LECTURE IL HOW STUDYING THE FUTURE EVOLUTION OF THE COSMOS AFFECTS OUR UNDERSTANDING OF THE UNIVERSE

Fred C. Adams University of Michigan UC Santa Cruz -- July 2013

Five Ages of the Universe

Primordial Era n < 6
Stelliferous Era n = 6 - 14
Degenerate Era n = 14 - 40
Black Hole Era n = 40 - 100
Dark Era n > 100

where *n* defined by $t = 10^n$ years

The Primordial Era

The `Big Bang Moment' Inflation Matter > Antimatter Quarks --- protons & neutrons Dark matter abundances are frozen Nuclear synthesis of the light elements Cosmic Microwave Background Universe continues to expand

The Stelliferous Era

 Stars dominate energy production Lowest mass stars increasingly important Biosphere ends in 3.5 Gyr Earth as a planet ends in 7 Gyr Odds of Earth escape (capture) are about one part in 10^5 (3x10^6) Most liquid water found inside frozen planets Star formation and stellar evolution end near (by) cosmological decade n = 14

The Degenerate Era

The Black Hole Era

Black holes are brightest stellar objects
Generation of energy via Hawking radiation
Every galaxy contributes one supermassive and about one million stellar black holes
Black hole lifetime is mass dependent:

One solar mass: Million solar mass: Galactic mass: Horizon mass: n=65 n=83 n=98 n=131

 $t_T \propto M_{bh}^3$

The Dark Era

No stellar objects of any kind
Inventory of elementary particles: electrons, positrons, neutrinos, & photons
Positronium formation and decay
Low level annihilation
Vacuum tunneling events?

Ashes to ashes, Dust to dust, Particles to particles, Our universe will evolve through a a one-way time-line. Other universes can live through their own time-lines, as parts of The MULTIVERSE.



Lecture I Summary

- Our current understanding of the laws of physics and astrophysics allow us to construct a working picture of the future.
- Studying physical processes of the future provides insight into current astrophysical problems, e.g., the reason for red giants, structure of dark matter halos, dynamical scattering problems, defining the masses of galaxies, etc.

Outline for Lecture II

 Why stars become red giants How to define the mass of a galaxy The asymptotic form of galactic halos How common are stars in alternate universes (other parts of the multiverse) Time-varying constants of nature Even longer term astrophysical processes Discussion issues

Why do stars become RED GIANTS ?



Long term Evolution of Red Dwarfs



Life Span of Red Dwarfs





Evolution of star with m = 0.10

Late time light curve for Milky Way



Radial size vs time for varying masses



Radial size vs time for varying masses



Photosphere evolution:





Why Do Stars Become RED GIANTS ?

Luminosity increases $L_* = 4\pi R_*^2 \sigma T_*^4$

Star can either become large or hot. If photosphere has an opacity wall, then the star cannot become hot, so it must become large => Giant!

Opacity Analysis

Radiation equation $L = -4\pi r^2 \frac{4acT^3}{3\kappa\rho} \frac{dT}{dr}$

Hydrostatic equilibrium $\frac{dT}{dr} = -\frac{1}{1+n} \frac{\mu GM_r}{r^2 R_g}$

$$L_* = \frac{16\pi}{3} \frac{acG}{R_g} \frac{\mu M_*}{1+n} \frac{T_*^3}{\kappa \rho}$$

 $\kappa = C \rho^{\alpha} T^{\omega}, \ \rho \propto R_*^{-\gamma}, \& L_* = 4 \pi R_*^2 \sigma T_*^4$

ΔT_{*}	α	ΔL_{*}	ΔR_{*} _	<i>ω</i> + 1	ΔL_{*}
T_*	ω + 1 + 4 α	L_*	R_*	$\overline{2(\omega+1+4\alpha)}$	L_*

What is the total mass of a galaxy? Why do dark matter halos have a nearly universal form?



(M. Busha, F. Adams, et al. 2003, 2005, 2007)



Dark matter halos approach a well-defined asymptotic form with unambiguous total mass, outer radius, & density profile



(Busha et al. 2005)

Phase Space of Dark Matter Halo



(M. Busha et al. 2005)

Spacetime Metric Attains Universal Form



ORBITS?

