# RESONANT STRIPPING AS THE ORIGIN OF DWARF SPHEROIDAL GALAXIES Elena D'Onghia

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Dwarf spheroidals (dSphs) challenge our understanding of galaxy formation and evolution because:

- dSphs are gas poor and have few stars (Mateo 1998)
- found in galaxy clusters and groups (Fergusson & Binggeli 1994)
- the most dark matter dominated galaxies (M/L $\sim$  30–100)
- the ultra-faint dwarf galaxies have L~ 1000 Lsun (Willman et al. 2005, Zucker et al. 2006;Belokurov et al.2009; Walker et al.2008)

# We need a mechanism to separate gas & stars from DM to explain high M/L ratio

## $\diamond$ Previous Theories:

- Gas photoheated during reionization or blown out by feedback, but:
  - few signatures of reionization in dSphs (Gallagher et al. 2003)
  - difficulties explaining the morphology
- Tidal shocking can convert a disk of stars into a spheroid but requires:
  - ram pressure to remove the gas (Mayer et al. 2007)
  - that dwarfs orbit close to Milky Way or Andromeda

## A small dwarf orbiting inside a larger system (D'Onghia et al., 2009, Nature, 460, 605)

The stripping of stars is caused by a gravitational process: "*Resonant Stripping*": stars and gas in the victim are removed by a resonance between the spin frequency of its disk and the angular frequency of its orbit around the perturber.





NEW: "*Resonant Stripping*" alters the M/L ratio in galaxies because stars and gas are removed more efficiently than the dark matter





#### Nearly Prograde Encounter

#### Nearly Retrograde Encounter

Stars

Prograde CCW orbit CCW rotation

Retrograde CCW orbit CW rotation

Courtesy of G. Besla









D'Onghia et al., 2009, Nature, 460, 605

Evolution of kinematic and structural properties



D'Onghia et al., 2009, Nature, 460, 605

### Predictions

*Resonant Stripping* should be visible in situ in groups of dwarfs nearby.

 Many dSphs should be found in groups of dwarfs along with detectable stellar tails and shells.



tadpole galaxy

NGC 2782

# *NOTE: Resonant Stripping depends on a combination of the rotation curve and orbital parameters*

$$\Omega_{dwarf} \approx \Omega_{Pert} \Rightarrow \frac{v}{r} \approx \frac{V}{R} \sqrt{1+e}$$



Fainter Dwarfs may be reproduced200assuming a slowly rising rotation curve



tainter dwarfs



✓ The tails and bridges of stars in major mergers are caused by a tidal resonance (Toomre & Toomre 1972)

Simulations have shown that 10% of stars are removed
during major mergers

## "Quasi-Resonance theory of tidal interactions"



$$\Delta v(x) = -GM_{pert} \int \left[ \frac{x}{|X|^3} - \frac{X(x \cdot X)}{|X|^5} \right] dt$$



$$\Delta v_x = -2 \frac{GM_{pert}}{b^2 V} r \cos \phi_0 (-\alpha^2 K_0 - \alpha (1 \pm \alpha) K_1$$
$$\Delta v_y = -2 \frac{GM_{pert}}{b^2 V} r \sin \phi_0 (\alpha^2 K_0 \pm \alpha^2 K_1)$$

 $\alpha = \frac{|\Omega|b}{V}$ 

angular frequency

D'Onghia et al., 2010, in prep

## Limits of "Tidal Quasi-Resonance approximation"

✓ For angular frequency=0 Tidal near-resonance  $\rightarrow$  Impulse Approximation

 $\alpha = 0$ 

$$\Delta v_x = -2 \frac{GM_{pert}}{b^2 V} r \cos \phi_0 (-\alpha^2 K_0 - \alpha (1 \pm \alpha) K_1)$$
$$\Delta v_y = -2 \frac{GM_{pert}}{b^2 V} r \sin \phi_0 (\alpha^2 K_0 \pm \alpha^2 K_1)$$

$$\Delta v_x = -2 \frac{GM_{pert}}{b^2 V} r \cos \phi_0$$
$$\Delta v_y = -2 \frac{GM_{pert}}{b^2 V} r \sin \phi_0$$

 $\alpha = \infty$ ✓ For Infinite angular frequency Tidal near-resonance  $\rightarrow$  Adiabatic Invariant

$$\Delta v_x = -2 \frac{GM_{pert}}{b^2 V} r \cos \phi_0 (-\alpha^2 K_0 - \alpha (1 \pm \alpha) K_1)$$
  
$$\Delta v_y = -2 \frac{GM_{pert}}{b^2 V} r \sin \phi_0 (\alpha^2 K_0 \pm \alpha^2 K_1)$$

$$\Delta v_x = -2 \frac{GM_{pert}}{b^2 V} x \alpha e^{-\alpha} \sqrt{\frac{\pi}{2\alpha}} (1 - \alpha \pm \alpha) \to 0$$
$$\Delta v_y = 2 \frac{GM_{pert}}{b^2 V} y \alpha e^{-\alpha} \sqrt{\frac{\pi}{2\alpha}} (1 - \alpha) \to 0$$

## "Tidal Quasi-Resonance approximation"



Perturber

on parabolic orbit

$$\alpha = \frac{|\Omega|b}{V}$$

$$\Delta v_x = -4 \frac{GM_{pert}}{b^2 V} r \cos \phi_0 (2I_{30} - 2(1 \pm 2\alpha)I_{20} \pm 3\alpha I_{10} - 4\alpha^2 I_0)$$
  
$$\Delta v_y = 4 \frac{GM_{pert}}{b^2 V} r \sin \phi_0 (2I_{30} - (1 \pm 4\alpha)I_{20} \pm 3\alpha I_{10} - 4\alpha^2 I_{00})$$

## The Tidal Quasi-Resonance approximation is more efficient than the impulse approximation



angular frequency

#### D'Onghia et al., 2010, in prep

## Efficiency of Resonant Stripping



The efficiency depends on: mass ratio (m/M) and impact parameter b Comparison between the quasi-resonance theory and simulations

Prograde Orbit



D'Onghia et al. 2010, in prep

## The energy and angular momentum distributions match !





## Retrograde Orbit







#### The Tidal Resonance Approximation







#### Numerical Simulation

## Non Coplanar Orbit



The Tidal Resonance Approximation



Numerical Simulation

## Parabolic Orbit



# The Tidal Resonance Approximation

Numerical Simulation

# The Magellanic Stream as the product of an LMC + SMC tidal encounter



## Conclusion

 Interactions between dwarf galaxies in small groups excite a resonant response: "*Resonant Stripping*" that rapidly transforms disks into dSphs.

 Resonant stripping is a gravitational process that removes gas & stars in a disk but affects less DM and can be described by tidal Quasi-Resonance Theory

✓ TEST: -- rotational velocity of the stars in data

-- subhalo angular momentum in cosmological simulations and if they are preferentially retrograde