

# Dense Matter and Neutrinos

J. Carlson - LANL

- Neutron Stars and QCD phase diagram
- Nuclear Interactions
- Quantum Monte Carlo
- Low-Density Equation of State
- High-Density Equation of State
- Neutron Star Matter  
(protons, hyperons, etc.)
- Mass/Radius relations and observations
- Neutrinos - neutron star cooling
- Future

Collaborators:

S. Gandolfi (LANL)  
A. Gezerlis (Guelph)  
D. Lonardoni (ANL)  
A. Lovato (ANL)  
S. Reddy (UW/INT)

# Neutron Stars

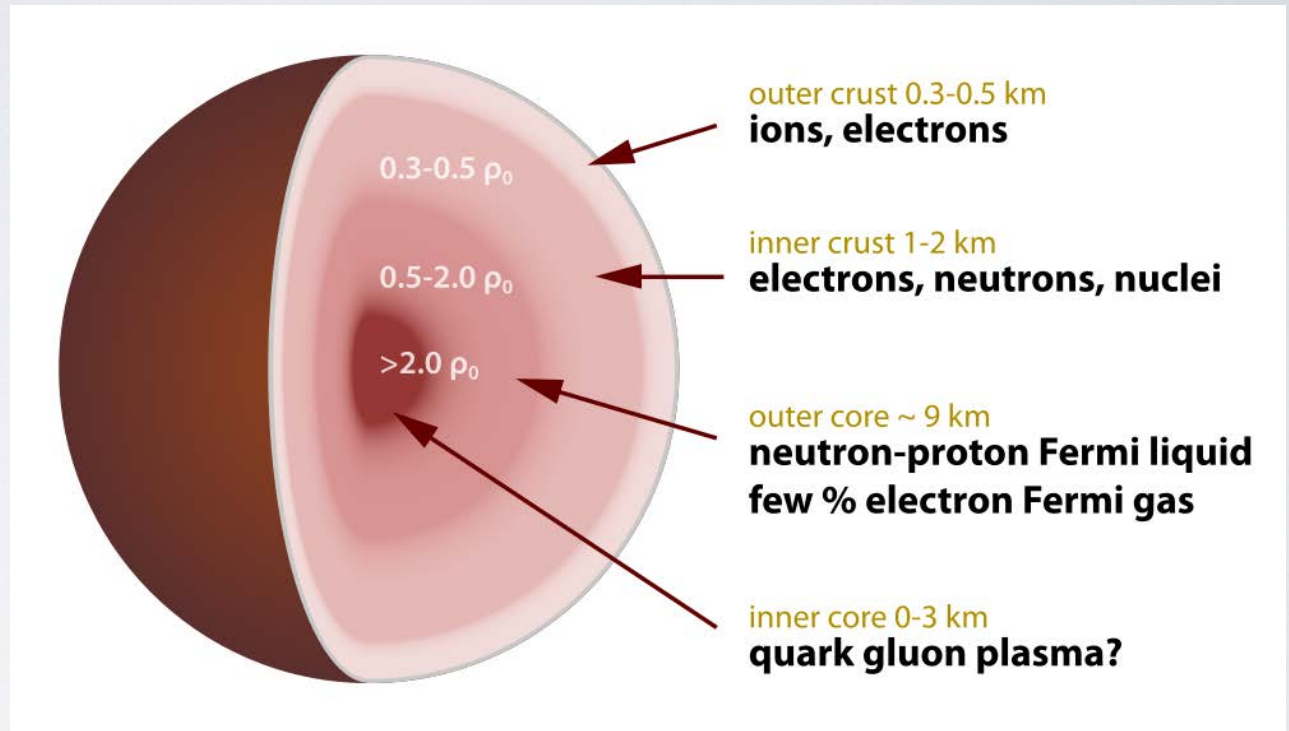
1-2 Solar Masses  
~12 km radius

Outer Crust  
nuclei + electrons

Inner Crust  
nuclei + neutrons +  
electrons

Core  
neutrons+protons+electrons+...

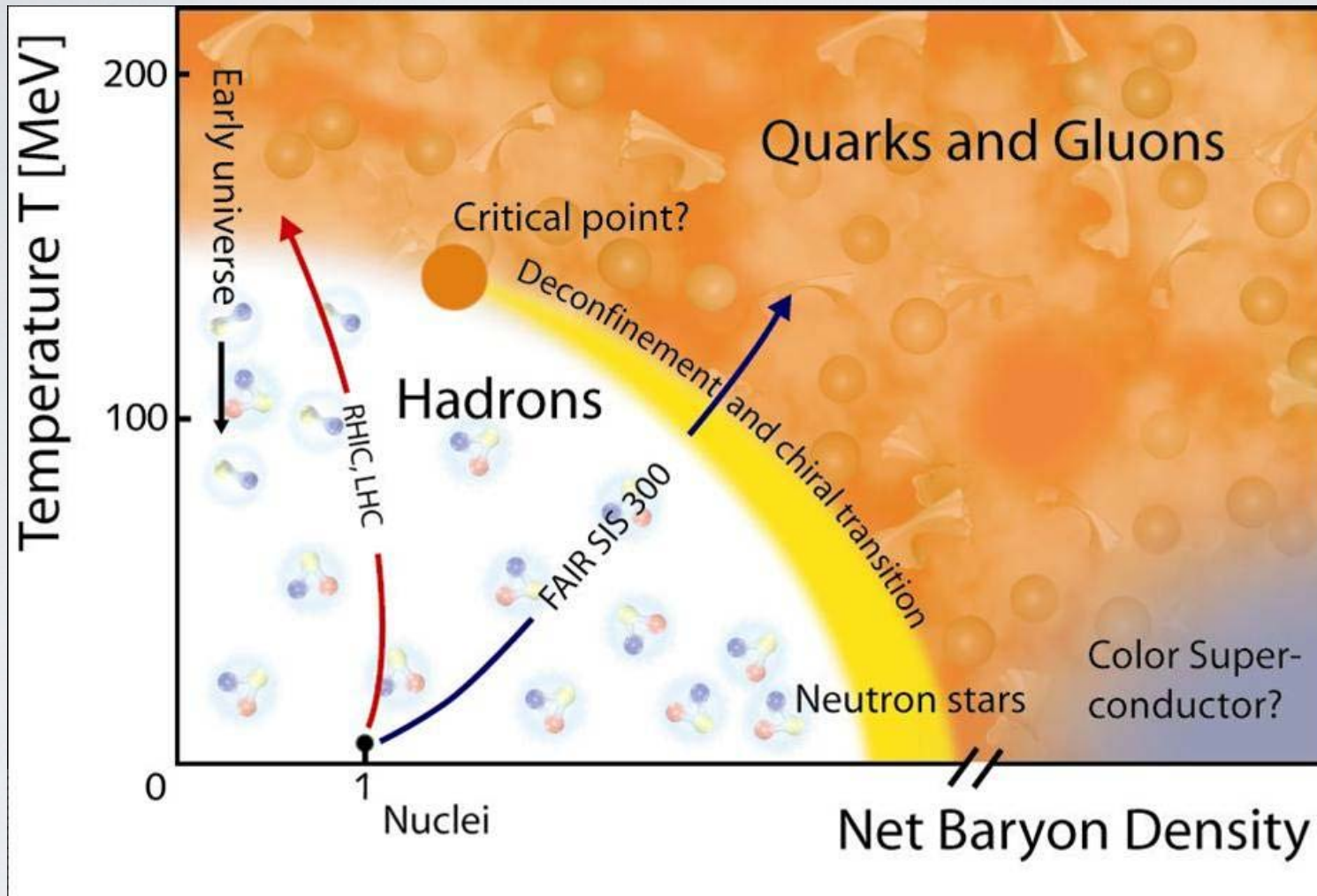
Predicted by Baade and Zwicky 1 year after discovery of the neutron



We will concentrate  
on the core: bulk of the star  
dominates the M/R curve  
important for neutrino cooling

charge neutrality + small electron mass  $\rightarrow$  ~10% electrons, protons

# QCD phase diagram (minimal)



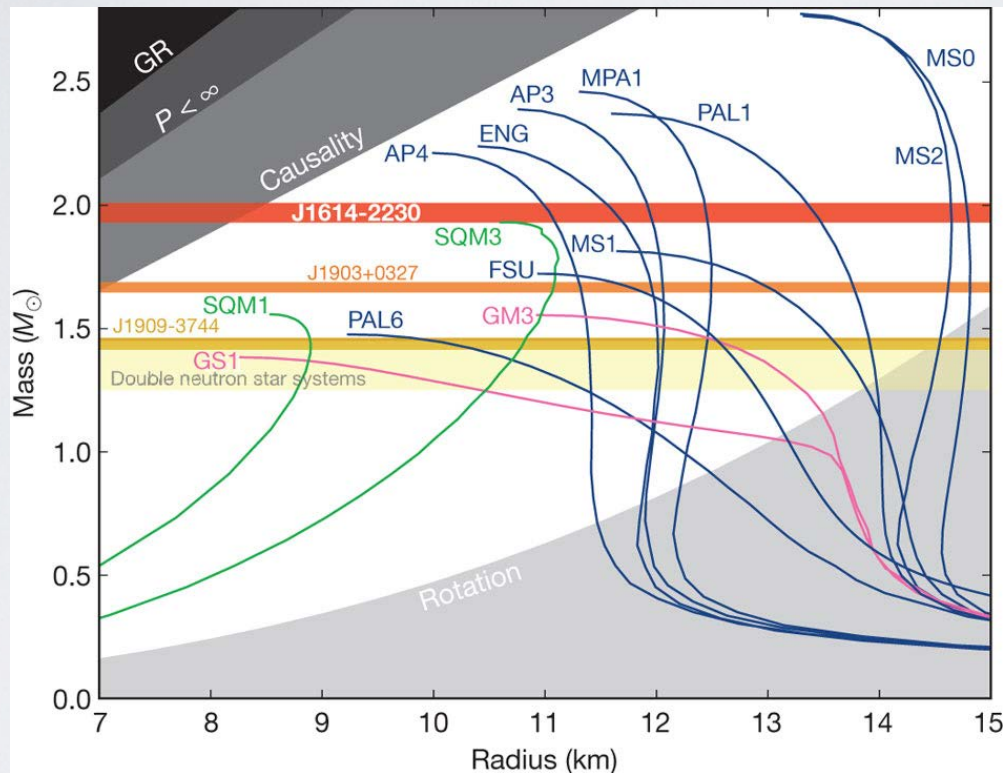
from FAIR, new facility in Darmstadt

high density and cold very difficult to reach in experiments

Color superconductor at very high density; important for neutron stars?

