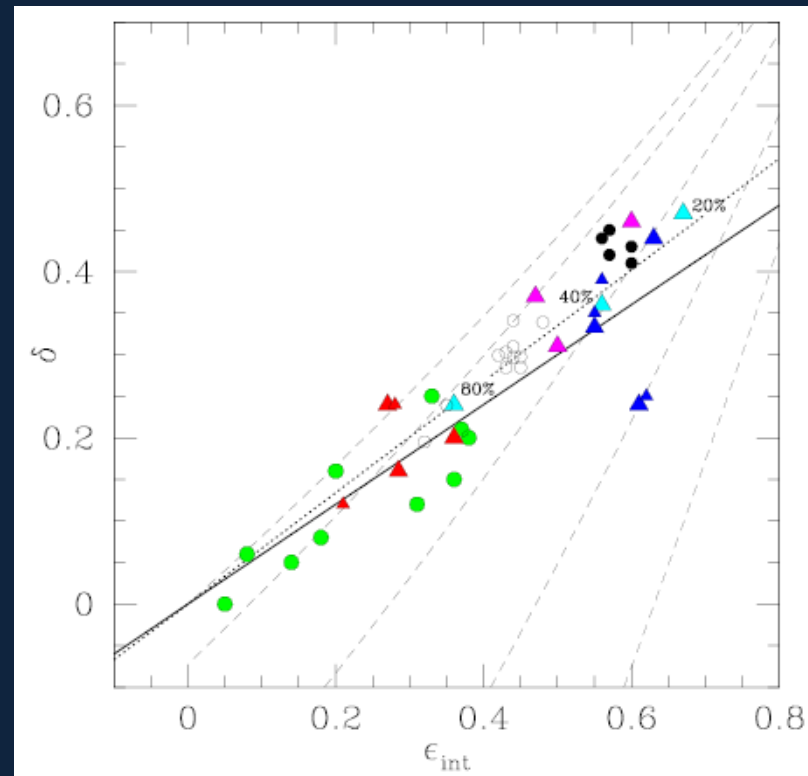
NGC 5846

NGC 4486



Cappellari et al. 2007

Burkert et al. 2008



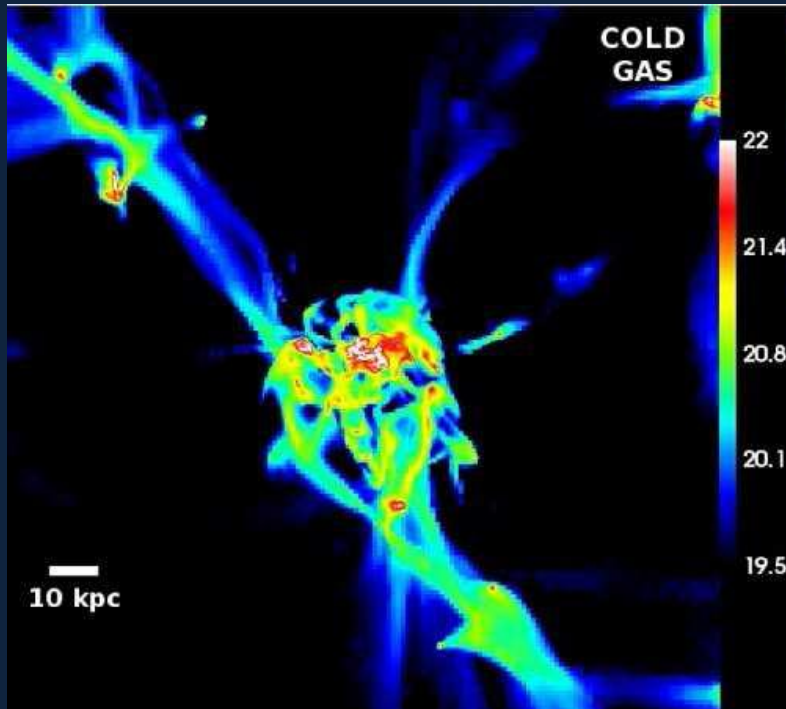
NGC 5813

NGC 3414

More general formation scenarios

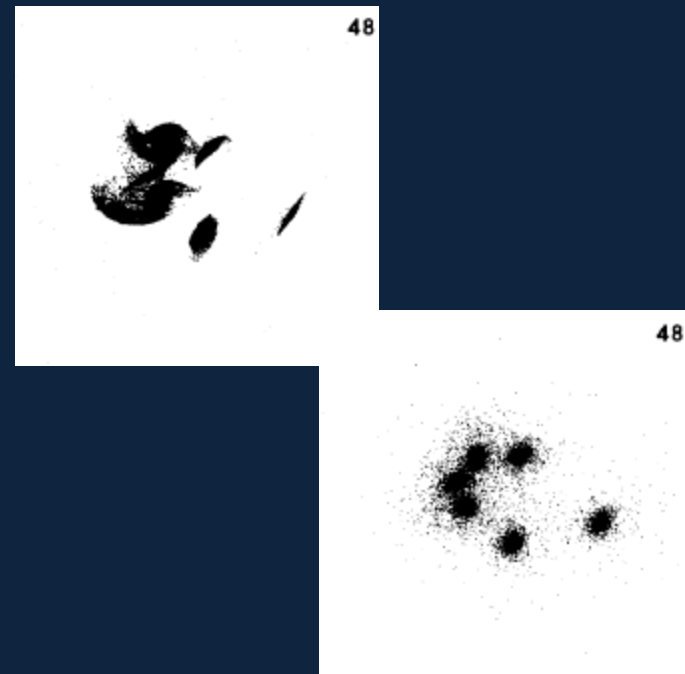
