RESONANT STRIPPING AS THE ORIGIN OF DWARF SPHEROIDAL GALAXIES

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Dwarf spheroidals (dSphs) challenge our understanding of galaxy formation and evolution because:

- dSphs are gas poor and have few stars (Mateo 1998)
- found in galaxy clusters and groups (Fergusson & Binggeli 1994)
- the most dark matter dominated galaxies (M/L~ 30-100)
- the ultra-faint dwarf galaxies have L~ 1000 Lsun (Willman et al. 2005, Zucker et al. 2006; Belokurov et al. 2009; Walker et al. 2008)
We need a mechanism to separate gas & stars from DM to explain high M/L ratio

✨ Previous Theories:

✓ Gas photoheated during reionization or blown out by feedback, but:
  • few signatures of reionization in dSphs (Gallagher et al. 2003)
  • difficulties explaining the morphology

✓ Tidal shocking can convert a disk of stars into a spheroid but requires:
  • \textbf{ram pressure} to remove the gas (Mayer et al. 2007)
  • that dwarfs orbit close to Milky Way or Andromeda
A small dwarf orbiting inside a larger system
(D’Onghia et al., 2009, Nature, 460, 605)

The stripping of stars is caused by a gravitational process:
“Resonant Stripping”: stars and gas in the victim are removed by a resonance between the spin frequency of its disk and the angular frequency of its orbit around the perturber.
Surprising outcome: baryons at the bottom of the potential well are removed!

NEW: “Resonant Stripping” alters the M/L ratio in galaxies because stars and gas are removed more efficiently than the dark matter
Resonant Stripping

Tidal Stripping

Stars

t=0

t=2 Gyrs

t=7 Gyrs

Resonant Stripping

Tidal Stripping

\[ M_{\text{dark}} / M_{\text{star}} \]

\[ t \text{ [billion years]} \]
Nearly Prograde Encounter

Nearly Retrograde Encounter
Evolution of mass surface density profile

D’Onghia et al., 2009, Nature, 460, 605
Evolution of kinematic and structural properties

D’Onghia et al., 2009, Nature, 460, 605
Predictions

✓ “Resonant Stripping” should be visible in situ in groups of dwarfs nearby.

✓ Many dSphs should be found in groups of dwarfs along with detectable stellar tails and shells.

 tadpole galaxy                  NGC 2782
NOTE: Resonant Stripping depends on a combination of the rotation curve and orbital parameters

$$\Omega_{dwarf} \approx \Omega_{Pert} \Rightarrow \frac{v}{r} \approx \frac{V}{R} \sqrt{1 + e}$$

Fainter Dwarfs may be reproduced assuming a slowly rising rotation curve
The tails and bridges of stars in major mergers are caused by a tidal resonance (Toomre & Toomre 1972)

Simulations have shown that 10% of stars are removed during major mergers
“Quasi-Resonance theory of tidal interactions”

\[ \Delta v(x) = -GM_{\text{pert}} \int \left[ \frac{x}{|X|^3} - \frac{X(x \cdot X)}{|X|^5} \right] dt \]

\[ \Delta v_x = -2 \frac{GM_{\text{pert}}}{b^2 V} r \cos \phi \left( -\alpha^2 K_0 - \alpha (1 \pm \alpha) K_1 \right) \]

\[ \Delta v_y = -2 \frac{GM_{\text{pert}}}{b^2 V} r \sin \phi \left( \alpha^2 K_0 \pm \alpha^2 K_1 \right) \]

\[ \alpha = \frac{|\Omega|}{V} \]

D’Onghia et al., 2010, in prep
Limits of "Tidal Quasi-Resonance approximation"

✓ For angular frequency=0

\[ \Delta v_x = -2 \frac{GM_{\text{pert}} b}{V} r \cos \phi_0 (-\alpha^2 K_0 - \alpha (1 \pm \alpha) K_1) \]
\[ \Delta v_y = -2 \frac{GM_{\text{pert}} b}{V} r \sin \phi_0 (\alpha^2 K_0 \pm \alpha^2 K_1) \]

\( \alpha = 0 \)

Tidal near-resonance \( \rightarrow \) Impulse Approximation

\[ \Delta v_x = -2 \frac{GM_{\text{pert}}}{b^2 V} r \cos \phi_0 \]
\[ \Delta v_y = -2 \frac{GM_{\text{pert}}}{b^2 V} r \sin \phi_0 \]

✓ For Infinite angular frequency

\( \alpha = \infty \)

Tidal near-resonance \( \rightarrow \) Adiabatic Invariant

\[ \Delta v_x = -2 \frac{GM_{\text{pert}} b}{V} r \cos \phi_0 (-\alpha^2 K_0 - \alpha (1 \pm \alpha) K_1) \]
\[ \Delta v_y = -2 \frac{GM_{\text{pert}} b}{V} r \sin \phi_0 (\alpha^2 K_0 \pm \alpha^2 K_1) \]

\[ \Delta v_x = -2 \frac{GM_{\text{pert}}}{b^2 V} xae^{-\alpha} \frac{\pi}{\sqrt{2\alpha}} (1 - \alpha \pm \alpha) \rightarrow 0 \]
\[ \Delta v_y = 2 \frac{GM_{\text{pert}}}{b^2 V} yae^{-\alpha} \frac{\pi}{\sqrt{2\alpha}} (1 - \alpha) \rightarrow 0 \]
“Tidal Quasi-Resonance approximation”

\[
\Delta v_x = -4 \frac{GM_{\text{pert}}}{b^2 V} r \cos \phi_0 \left( 2I_{30} - 2(1 \pm 2\alpha)I_{20} \pm 3\alpha I_{10} - 4\alpha^2 I_{00} \right)
\]

\[
\Delta v_y = 4 \frac{GM_{\text{pert}}}{b^2 V} r \sin \phi_0 \left( 2I_{30} - (1 \pm 4\alpha)I_{20} \pm 3\alpha I_{10} - 4\alpha^2 I_{00} \right)
\]

\[
\alpha = \frac{|\Omega|b}{V}
\]
The Tidal Quasi-Resonance approximation is more efficient than the impulse approximation.

The resonance is broad.

\[ \Omega_{\text{dwarf}} \approx \Omega_{\text{Pert}} \Rightarrow \frac{v}{r} \approx \frac{V}{R} \sqrt{1 + e} \]

D'Onghia et al., 2010, in prep
Efficiency of Resonant Stripping

The efficiency depends on:

mass ratio \((m/M)\) and

impact parameter \(b\)
Comparison between the quasi-resonance theory and simulations

Prograde Orbit

The Tidal Resonance Approximation

Numerical Simulation

D’Onghia et al. 2010, in prep
The energy and angular momentum distributions match!
Retrograde Orbit

The Tidal Resonance Approximation

Numerical Simulation
Non Coplanar Orbit

The Tidal Resonance Approximation

Numerical Simulation
Parabolic Orbit

The Tidal Resonance Approximation

Numerical Simulation
The Magellanic Stream as the product of an LMC + SMC tidal encounter
Interactions between dwarf galaxies in small groups excite a resonant response: "Resonant Stripping" that rapidly transforms disks into dSphs.

Resonant stripping is a gravitational process that removes gas & stars in a disk but affects less DM and can be described by tidal Quasi-Resonance Theory.

TEST: -- rotational velocity of the stars in data

-- subhalo angular momentum in cosmological simulations and if they are preferentially retrograde