# Neutron Star Mass/Radius Relations

For many years only  $\sim 1.4$ - $1.5$  solar mass neutron stars observed  
Recently several observed with  $\sim 2$  solar masses!



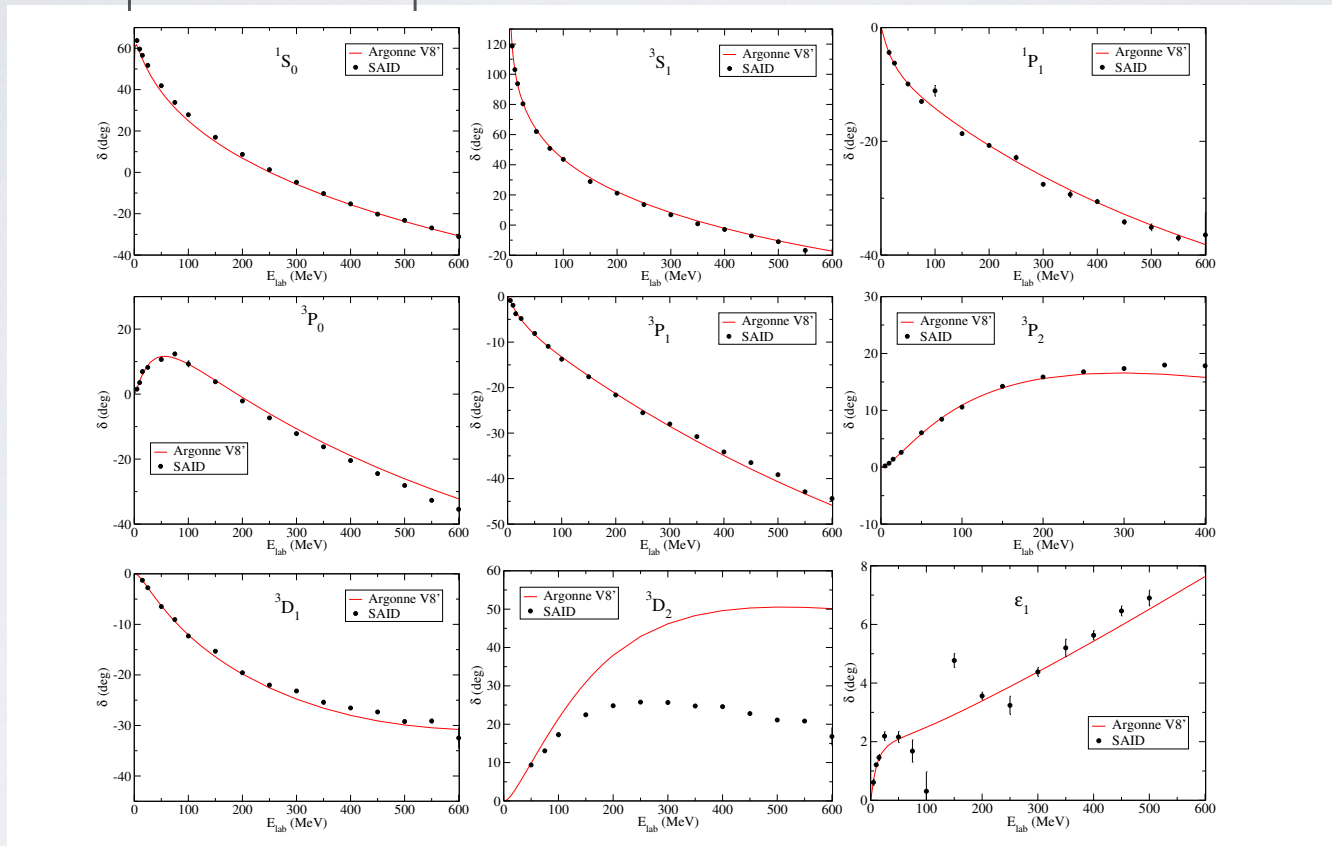
see Demorest, et al  
Nature 467: 2081 (2010)

Transitions to superconducting quark matter  
Wide range of predictions for mass/radius relationship

# Nuclear Interactions

Up to  $\sim 2-3 \times$  nuclear densities, matter can be described as a system of interacting nucleons

phase shifts for NN scattering - simple model (AV8') compared to experiment



At  $r=2 \rho_0 : k_F \sim 2 \text{ fm}^{-1}$

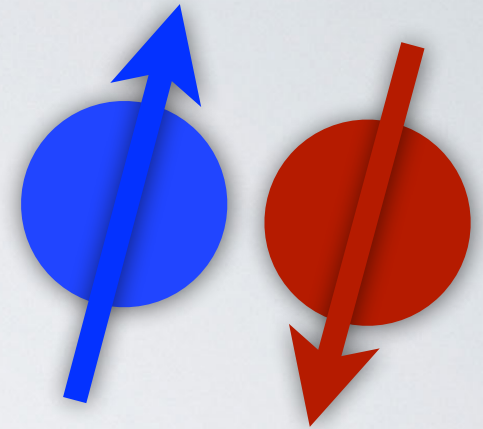
implies 2 nucleons at Fermi surface have  $E_{CM} = 160 \text{ MeV}; E_{lab} \sim 320 \text{ MeV}$

# Nuclear Interactions

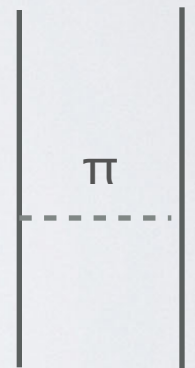
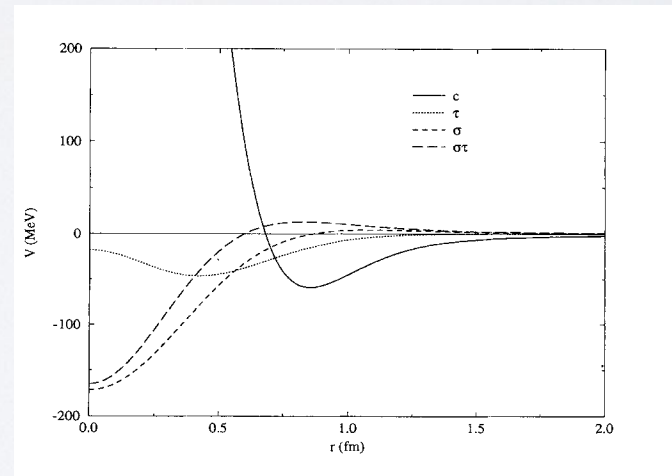
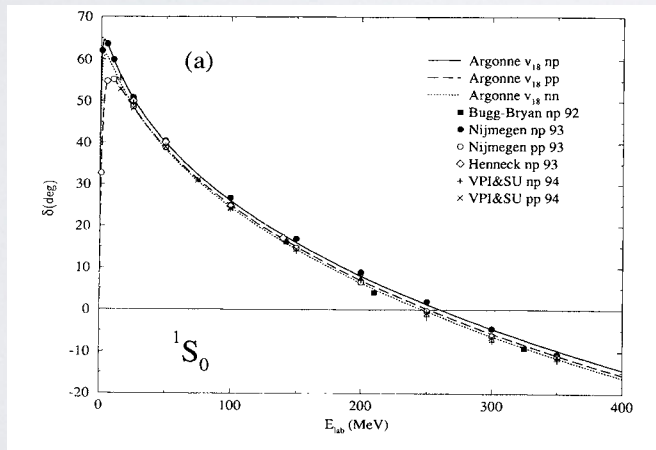
Very low densities dominated by  $^1S_0$  interaction

Very similar to cold atomic Fermi Gases

$$H = \sum_i \frac{p_i^2}{2m} + \sum_{i < j} V_0 \delta(\mathbf{r}_{ij})$$



Neutron-Neutron Scattering length  $\sim -18$  fm



pion + 2-pion + short-range repulsion

## Quantum Monte Carlo Methods

$$H = \sum_i \frac{p_i^2}{2m} + \sum_{i < j} V_{ij} + \dots$$

$$V_{ij} = \sum_k V_{ij}^k(r_{ij}) O_{ij}^k$$

$$O_{ij}^k = [1, \sigma_i \cdot \sigma_j, \sigma_i \cdot r_{ij} \sigma_j \cdot r_{ij}, L \cdot S_{ij}] \times [1, \tau_i \cdot \tau_j]$$

$$H \Psi = E \Psi$$

$$\Psi = \sum_{i=1}^{2^A \binom{A}{Z}} \psi(i)(\mathbf{R})$$

$2^A = 7 \times 10^{19}$  amplitudes  
for 66 neutrons  
in  $3A = 198$  dimensions