1) Complete exploration of merger phase space

2) Bulge formation through
Clump migration in "wild disks"



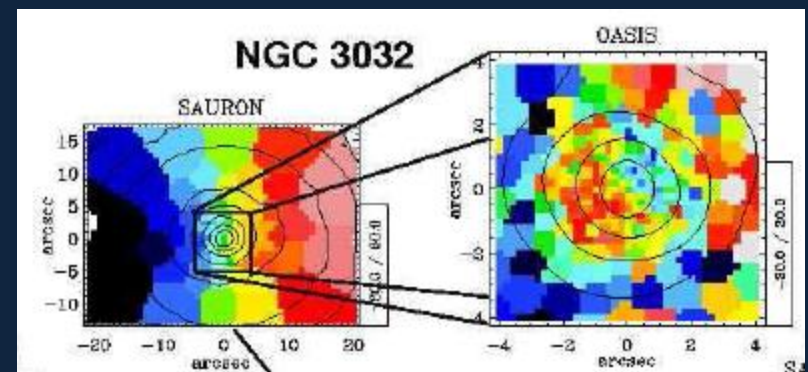
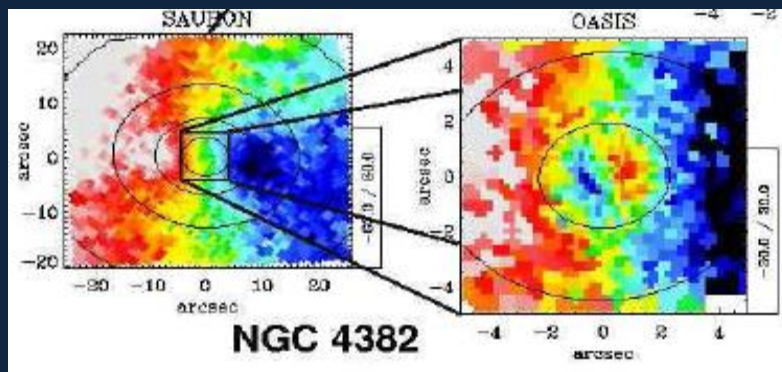
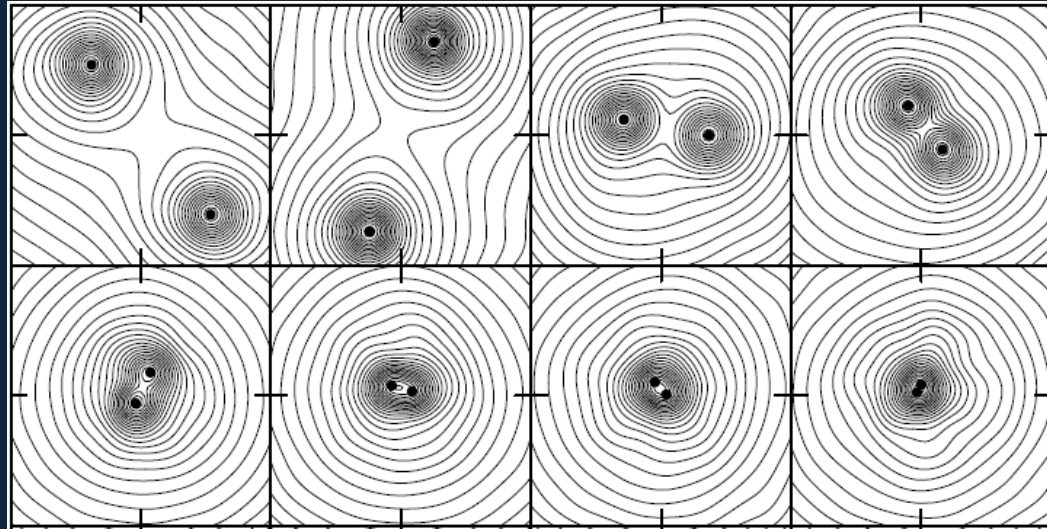
Ceverino et al. 2009

