

Nuclear reactions in the early universe II

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Organization

Nuclear reactions in the early universe

- Lectures (Paris/E. Grohs)
 - I. Overview of cosmology/Kinetic theory/Big bang nucleosynthesis (**BBN**)
 - II. Scattering & reaction formalism/Neutrino energy transport
- Workshop sessions (E. Grohs/Paris)
 - I. BBN exercises: compute Nuclear Statistical Equilibrium/electron fraction
 - II. Compute primordial abundances vs $\Omega_b h^2$: code parallelization
- Lecture notes
 - Will be available online (URL TBA)

Outline

Lecture I

- Overview
- Cosmological dynamics in GR
- Big bang nucleosynthesis (BBN)
- Boltzmann equation
 - Flat & curved spacetime

Lecture II

- Unitary reaction network (URN) of light nuclei
- Neutrino energy transport
- Evan Grohs: observations of primordial abundances

Light nuclear reaction program @ LANL

□ Motivation

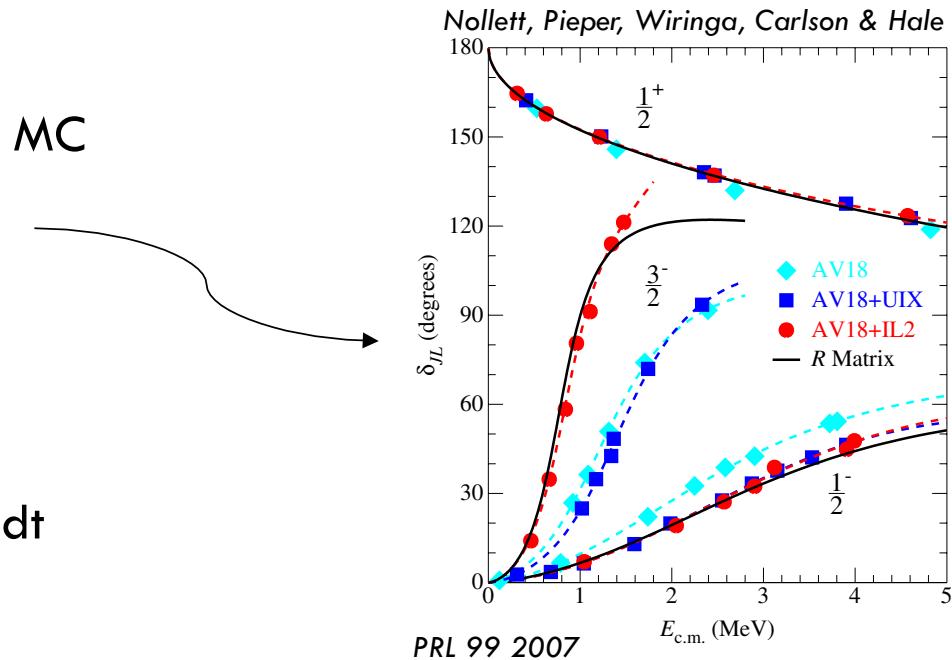
- Data sets: σ , $\sigma(\theta)$, $A_i(\theta)$, $C_{i,j}$, K_i^j , $\Sigma(\gamma)$, ... $\rightarrow T$ matrix \rightarrow resonance spectrum
- **Unitary** parametrization of compound nuclear system
- Applications: **astrophysical**, nuclear security, inertial confinement fusion, **criticality safety**, charge-particle transport, nuclear data (ENDF, ENSDF)

□ Ab initio

- Variational MC; Green's function MC
- GFMC [PRL 99, 022502 (2007)]
 - $n-{}^4He$ phase shifts
 - comparison GFMC/R-matrix
- challenge: multichannel
 - e.g. $n\alpha \rightarrow n\alpha$, $n\alpha \rightarrow dt$ & $dt \rightarrow dt$

□ Phenomenology

- R matrix (2 \rightarrow 2 body scatt/reacs)
- 3-body channels being incorporated



EDA Analyses of Light Systems

A	System	Channels	Energy Range (MeV)
2	N-N	p+p; n+p, γ +d	0-30 0-40
3	N-d	p+d; n+d	0-4
4	^4H ^4Li	n+t p+ ^3He	0-20
	^4He	p+t n+ ^3He d+d	0-11 0-10 0-10
5	^5He	n+ α d+t $^5\text{He}+\gamma$	0-28 0-10
	^5Li	p+ α d+ ^3He	0-24 0-1.4

Analyses of Light Systems, Cont.