## Quantum Monte Carlo (Auxiliary Field Diffusion Monte Carlo)

$$\Psi_0 = \exp[-H\tau] \Psi_T$$

$$\exp[-H\tau] \approx \exp[-V\tau/2] \exp[-T\tau] \exp[-V\tau/2]$$

Kinetic Term is a diffusion process in 3A coordinates

Spin-dependent potential terms rewritten as coupled

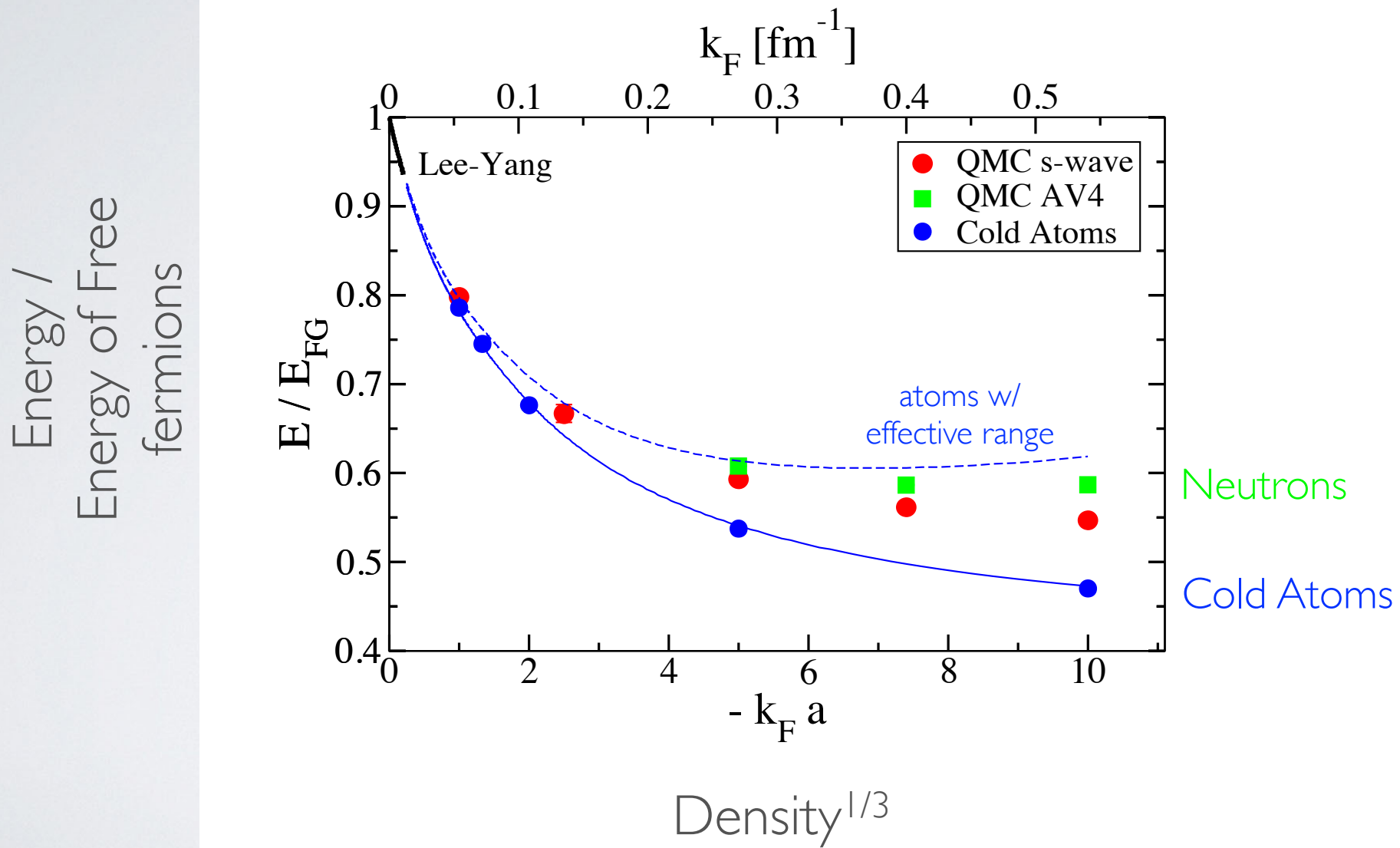
to an auxiliary field which is sampled by Monte Carlo,  
giving rotations of spins (and isospins)

$$\exp[-V\sigma_i \cdot \sigma_j \tau] = \sum_{x=\pm 1} \exp[-V^{1/2}\tau^{1/2}\sigma_i \cdot x] \exp[-V^{1/2}\tau^{1/2}\sigma_j \cdot x]$$

The simulation is a branching random walk in 3A coordinates  
and A spins and isospins.

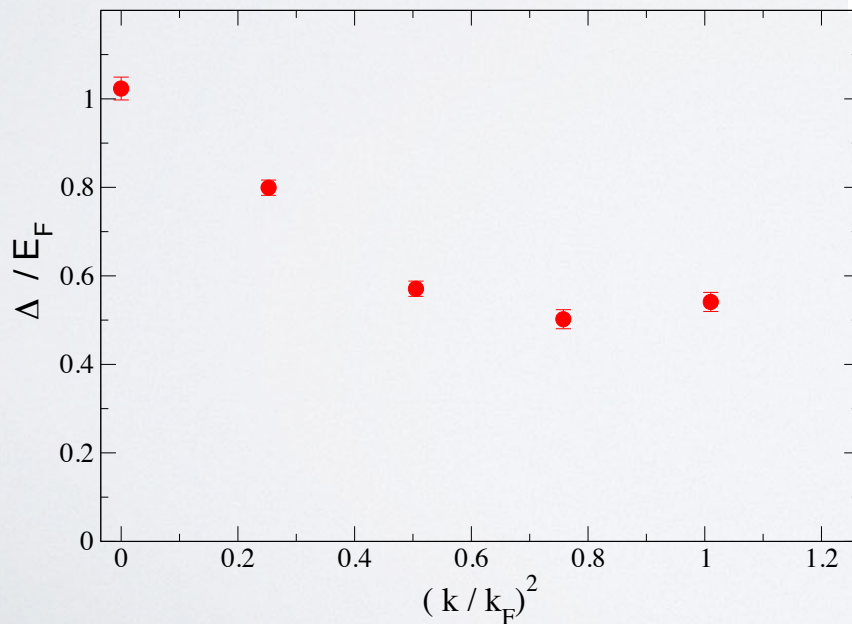
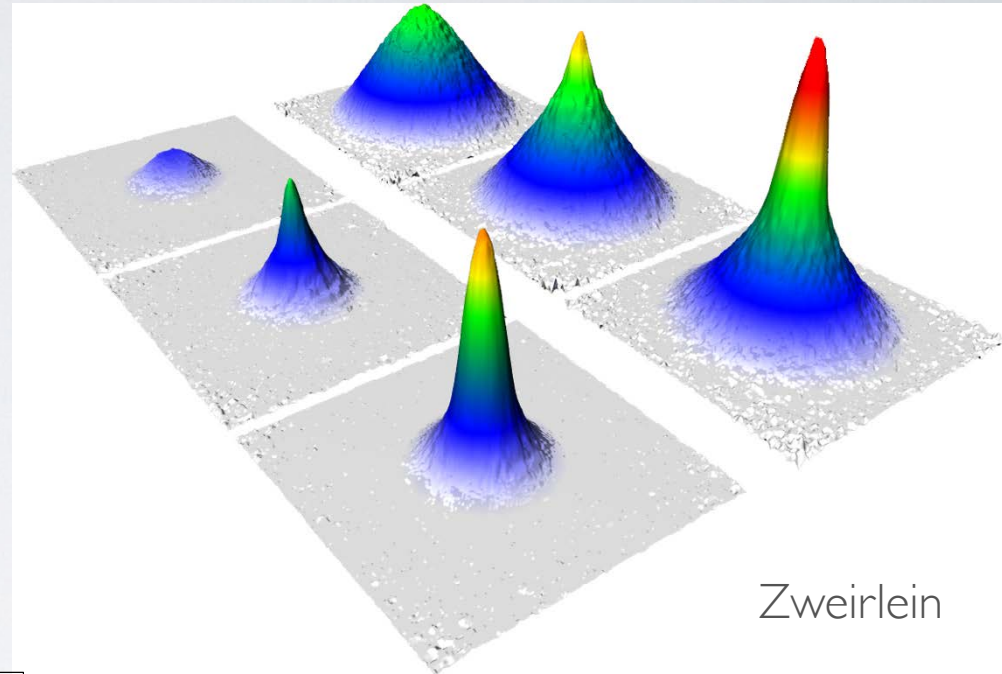
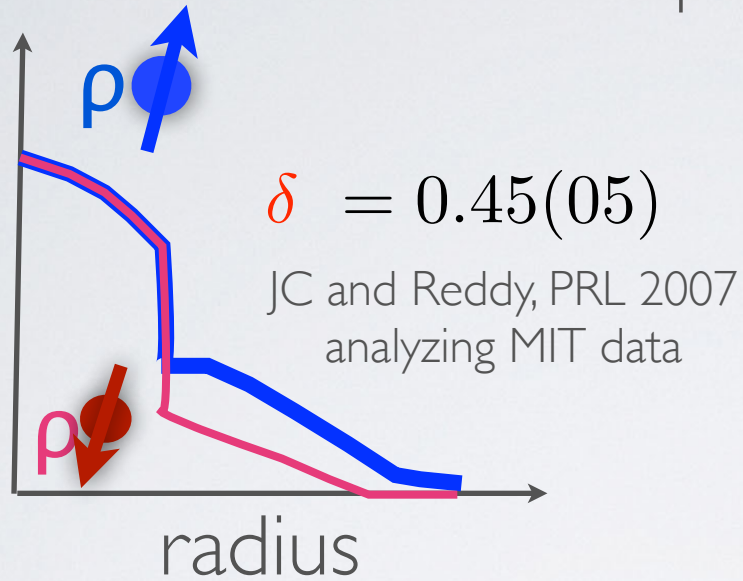