3) Multiple mergers /
cosmologically motivated
sequences



Weil & Hernquist 1996

Binary SMBHs in collisionless galactic nuclei



Requirements: Integrate cuspy system with $N \approx 10^7-8$ for up to 500-1000 t_{dyn} ; suppress relaxation enough to be in empty loss cone regime.

Ideal code: Pure triaxial NBSCF (no kicks) until initial spherical loss cone cleared out, switching to MCSCF after that.

The Self-consistent field (SCF) method

$$\rho(r, \theta, \phi) = \sum_{l=0}^{\infty} \sum_{m=0}^l P_{lm}(\cos \theta) [A_{lm}(r) \cos m\phi + B_{lm}(r) \sin m\phi]$$

(A)

$$\Phi(r, \theta, \phi) = \sum_{l=0}^{\infty} \sum_{m=0}^l P_{lm}(\cos \theta) [C_{lm}(r) \cos m\phi + D_{lm}(r) \sin m\phi]$$

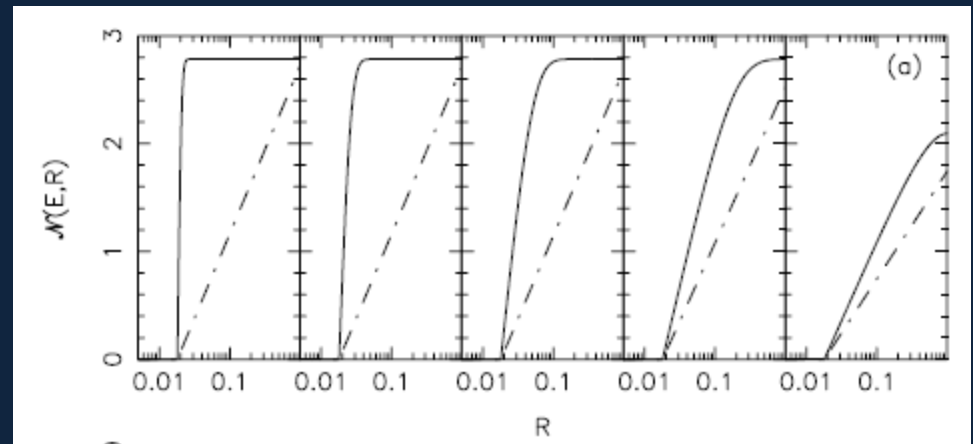
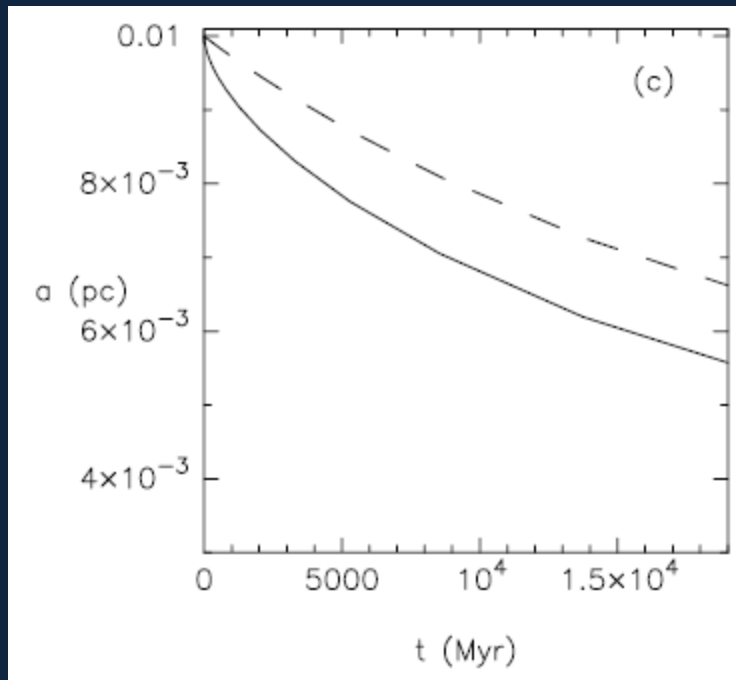
$$\begin{pmatrix} A_{lm}(r) \\ B_{lm}(r) \\ C_{lm}(r) \\ D_{lm}(r) \end{pmatrix} = N_{lm} \sum_{n=0}^{\infty} \tilde{A}_{nl} \begin{pmatrix} \tilde{\rho}_{nl}(r) \\ \tilde{\rho}_{nl}(r) \\ \tilde{\Phi}_{nl}(r) \\ \tilde{\Phi}_{nl}(r) \end{pmatrix} \sum_k m_k \tilde{\Phi}_{nl}(r_k) P_{lm}(\cos \theta_k) \begin{pmatrix} \cos m\phi_k \\ \sin m\phi_k \\ \cos m\phi_k \\ \sin m\phi_k \end{pmatrix}$$