Most of the mass is in dark matter
Most dark matter resides in these halos
Halos have the universal form found here (nfw/hq) for most of their lives
Most orbital motion that will EVER occur will be THIS orbital motion in DM halos

$$\left(factor of 10^{74}\right)$$

Triaxial Forces

$$\begin{split} F_{x} &= \frac{-2 \operatorname{sgn}(x)}{\sqrt{\left(a^{2}-b^{2}\right)\left(a^{2}-c^{2}\right)}} \ln \left(\frac{2G(a)\sqrt{\Gamma}+2\Gamma-a^{2}\Lambda}{2a^{2}\xi G(a)+\Lambda a^{2}-2a^{4}\xi^{2}}\right) \\ F_{y} &= \frac{-2 \operatorname{sgn}(y)}{\sqrt{\left(a^{2}-b^{2}\right)\left(b^{2}-c^{2}\right)}} \left[\sin^{-1} \left(\frac{\Lambda-2b^{2}\xi^{2}}{\sqrt{\Lambda^{2}-4\Gamma\xi^{2}}}\right) - \sin^{-1} \left(\frac{2\Gamma/b^{2}-\Lambda}{\sqrt{\Lambda^{2}-4\xi^{2}\Gamma}}\right) \right] \\ F_{z} &= \frac{-2 \operatorname{sgn}(z)}{\sqrt{\left(a^{2}-c^{2}\right)\left(b^{2}-c^{2}\right)}} \ln \left(\frac{2G(c)\sqrt{\Gamma}+2\Gamma-c^{2}\Lambda}{2c^{2}\xi G(c)+\Lambda c^{2}-2c^{4}\xi^{2}}\right) \end{split}$$

(Adams, Bloch, Butler, Druce, Ketchum 2007)

$$G(u) = \xi^{2}u^{4} - \Lambda u^{2} + \Gamma$$

$$\xi^{2} = x^{2} + y^{2} + z^{2}$$

$$\Lambda = (b^{2} + c^{2})x^{2} + (a^{2} + c^{2})y^{2} + (a^{2} + b^{2})z^{2}$$

$$\Gamma = b^{2}c^{2}x^{2} + a^{2}c^{2}y^{2} + a^{2}b^{2}z^{2}$$

Orbit Gallery







Orbit Gallery







INSTABILITIES



Unstable motion shows:
(1) exponential growth,
(2) quasi-periodicity,
(3) chaotic variations,
(4) eventual saturation

Orbits in any of the principal planes are unstable to motion perpendicular to the plane.





Given the possible existence of multiple universes, we face the question: Do other universes have different versions of the laws of physics? Do other versions of the laws of physics allow life to develop?

Stars and Stellar Structure in Other Universes

F. Adams (2008) JCAP, 08010

OVERVIEW

 Build robust, analytic stellar model Polytropic model of mechanical structure Only one nuclear burning species Radiative energy transport Define solution space: 4 forces of nature reduced to 3 parameters: (G, α, C)

Large fraction of this parameter space allows for the existence of working stars

Stellar Structure Equations

 $\frac{\partial P}{\partial r} = -\frac{GM(r)\rho}{r^2}, \quad \frac{\partial M}{\partial r} = 4\pi r^2\rho$ $\frac{\partial T}{\partial r} = -\frac{3\kappa\rho\ell}{64\pi r^2\sigma T^3}$ $\frac{\partial \ell}{\partial r} = 4\pi r^2 (\varepsilon - \varepsilon_v),$ $P(\rho,T,\ldots) \quad \kappa(\rho,T,\ldots) \quad \varepsilon(\rho,T,\ldots)$

Polytropic Stars $P = K \rho^{\Gamma} = K \rho^{1+1/n}$ $\xi = r/R, \ \rho = \rho_C f^n, \ R^2 = \frac{K\Gamma/(\Gamma-1)}{4\pi G\rho_C^{2-\Gamma}}$ $\frac{d}{d\xi} \left(\xi^2 \frac{df}{d\xi} \right) + \xi^2 f^n = 0$