# Equation of State ( $E/A$ ) for neutrons and cold Fermi atoms



# Superfluidity (s-wave)

Spin up, down densities in a trap



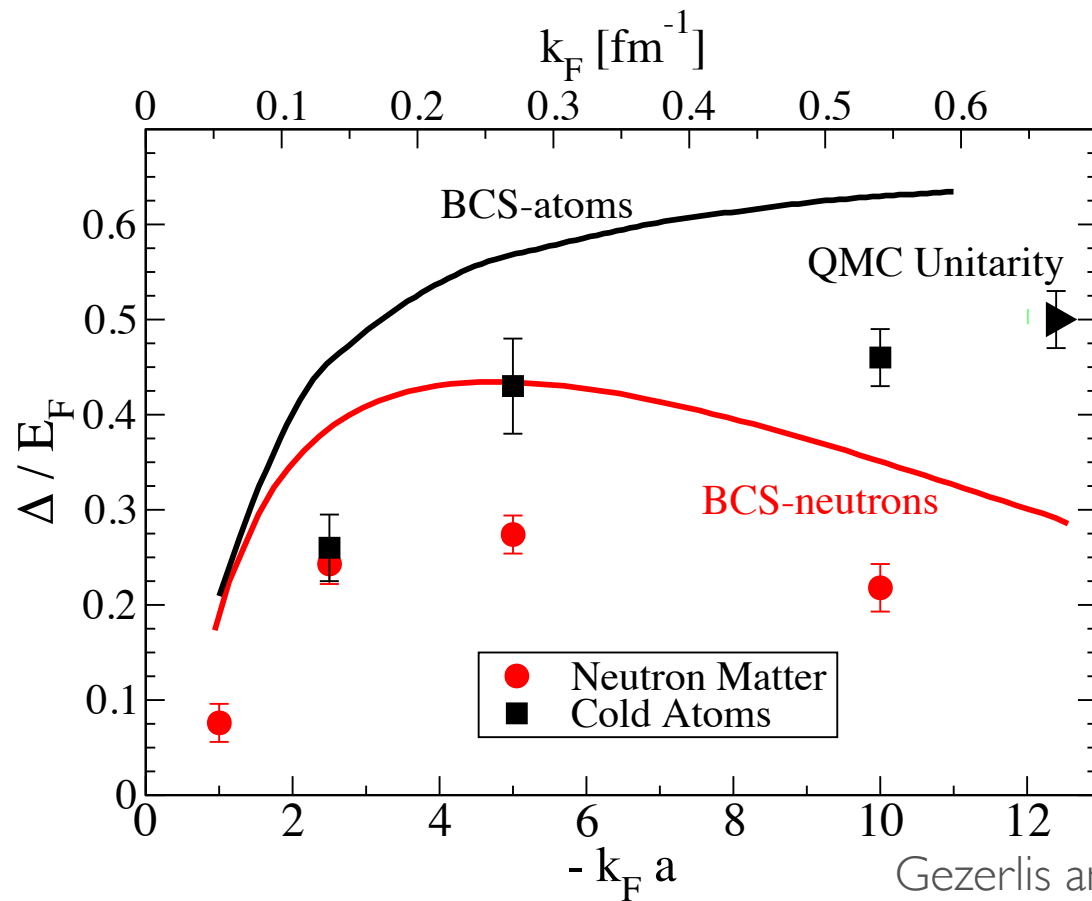
$$\Delta = \delta \frac{\hbar^2 k_F^2}{2m}$$

$$\delta = 0.50(03)$$

$$(k_{min}/k_f)^2 = 0.80(10)$$

JC and Reddy, PRL 2005

# Superfluid Pairing Gap

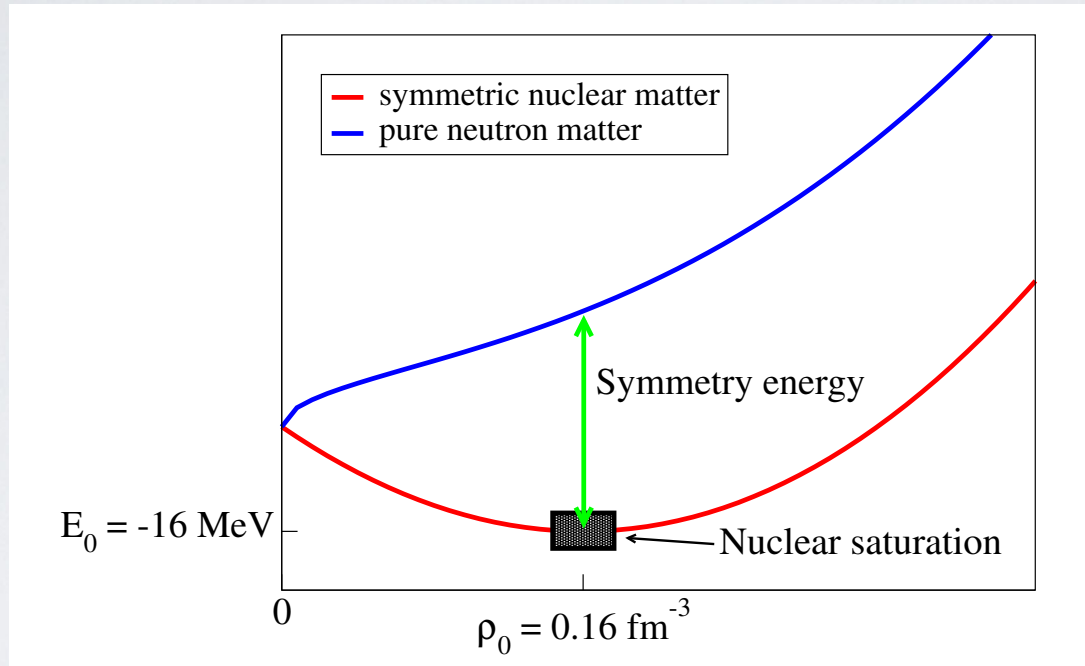


Solid lines:  
BCS mean-field theory  
Points w/ error bars:  
QMC

Gezerlis and Carlson, PRC 2010, 2012

Cold Atoms have highest superfluid gap /  $E_F$  of any system;  
Neutrons have highest pairing gap /  $E_F$  in nature.

## Equation of State at Higher Densities: near nuclear saturation

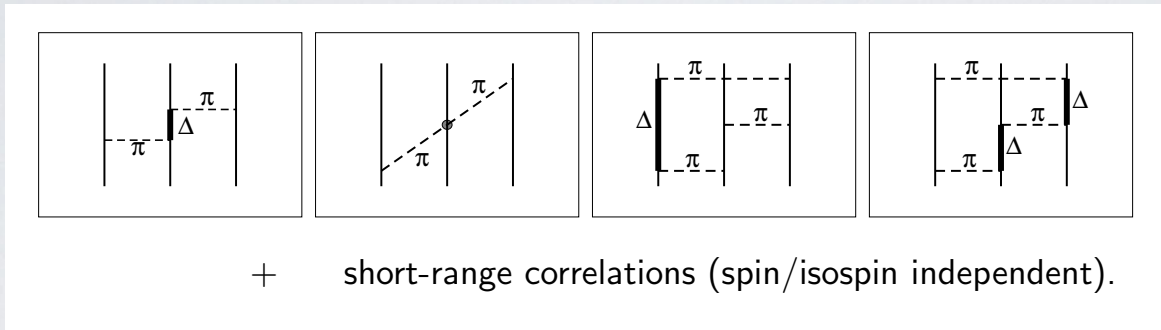


From experiments:

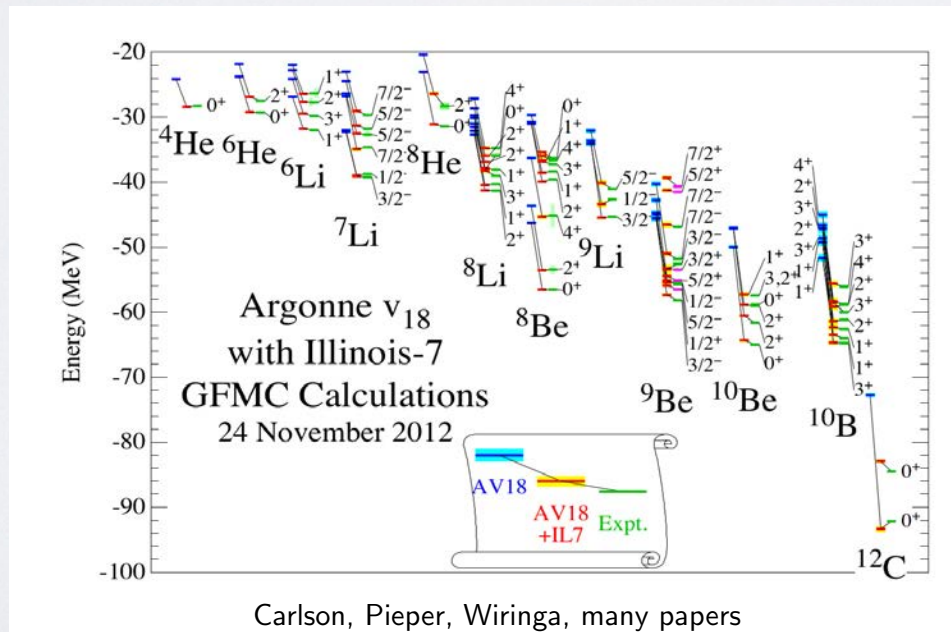
$$E_{SNM}(\rho_0) = -16 \text{ MeV}, \quad \rho_0 = 0.16 \text{ fm}^{-3}, \quad E_{sym} = E_{PNM}(\rho_0) + 16$$

The symmetry energy is accessible (indirectly) by experiment

At higher densities three-nucleon interactions start to become important

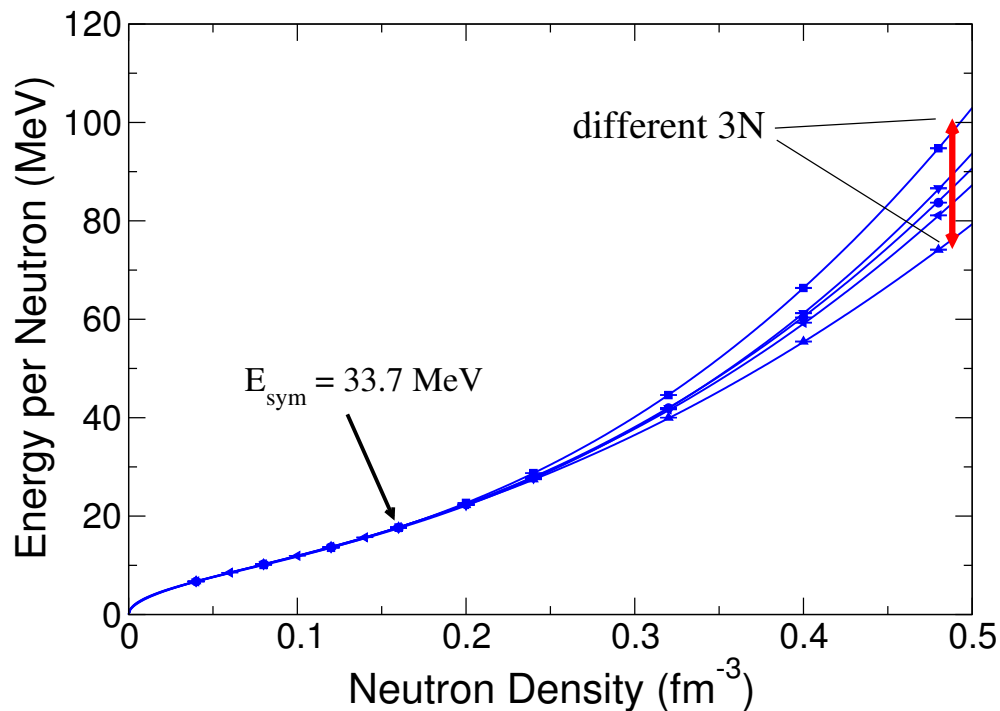


Calibrated to light nuclei



Consider a wide range of three-nucleon forces that give the same symmetry energy and then see how they extrapolate to high density