(B)

- 1) $O(N)$
- 2) Completely parallel
- 3) No Cartesian softening

Most astrophysical systems modeled with N-body simulations undergo brief episodes of rapid evolution, long periods of slow evolution ($t_{ev} \gg t_{dyn}$).

Binary S/IMBHs in *collisional* nuclei



Milosavljevic & Merritt 2003
Sharp phase space gradients near the loss cone.

Requirements: Integrate cuspy system with $N \approx 10^7$ for up to 500-1000 t_{dyn} ; suppress relaxation enough to be in empty loss cone regime *and* accurately treat 2-body relaxation.

Ideal code: Spherical, triaxial MCSCF with Henon kicks (note: unlike Spitzer MC, Henon method does not require a Maxwellian distribution!).

Relaxation in N-body simulations

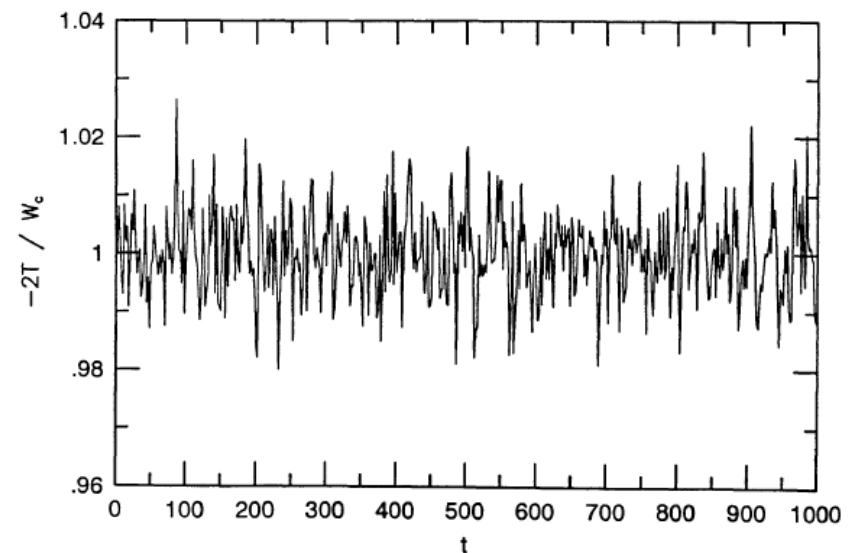
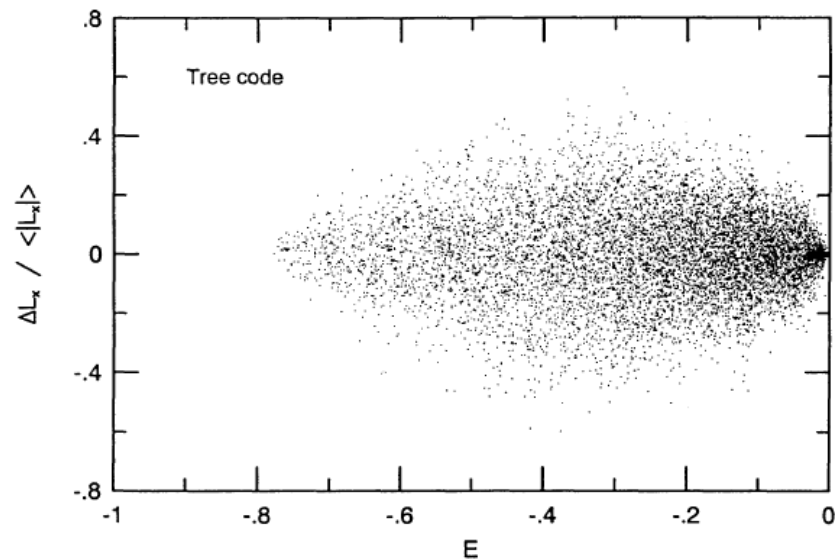
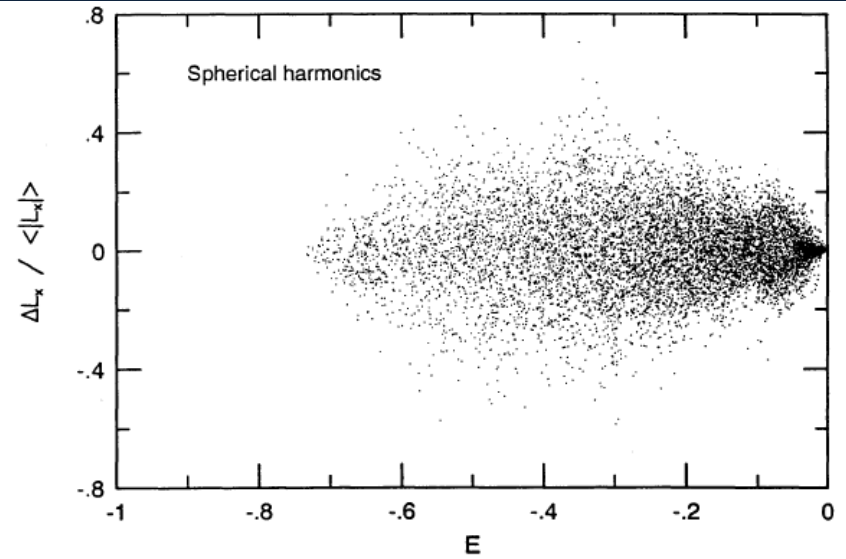
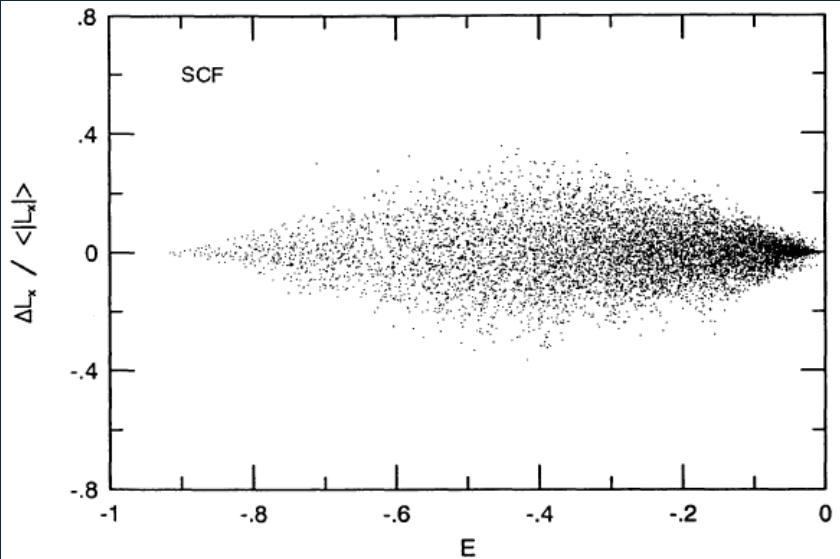