A	System (Channels)
6	^6He ($^5\text{He} + \text{n}$, $\text{t} + \text{t}$); ^6Li ($\text{d} + ^4\text{He}$, $\text{t} + ^3\text{He}$); ^6Be ($^5\text{Li} + \text{p}$, $^3\text{He} + ^3\text{He}$)
7	^7Li ($\text{t} + ^4\text{He}$, $\text{n} + ^6\text{Li}$); ^7Be ($\gamma + ^7\text{Be}$, $^3\text{He} + ^4\text{He}$, $\text{p} + ^6\text{Li}$)
8	^8Be ($^4\text{He} + ^4\text{He}$, $\text{p} + ^7\text{Li}$, $\text{n} + ^7\text{Be}$, $\text{p} + ^7\text{Li}^*$, $\text{n} + ^7\text{Be}^*$, $\text{d} + ^6\text{Li}$)
9	^9Be ($^8\text{Be} + \text{n}$, $\text{d} + ^7\text{Li}$, $\text{t} + ^6\text{Li}$); ^9B ($\gamma + ^9\text{B}$, $^8\text{Be} + \text{p}$, $\text{d} + ^7\text{Be}$, $^3\text{He} + ^6\text{Li}$)
10	^{10}Be ($\text{n} + ^9\text{Be}$, $^6\text{He} + \alpha$, $^8\text{Be} + \text{nn}$, $\text{t} + ^7\text{Li}$); ^{10}B ($\alpha + ^6\text{Li}$, $\text{p} + ^9\text{Be}$, $^3\text{He} + ^7\text{Li}$)
11	^{11}B ($\alpha + ^7\text{Li}$, $\alpha + ^7\text{Li}^*$, $^8\text{Be} + \text{t}$, $\text{n} + ^{10}\text{B}$); ^{11}C ($\alpha + ^7\text{Be}$, $\text{p} + ^{10}\text{B}$)
12	^{12}C ($^8\text{Be} + \alpha$, $\text{p} + ^{11}\text{B}$)
13	^{13}C ($\text{n} + ^{12}\text{C}$, $\text{n} + ^{12}\text{C}^*$)
14	^{14}C ($\text{n} + ^{13}\text{C}$)
15	^{15}N ($\text{p} + ^{14}\text{C}$, $\text{n} + ^{14}\text{N}$, $\alpha + ^{11}\text{B}$)
16	^{16}O ($\gamma + ^{16}\text{O}$, $\alpha + ^{12}\text{C}$)
17	^{17}O ($\text{n} + ^{16}\text{O}$, $\alpha + ^{13}\text{C}$)
18	^{18}Ne ($\text{p} + ^{17}\text{F}$, $\text{p} + ^{17}\text{F}^*$, $\alpha + ^{14}\text{O}$)

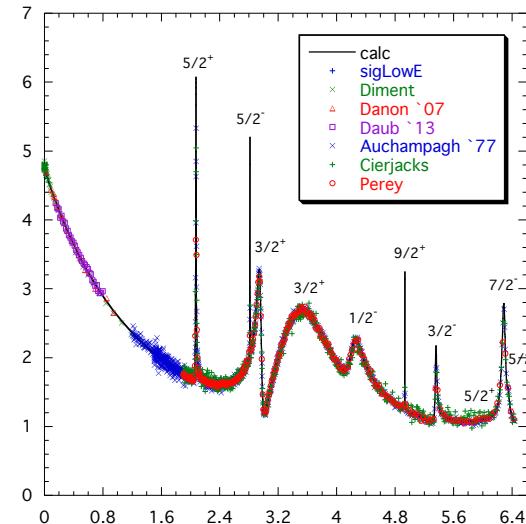
26 tabulated analyses

Paris BBN 2014 May 15



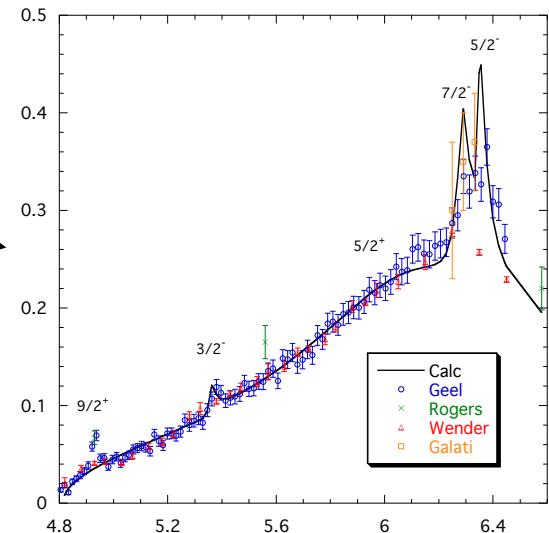
$^{13,14}\text{C}$ system analyses: σ_T (b) vs. E_n (MeV)