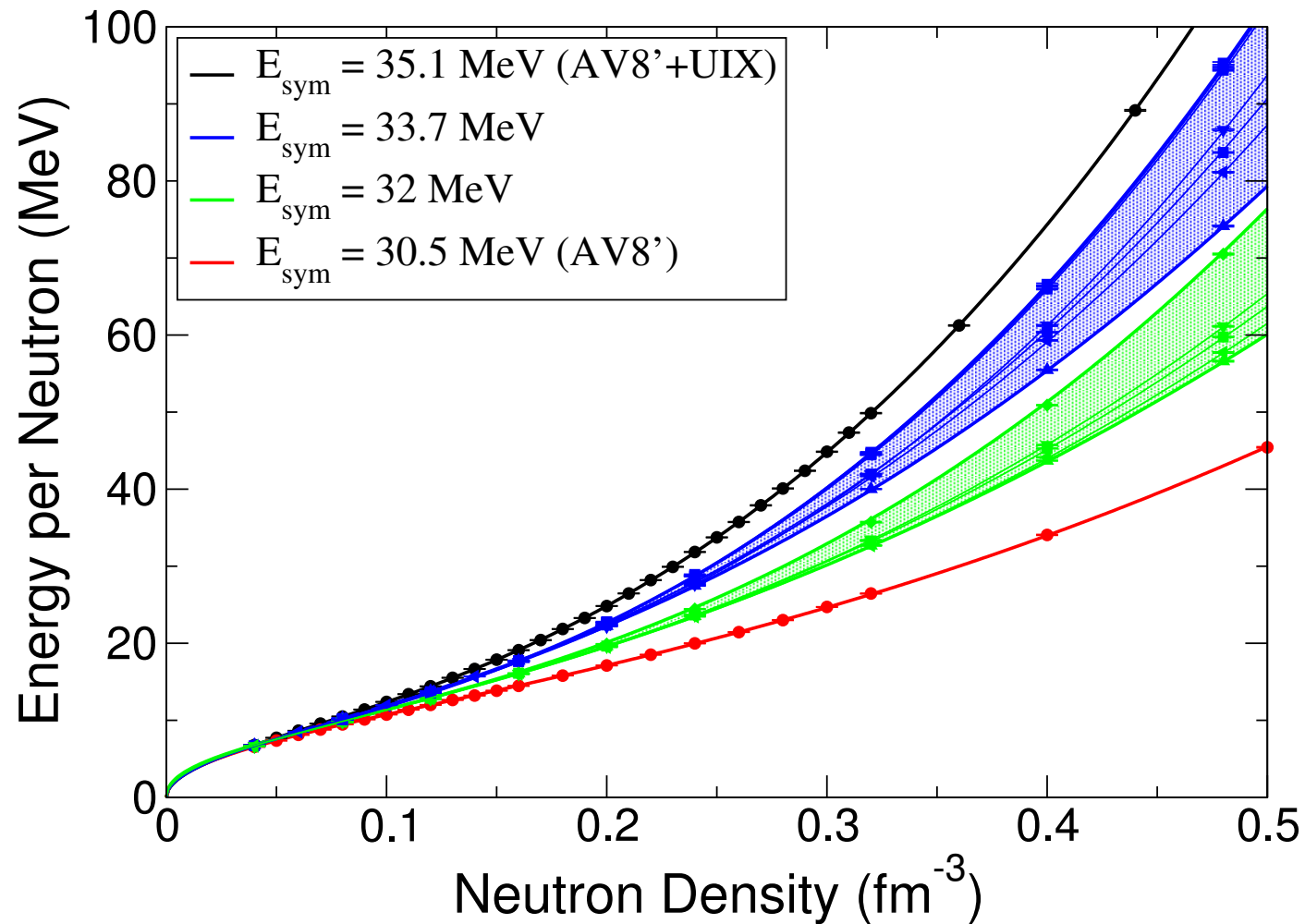
We consider different forms of three-neutron interaction by only requiring a particular value of  $E_{sym}$  at saturation.



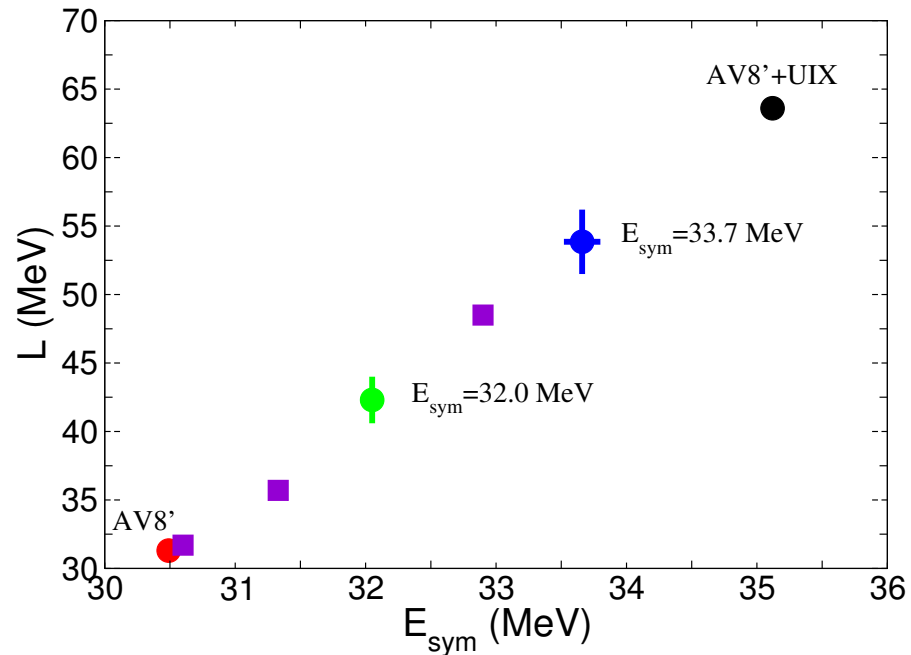
different 3N:

- $V_{2\pi} + \alpha V_R$
- $V_{2\pi} + \alpha V_R^\mu$   
(several  $\mu$ )
- $V_{2\pi} + \alpha \tilde{V}_R$
- $V_{3\pi} + \alpha V_R$

# Equations of state with a fixed symmetry energy



# Strong Correlation between Symmetry Energy and its Derivative

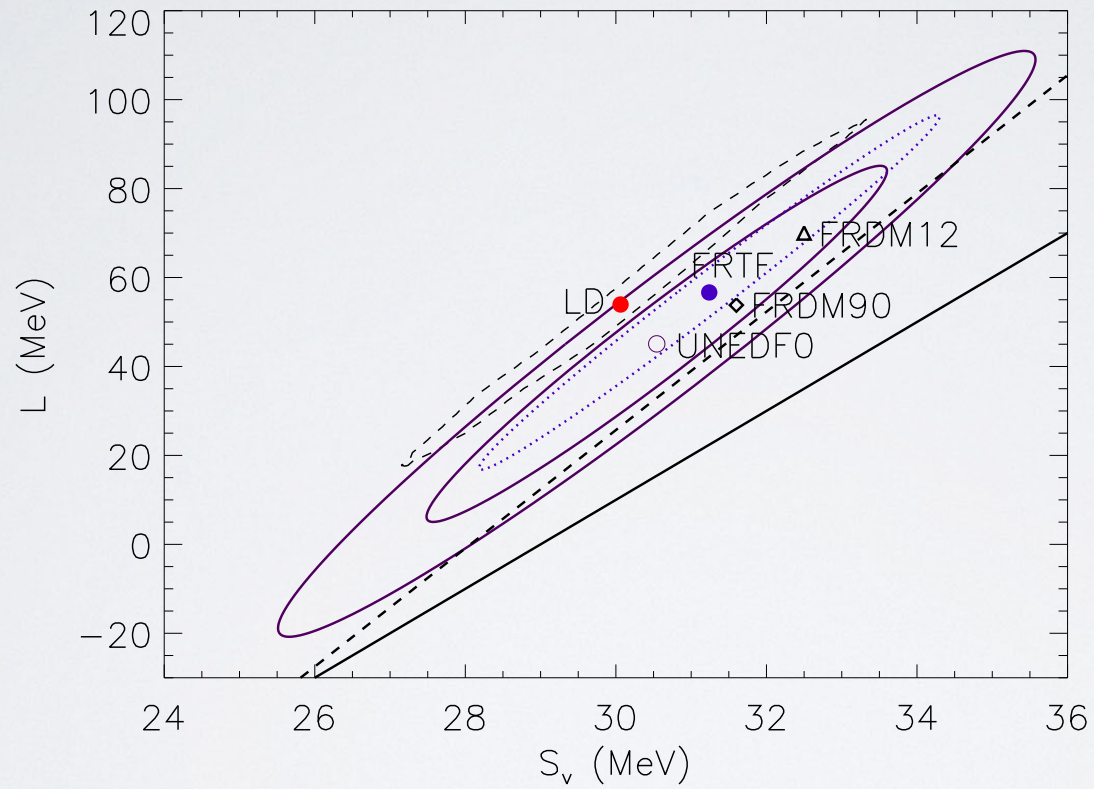


Gandolfi *et al.*, EPJ (2014)