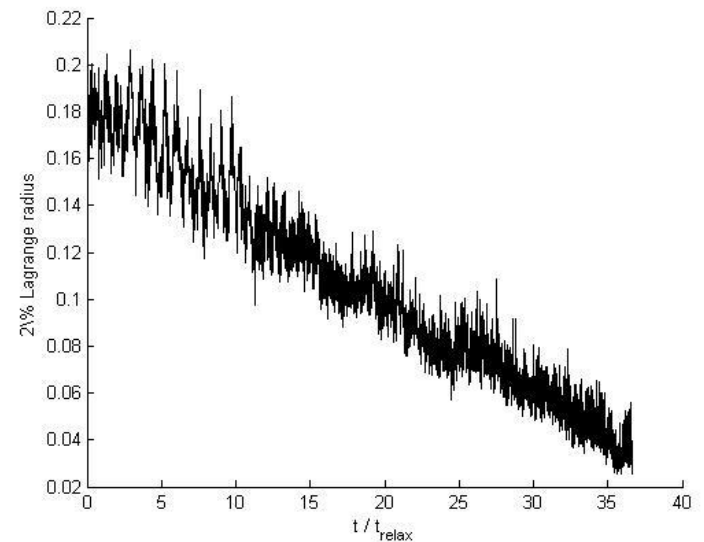
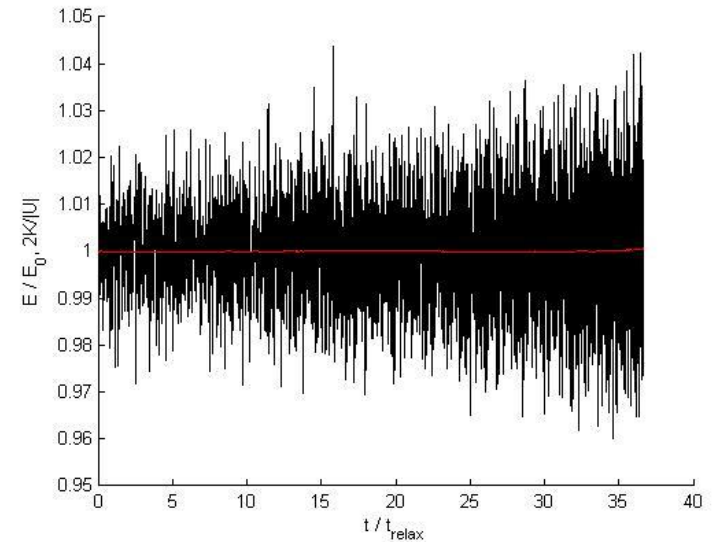
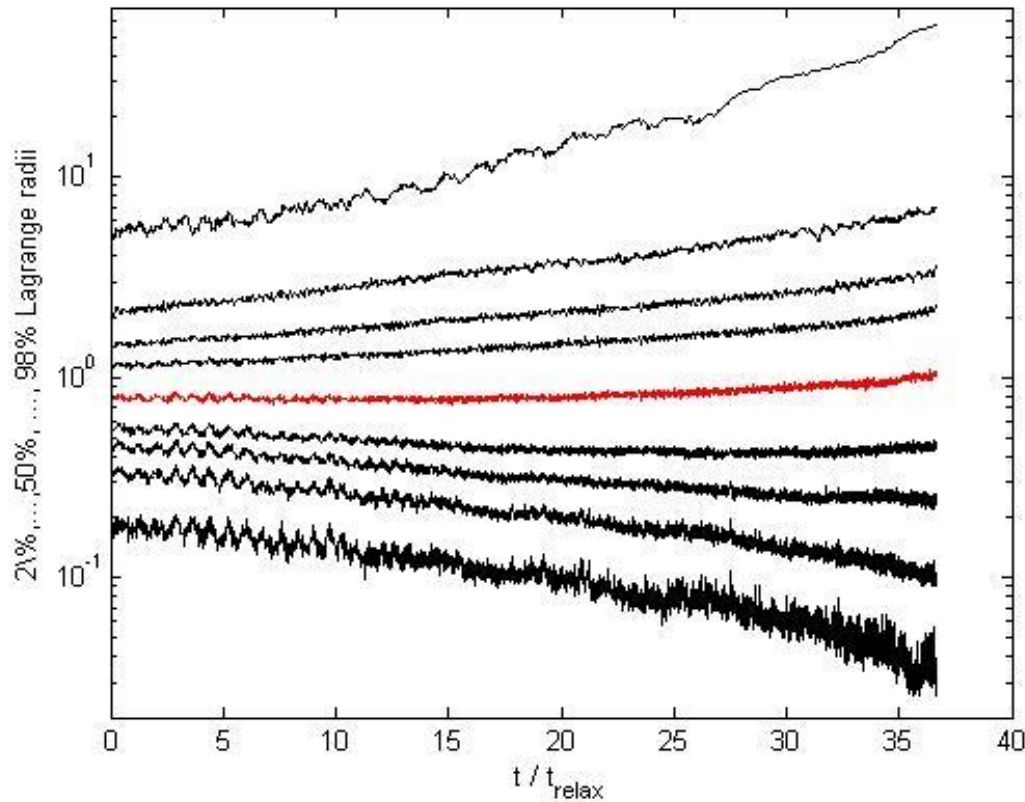
FIG. 12

