



Theoretical Modeling of Cosmic Structures



Growth Spurting Baby Galaxies

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with

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The First Billion Years



Direct Observational Evidence



Stellar assembly rate

Approximating stellar mass growth due to equal mass mergers by: $\dot{M}_{grow} = \frac{M_*}{t_{df}}$

At z > 6 one equal mass merger can grow the stellar mass more than constant SF over the next ~0.6 Gyr. Mergers could be more important to the evolution of the mass function than star formation depending on their frequency.



Dutton et al 2010

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Simulation

- GADGET-2 version used for the OWLS project (Schaye et al. 2010): SF; metal enrichment; metal line cooling from 11 elements; BH growth and feedback; SN feedback
- Added molecular networks and cooling
- Added POPIII formation and evolution, seed BHs
- Added dust from PISN, AGB & SNII; thermal sputtering
- Inclusion of Lyman-Werner background (11.3 - 13.6 eV)
- Coupled to radiative transfer scheme SIMPLEX in postprocessing

ΙΜοΟς



The First Billion Years Simulation

Theoretical Modeling of Cosmic Structures Max Planck Research Group Max Planck Institute for Extraterrestrial Physics



http://www.mpe.mpg.de/tmox/



 $V = (8 \text{Mpc})^3$ $N = 2 \times 1368^3$ $m_{gas} = 890 M_{\odot} h^{-1}$ $m_{DM} = 4375 M_{\odot} h^{-1}$



The Mass Function



Star formation



Khochfar et al. 2012

Star formation



Khochfar et al. 2012

Star formation histories

Mass



SN-Feedback

Density + Temp

Metallicity



Movie: C. Dalla Vecchia

Reionization

SimpleX code including H + He ionization, following 10 frequency bands

- Galaxies in haloes below 10⁸ M_{sun} are able to reionize the Universe initially
- Massive haloes only host dominant sources of photons at later times

$$\begin{aligned} \frac{\mathrm{d}Q_{\mathrm{H\,II}}}{\mathrm{d}t} &= \frac{\dot{N}_{\mathrm{ion}}}{\bar{n}_{\mathrm{H,0}}} - Q_{\mathrm{H\,II}}C\bar{n}_{\mathrm{H,0}}\alpha(T)(1+z)^{3}\\ \tau_{\mathrm{e}} &= \int_{0}^{z_{\mathrm{rec}}} \mathrm{d}z \left|\frac{\mathrm{d}t}{\mathrm{d}z}\right| Q_{\mathrm{H\,II}}(z)\bar{n}_{\mathrm{H,0}}(1+z)^{3}\sigma_{\mathrm{T}} \end{aligned}$$



J.-P. Paardekooper, SK, Dalla Vecchia 2012

Sources Reionizing the Universe



J.-P. Paardekooper, SK, Dalla Vecchia 2012

Escape fractions

- Escape fraction are a strong function of time
- Feedback
 increases the
 escape fraction
- Low mass galaxies have higher escape fraction due to a more dramatic effect of feedback



J.-P. Paardekooper, SK, Dalla Vecchia 2012

Dark Matter Profiles



Distribution of profiles

Roughly 50% of massive haloes at z~6 show a density decrement, with a long tail to density enhancements.



Davis, SK 2012

Baryon-DM connection

Adiabatic contraction:

 $r_i M_i = r_f (M_i + \Delta M)$

Enhancement of baryons leads to dark matter response, but not always! Some haloes have less dark matter in them even if the baryon fraction is increased. Adiabatic contraction is clearly not working in those cases.



Davis, SK 2012

Feedback impact on the halo



Davis, SK et al 2012

Off-set galaxies

- Galaxies show an offset with respect to the DM potential minimum. Systematic larger off-set in haloes with reduced dark matter fraction.
- Galaxies could heat dark matter halo via dynamical friction



Davis, SK 2012

Conclusions

- The Universe is initially reionized by baby galaxies and kept ionized by massive galaxies.
- Baby galaxies with M* <10⁶ M_{sun} can reionize the Universe due to higher escape fraction then massive galaxies
- SN-feedback drives the escape fractions
- Dark matter haloes can contract due to the presence of baryons consistent with AC
- 'Offset' galaxies can create cores in DM haloes due to the exchange of angular momentum
- Inner dark matter profiles have transient feature