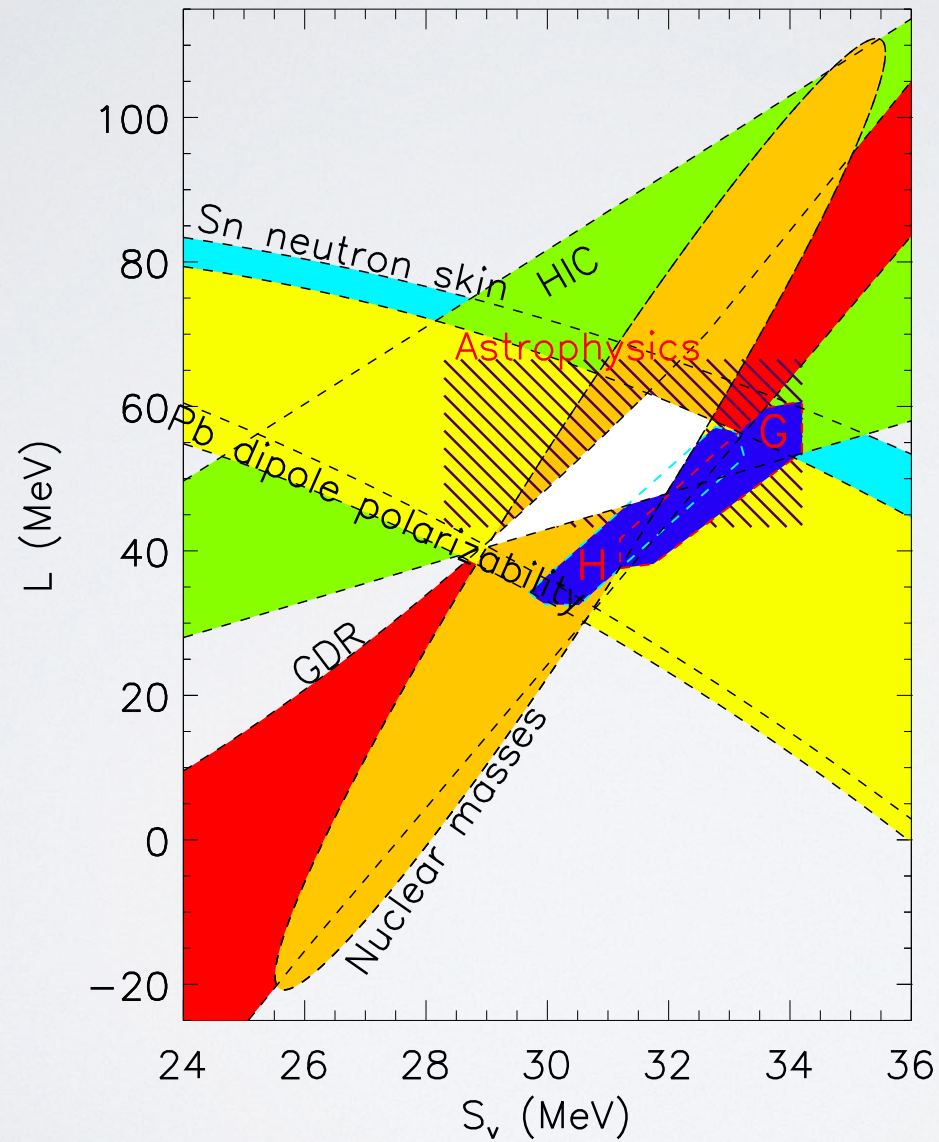
New chiral interaction models give very similar results



# Fits to nuclear masses



# Variety of Experimental Constraints

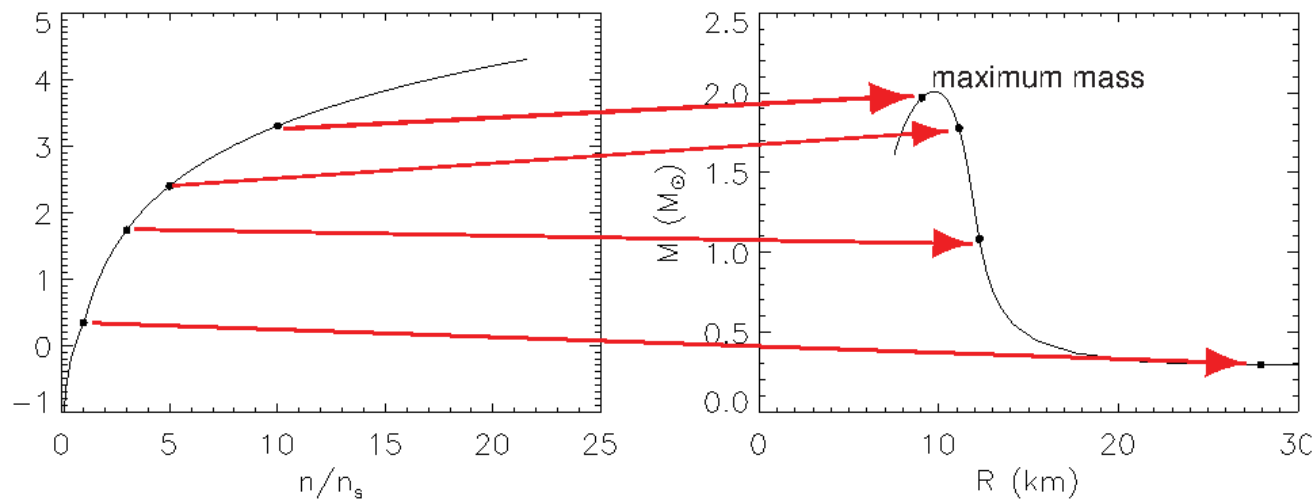


# Equation of State to Mass / Radius

TOV equations:

$$\frac{dP}{dr} = -\frac{G[m(r) + 4\pi r^3 P/c^2][\epsilon + P/c^2]}{r[r - 2Gm(r)/c^2]},$$

$$\frac{dm(r)}{dr} = 4\pi\epsilon r^2,$$



from Lattimer

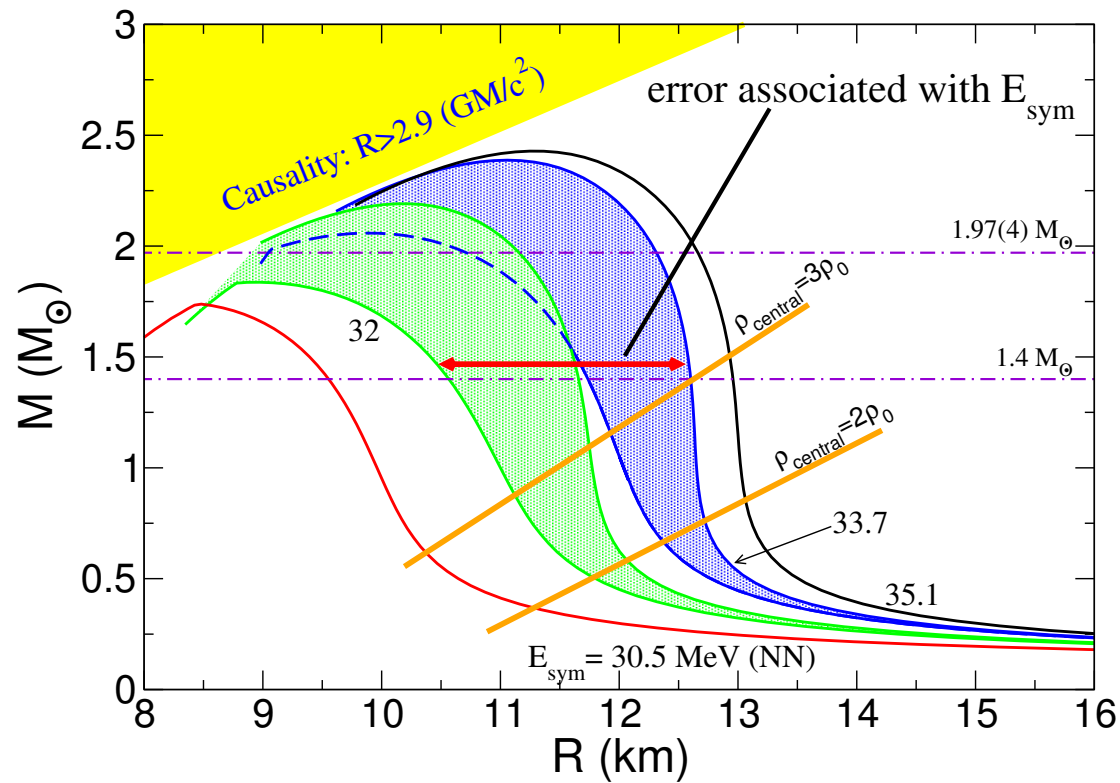
Tolman Oppenheimer Volkov equations: 1939

used free neutron gas to estimate upper bound of 0.7 solar masses

see Silbar and Reddy: [arXiv:nucl-th/0309041](https://arxiv.org/abs/nucl-th/0309041) for an introduction

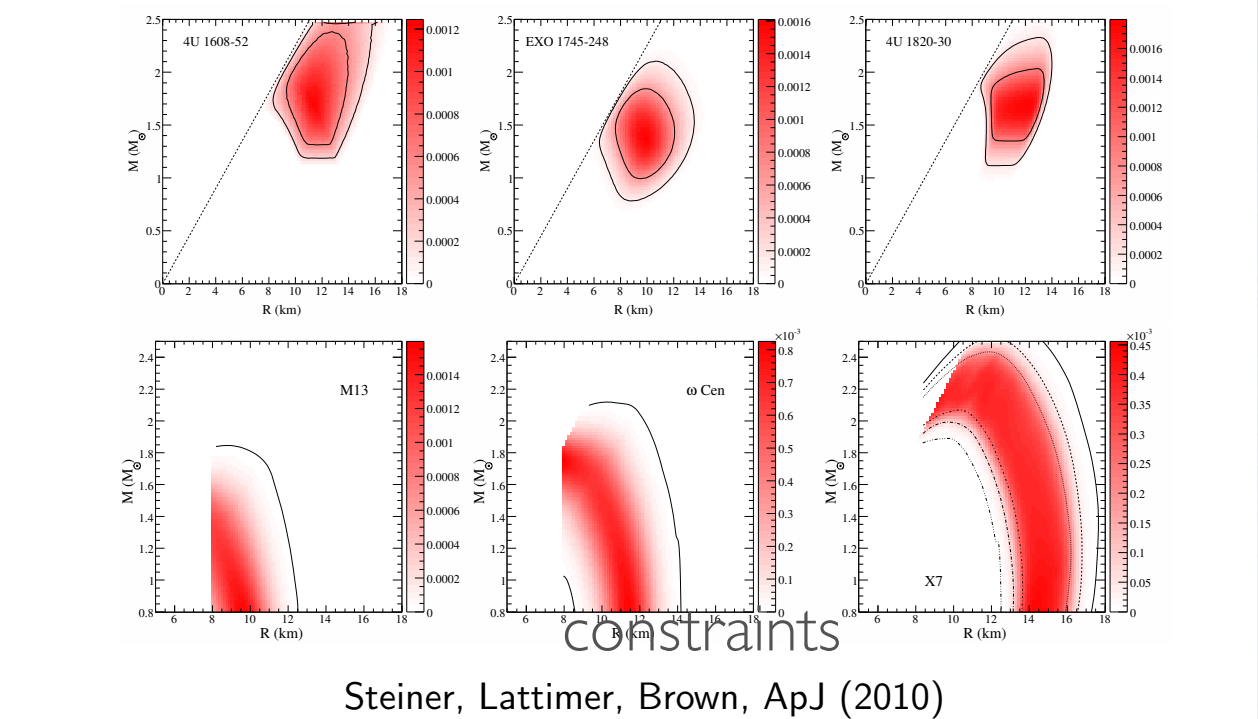
# Neutron Star Mass/Radius: Calculations

EOS used to solve the TOV equations.