FIG. 13

Evolution owing to time fluctuations of coefficients

$$N = 2.5 \times 10^3$$

$$N_{\max} = 6, l_{\max} = 4$$



Temporal smoothing

Figure 14 motivates a modification to the SCF method that would be impossible or, at best, very difficult to implement in N -body codes, which we refer to as “temporal smoothing.” That is, one could imagine setting A'_{000} equal to the global time-average of the values plotted in Figure 14, and then integrating particle orbits in this fixed potential. As implied earlier, particle energies would

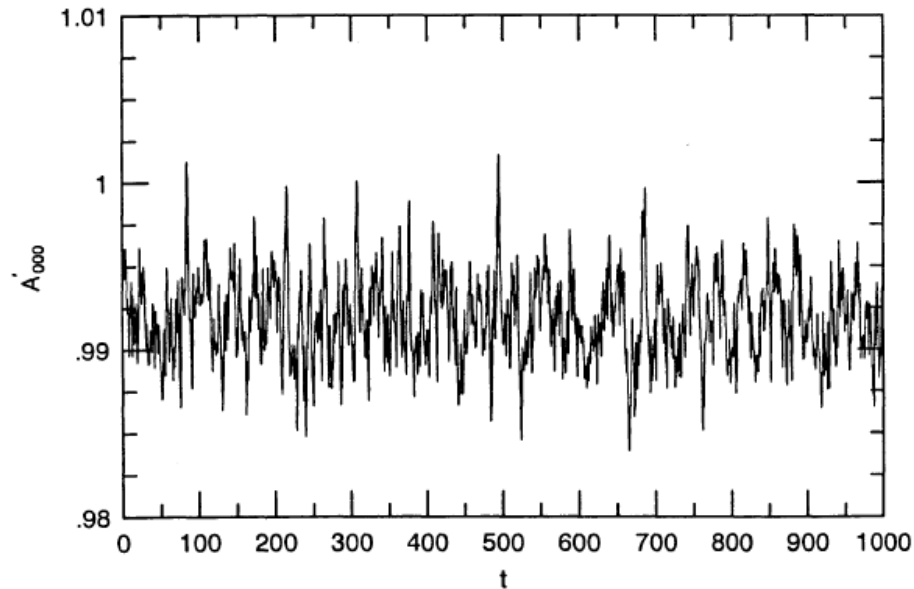
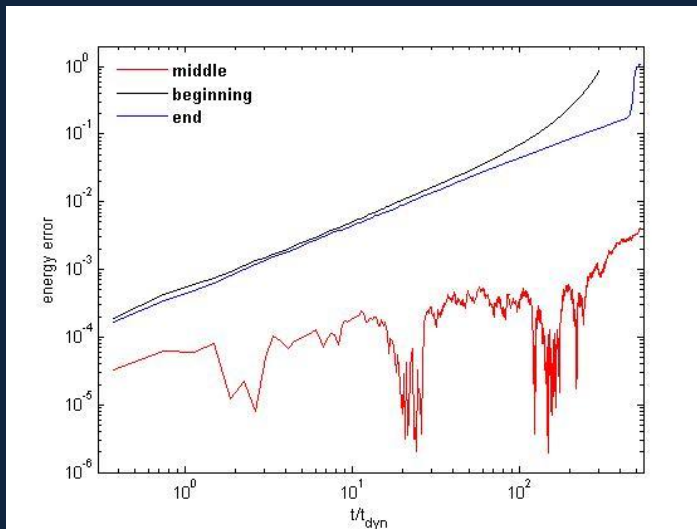
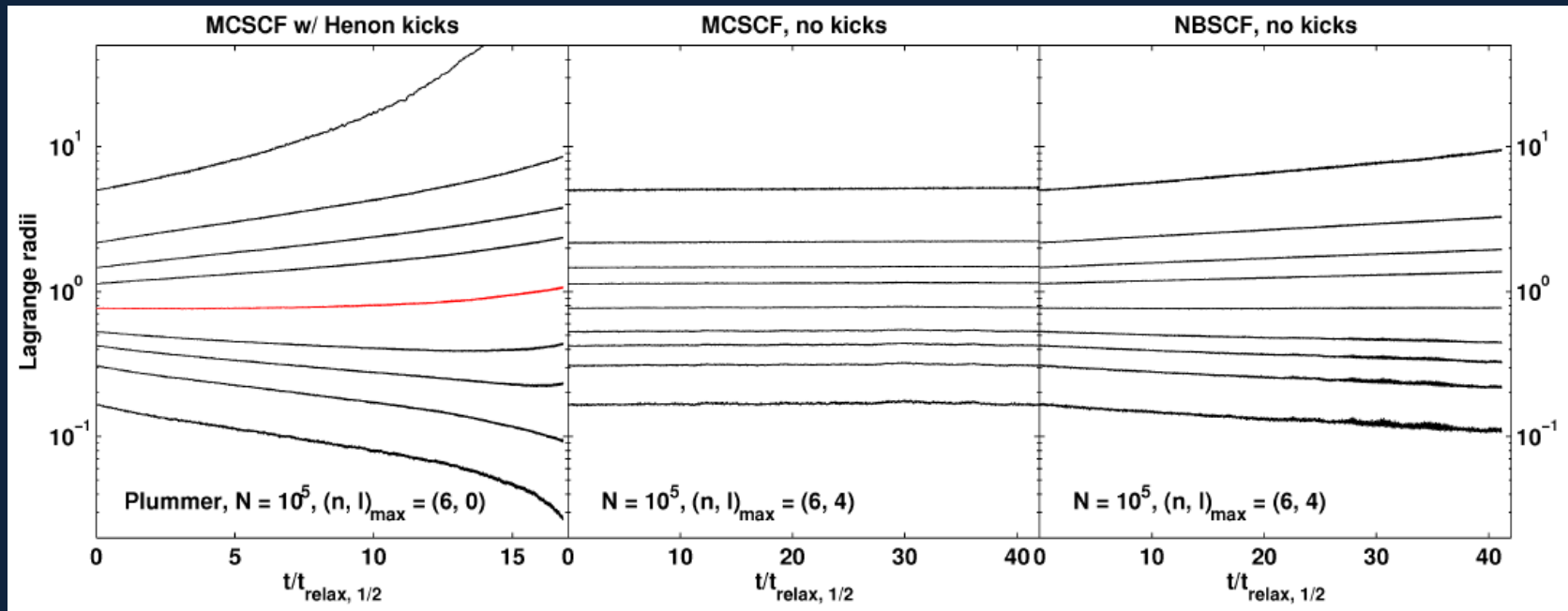


FIG. 14.—Expansion coefficient in a purely radial version of the run shown in Fig. 10, as a function of time. A'_{000} differs from the value of A_{000} given by eq. (2.33) by a constant factor so that A'_{000} would be precisely one in the continuum limit.

“Monte-Carlo” implementation

- 1) Initialize $\{m_k, \mathbf{r}_k, \mathbf{v}_k\}$.
- 2) Compute $\mathbf{a}(\mathbf{r})$ from eq'n (B).
- 3) Integrate each particle's orbit *separately* through $\delta t = \epsilon t_{\text{relax}}$.
- 4) Apply Henon kicks.
- 5) Go back to step 2.

Suppressing relaxation with rare potential updates



Energy conservation is currently an issue for long-term integrations of 10^4 - 10^5 dynamical times.

Summary

- Merger model predicts a characteristic, physically intuitive orbital structure (though not necessarily unique), observable with SAURON + SMEAGOL
- Power of direct comparisons with dynamical models of observed systems: Physical intuition, some features detectable only this way
- Need for extension to more general formation models - cosmological merger trees, smooth/clumpy gas accretion, ...
- SCF / MC technique ideal for studying nuclear kinematics of SMBH binary nuclei; code under development.