Radial Profiles: Mechanical Structure



Nuclear Reaction Rates

$$\sigma(E) = \frac{S(E)}{E} \exp\left[-\left(\frac{E_G}{E}\right)^{1/2}\right] \qquad E_G = \left(\pi \alpha Z_1 Z_2\right)^2 2m_R c^2$$

$$\varepsilon = C\rho^2 \Theta^2 \exp[-3\Theta] \quad \Theta = (E_G/4kT)^{1/3}$$

where
$$C = \frac{8(\Delta E)S(E_0)}{\sqrt{3} \pi \alpha m_1 m_2 Z_1 Z_2 m_R c}$$

(thermal distribution of particle energies + Coulomb barrier + quantum tunneling)



Central Temperature

specify opacity:
$$\kappa \approx A\rho T^{-7/2}$$
 (Kramer)

after some algebra:

$$\Theta_{C}I(\Theta_{C})T_{C}^{3} = \frac{(4\pi)^{3}ac}{3\beta\kappa_{c}C} \left(\frac{M_{*}}{\mu_{0}}\right)^{4} \left[\frac{G}{(n+1)R_{gas}}\right]^{7}$$
where $I(\Theta_{C}) = \int_{0}^{\xi_{*}} f^{2n}\xi^{2}\Theta^{2}\exp[-3\Theta]d\xi$

Stellar Structure Solutions

$$R_* = \frac{GM_* \langle m \rangle}{kT_C} \frac{\xi_*}{(n+1)\mu_0}$$
$$L_* = \frac{(2\pi)^7}{15h^3c^2\beta\kappa_0\Theta_C} \left(\frac{M_*}{\mu_0}\right)^3 \left[\frac{G\langle m \rangle}{(n+1)}\right]^4$$
$$T_* = \left[\frac{L_*}{4\pi R_*^2\sigma}\right]^{1/4}$$

Hertzsprung-Russell Diagram



Luminosity-Mass Relationship



Minimum & Maximum Masses

characteristic mass scale:

$$M_{0} = \left(\frac{hc}{2\pi G}\right)^{3/2} m_{P}^{-2} = \frac{M_{Pl}^{3}}{m_{P}^{2}} \approx 1.85 M_{SUN}$$

$$M_{*\min} = 6(3\pi)^{1/2} \left(\frac{4}{5}\right)^{3/4} \left(\frac{m_P}{m_{ion}}\right)^2 \left(\frac{kT_{nuc}}{m_e c^2}\right)^{3/4} M_0 \approx 0.07 M_0$$

$$M_{*\max} = \left(\frac{18\sqrt{5}}{\pi^{3/2}}\right)^{3/4} \left(\frac{1-f_g}{f_g^4}\right)^{1/2} \left(\frac{m_P}{\langle m \rangle}\right)^2 M_0 \approx 56M_0$$

Chandrasekhar: $M_{Ch} = \frac{1}{5} (2\pi)^{3/2} (Z/A)^2 M_0 \approx 0.76 M_0 \approx 1.4 M_{SUN}$

Combined Constraint

We have required central temperature as function of stellar mass, and the relation between minimum stellar mass and given nuke-burning temperature. Combine to eliminate mass, get constraint:

$$\Theta_C I(\Theta_C) = \left(\frac{2^{20}\pi^4 3^4}{5^{11}}\right) \left(\frac{h^3}{c^2}\right) \left(\frac{1}{\beta\mu_0^4}\right) \left(\frac{1}{m m_e^3}\right) \left(\frac{G}{\kappa_0 C}\right)$$

$$I(\Theta_C) = \int_{0}^{\xi_*} f^{2n} \xi^2 \Theta^2 \exp[-3\Theta] d\xi \quad \& \quad \Theta = \left(\frac{E_G}{4kT}\right)^{1/3}$$

Central Temperature Solution



Allowed Region of Parameter Space



CONCLUSION

Stars in other universes are not as rare as sometimes claimed: Substantial fraction of parameter space allows for the existence of working stars.