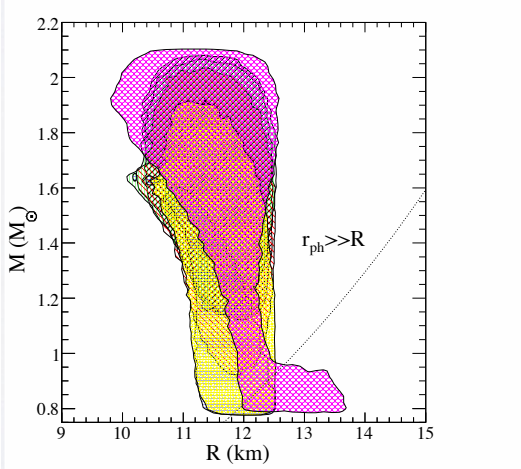
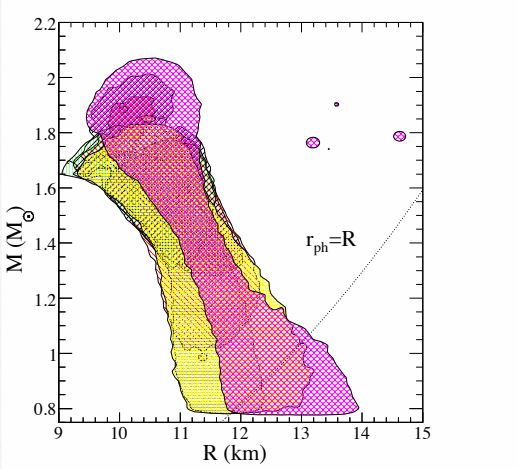
Gandolfi, Carlson, Reddy, PRC (2012).

# Observations - still controversial

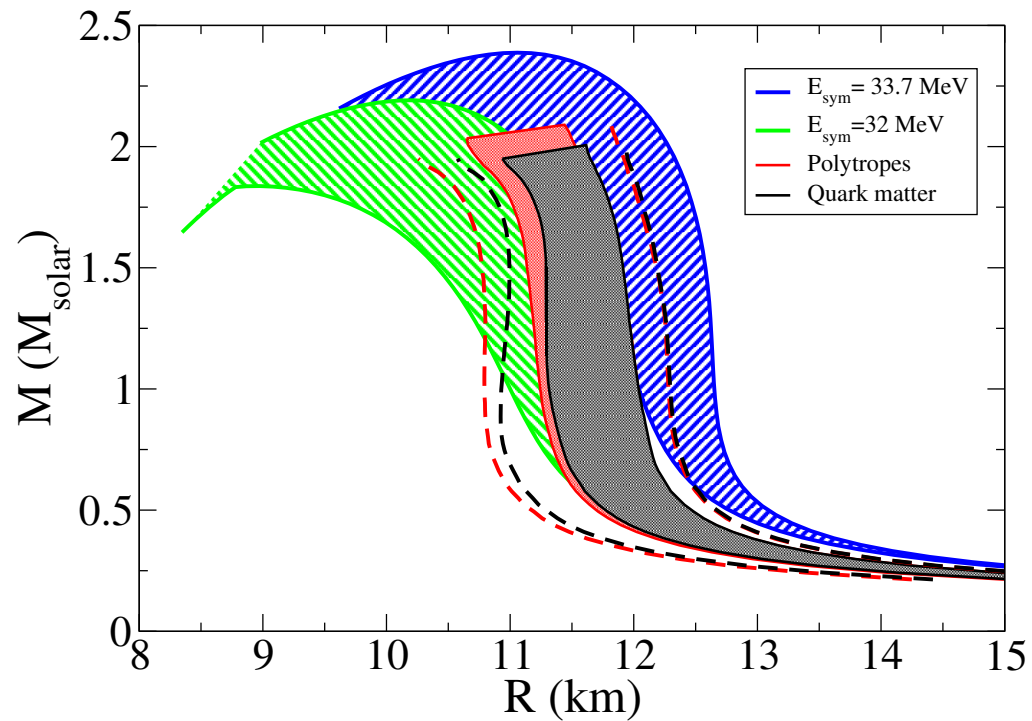


constraints from individual stars  
 observations from  
 3 X-ray bursters  
 plus 3 low-mass X-ray binaries

Mass radius constraints  
 subject to assumptions



# Comparison of theory and observations



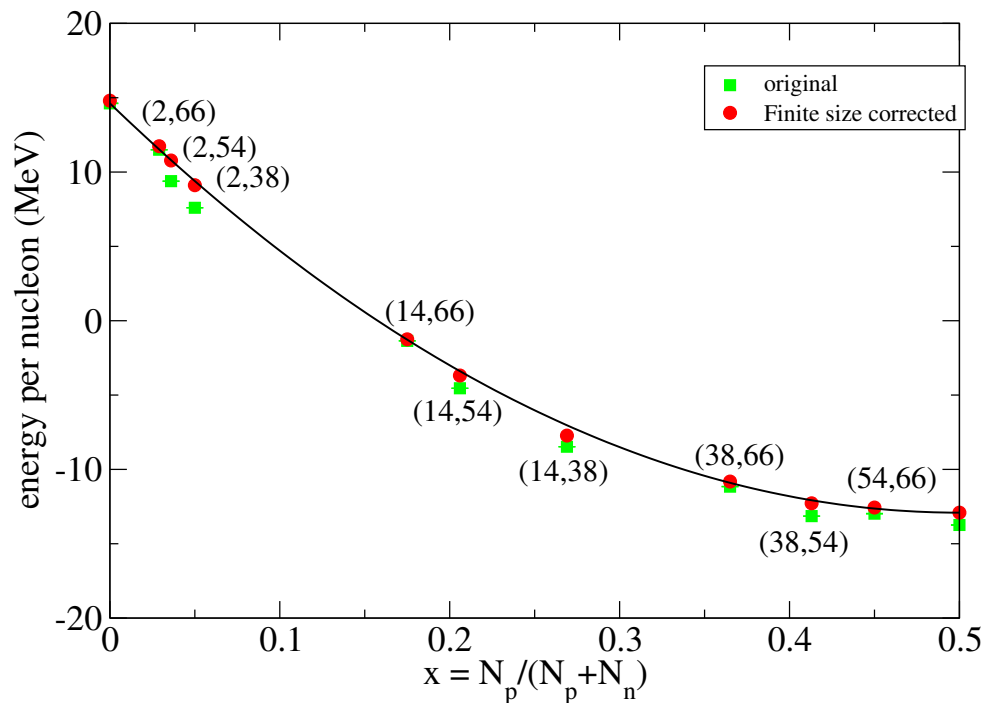
$$32 < E_{\text{sym}} < 34 \text{ MeV}, \quad 43 < L < 52 \text{ MeV}$$

Steiner, Gandolfi, PRL (2012).

What about other particles? protons

Quadratic dependence of  $E$  versus  $n/p$  imbalance

Asymmetric nuclear matter  $E(\rho, x) = E_{SNM}(\rho) + E_{sym}^{(2)}(\rho)(1 - 2x)^2$

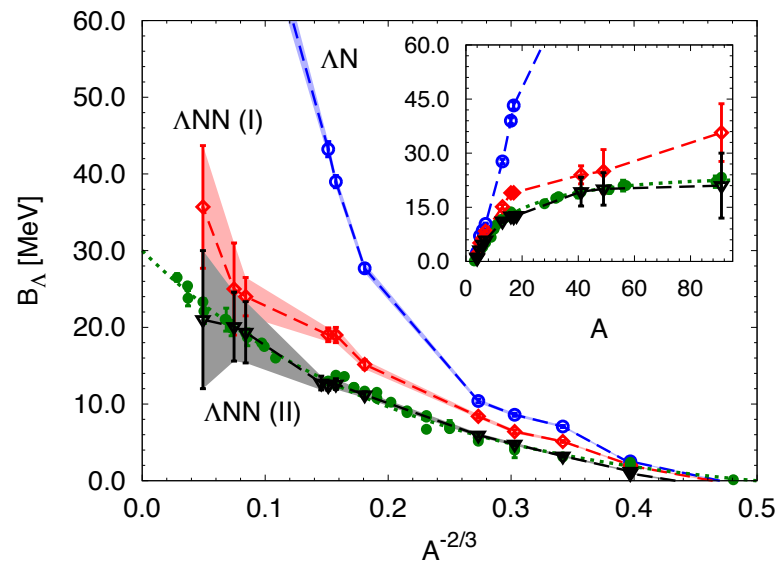


Gandolfi, Lovato, Carlson, Schmidt, arXiv:1406.3388

proton fraction also important for neutrino processes

What about other particles? hyperons, ...

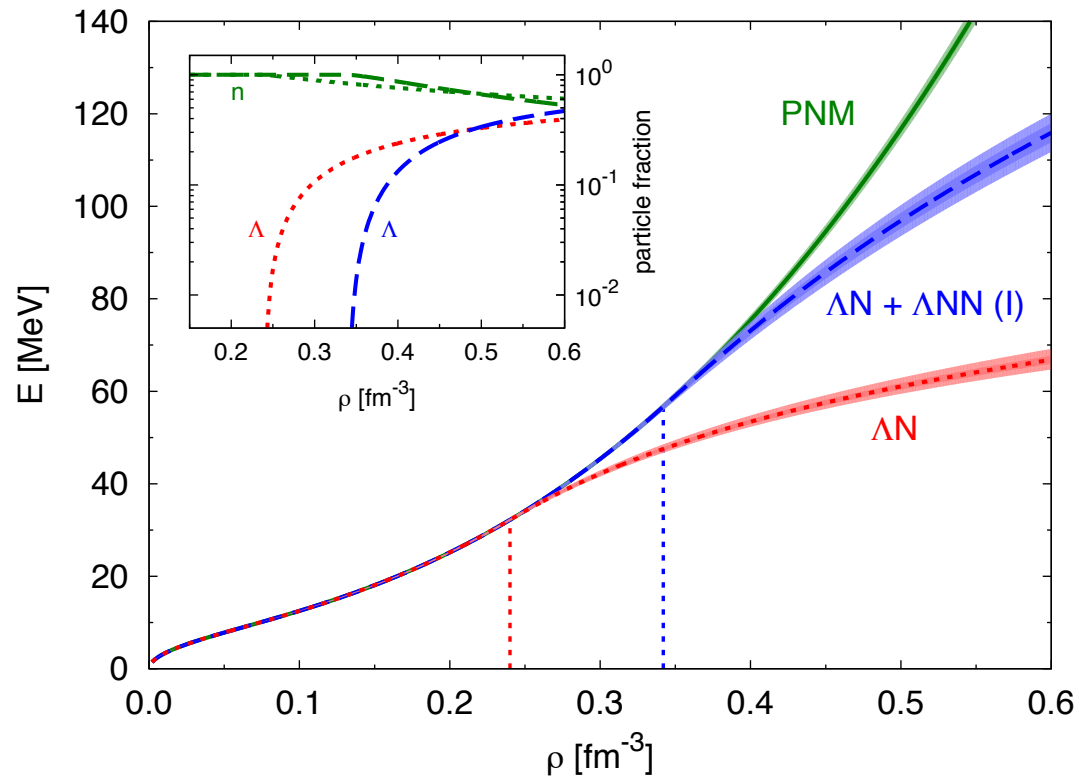
Hyperons are bound in nuclei by  $\sim 30$  MeV.  
What happens in dense matter?



Lonardonì, Pederiva, SG, PRC (2013) and PRC (2014).



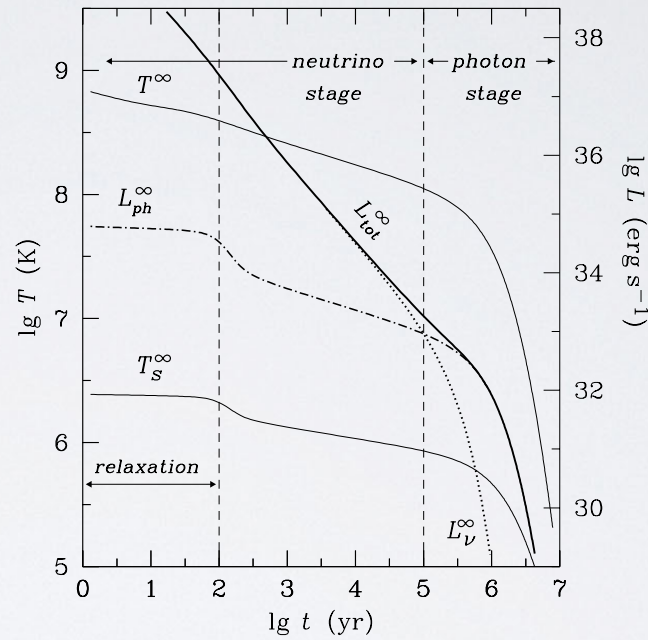
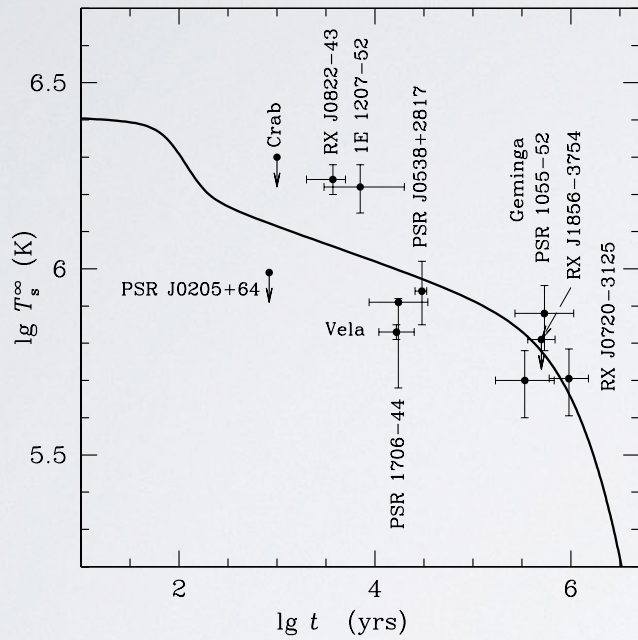
# Hyperons in Neutron Matter



Lonardonì, Lovato, Gandolfi, Pederiva, arXiv:1407.4448 (2014)

Best model gives no hyperons  
up to 3-4  $\times$  saturation density

# Neutrinos in neutron stars and proto-neutron stars



Yakovlev 2004

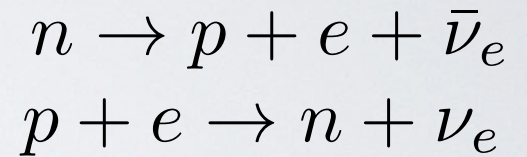
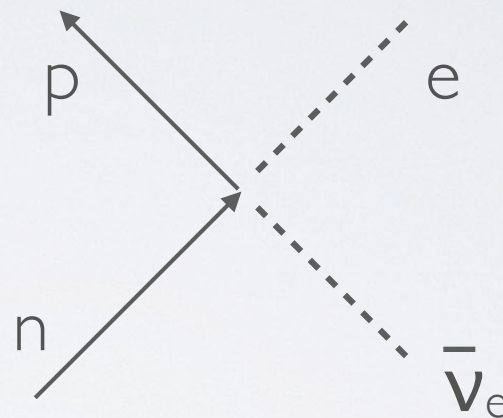
# Neutron Star Cooling Introduction

Sensitive to:

Equation of state  
Neutrino Emission  
Superfluidity  
Magnetic Fields  
Surface

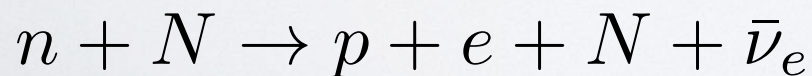
Direct Urca:

Lattimer, Pethick, Prakash, Haensel (1991)



threshold associated with Fermi surfaces limit this to  $\rho > 2 \rho_0$   
Requires  $\sim 15\%$  proton fraction  
to satisfy energy and momentum conservation

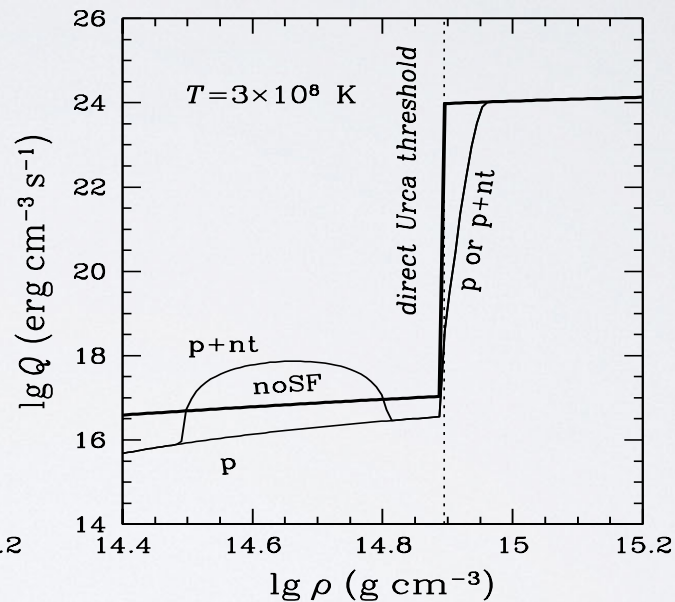
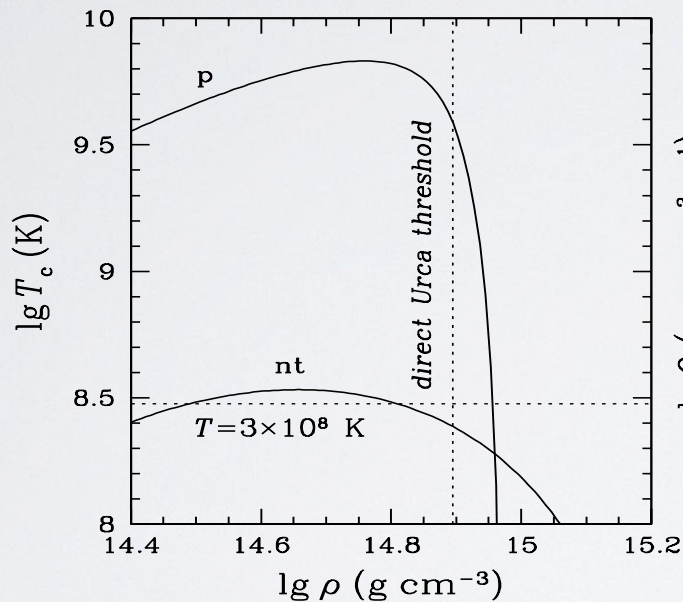
modified Urca works throughout the core



much slower

# Superfluidity

suppresses familiar neutrino processes  
 creates new process: production through Cooper pairing  
 3P2 - 3F2 pairing particularly important but not well constrained



$3 \times 10^8 \text{ K} \sim .026 \text{ MeV}$

typical s-wave pairing gaps  $\sim 1 \text{ MeV}$

angle average 3P2  $\sim .01 \text{ MeV}$

$\log(\text{nuclear saturation density in } \text{g/cm}^3) \sim 14.4$

Yakovlev, et al, 2004

also see Gezerlis, Pethick, Schwenk arXiv:1406.6109

# Neutron and Proto-Neutron Star Cooling

Neutron star cooling depends upon

Equation of State

Neutrino Emission and Propagation

Neutron (and proton) Superfluidity

+ ...

Supernovae neutrino emission also depends upon

weak response of matter

interesting regime at low densities ( $0.1 \rho_0$ )

and moderate temperatures (non-degenerate matter)

Rapid progress in theory and observations

# Summary/ Outlook

Rapid progress in our understanding of cold dense matter

Excellent connections to  
Theory of strongly-correlated matter  
Experiments in cold atom physics  
Astrophysical observations  
Future measurements of gravitational waves  
Supernovae physics and neutrino physics