This result is only the first step:

Habitable universes require more constraints
Expand parameter space (particle masses)
Need the probability of realizing parameters
Consider more complicated stellar models
Some universes could have alternative stars, including dark matter stars and black holes

Region of Parameter Space where Black Holes act as Stars



Instead of having different values in other universes, the fundamental constants can have time-varying values within our own universe

Fate of Degenerate Objects



White Dwarf Evolution with Proton Decay and Time-varying G



$$G(t) = G_0 (1 + t/t_*)^{-p}$$

three regimes:

$$t_* >> \Gamma^{-1} (old \ track)$$

 $t_* << \Gamma^{-1} (horz \ track)$

 $t_* \approx \Gamma^{-1}$ (new track)

here : p = 0, 0.25, 0.5

Ketchum & Adams 2008

Uncertainties

As one journeys deeper into future time, projections necessarily become more uncertain (our basic timeline stops at cosmological decade n = 100).
As we learn more about the fundamental laws of physics, or if the laws change with cosmological time, corrections (both

large and small) to this timeline must be made.

Higher Order Proton Decay

 If baryon decay is suppressed at leading order, it can still (sometimes) take place at higher order: n = 100 – 160

• The vacuum state of the standard model of particle physics allows tunneling transitions – sphalerons – that have timescale: n = 141



NOTE: Higher order proton decay is hard to observe

 The universe contains `only' about 10^78 protons in total. If the proton decay timescale is n = 100, e.g., then the probability that a single proton has decayed in the whole universe, thus far in time, is only about 0.00000000001

Liquid Rocks

 Since atoms within a lattice can tunnel via quantum mechanics, they can change places, so that any solid is really a liquid over the long term. Time required for rock to be a liquid:



$$t = 10^{65} yr$$
$$n = 65$$

Infinite Monkey Theorem

 Given enough time, a hypothetical monkey typing at random would, as part of his output, almost surely, produce all of the plays of Shakespeare



$$t = 10^{500} yr$$
$$n = 500$$

The Iron Age

 Because iron nuclei have the highest binding energy per particle, all nuclei will decay to iron if you wait long enough:

$$t = 10^{1500} yr \rightarrow n = 1500$$



Tunneling to Black Holes

 In the absence of proton decay, white dwarfs would live forever in the absence of a lower energy state, i.e., a black hole state. But it takes a long time for a white dwarf to tunnel into a black hole:

$$t = 10^{10^{76}} yr \rightarrow n = 10^{76}$$

A Basic Lesson

 The time scales of the last three processes are much longer than the time scales considered thus far.

 None of the last three processes are likely to happen – ever. Another process, in this case proton decay, will occur first (most likely) and thereby prevent them from happening.

Vacuum Tunneling Time Scales

φ

$$P = K \exp[-S_4]$$

$$V(\phi) = \lambda \phi^4 - a\phi^3 + b\phi^2 + c\phi + d$$

$$need \ S_4 \ge 231 \ \log 10 \approx 532 \ (survival)$$

$$for \ \lambda = 0.1 - 1:$$

$$0.5 < S_4 < 30,000$$

$$t_{max} \approx 10^{12,800} yr$$

$$n \approx 12,800$$

Wrapping Up

 The future evolution of the universe is a rich subject – many many topics Studies of the future universe provide us with a deeper understanding of the present-day universe As our theories get farther from the experimental constraints: Be careful!

 We have introduced the concept of a Copernican Time Principle. But: The cosmos does indeed evolve with time, so that life – as we know it – is certainly more likely to arise in a particular range of epochs (maybe n = 6 - 30). To what extent is such a principle valid or useful?

 What happens to this (or any) projection of the future history of the universe if the laws of physics are time dependent?

 Experiments => we are good up to cosmological decade n = 15 or so; but there is a lot of real estate at later times

 Given the uncertainties inherent in studying the future universe, at what point does the exercise stop being science?

 Similarly, as we go back in time toward the moment of the Big Bang, at what point does that exercise stop being science?

 (in both cases, one moves farther away from experimental confirmation/refutation)

 There is an asymmetry between the past history and future history of the universe. We can look for signatures of past events but not for future events. To what extent does this asymmetry affect the answers to the previous questions?