$n + ^{12}\text{C}$ Total Cross Section

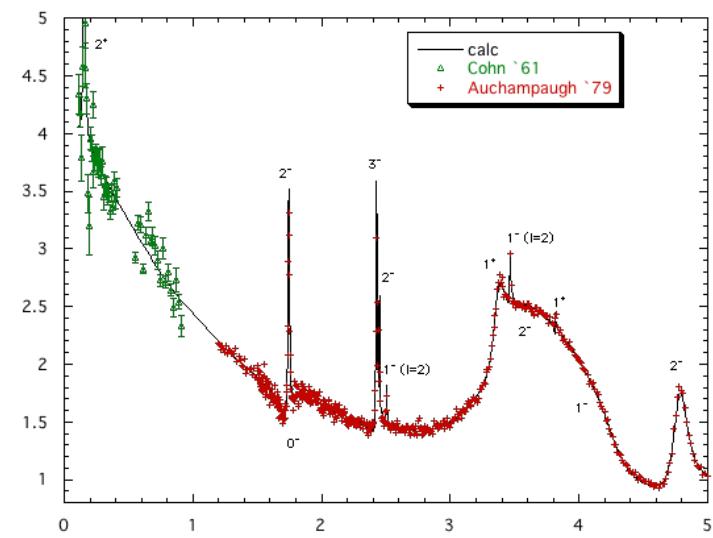


Two-channel analysis

$^{12}\text{C}(n,n')$ Cross Section



$n + ^{13}\text{C}$ Total Cross Section



Analyses by GMH/MWP

Single-channel analysis

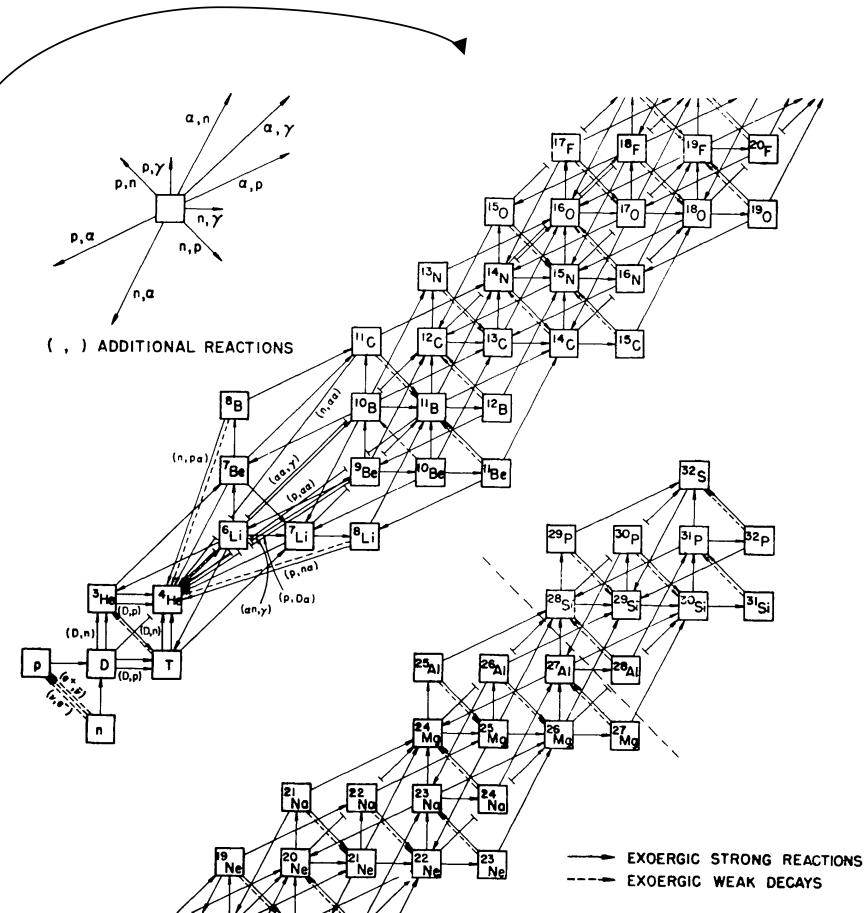
Unitary, self-consistent primordial nucleosynthesis

- State of standard big-bang nucleosynthesis (BBN)
 - d & ^4He abundances: signature success cosmology+nucl astro+astroparticle
 - but there's at least one **Lithium (^7Li) Problem** [^6Li too? See: [Lind et.al. 2013](#)]
 - coming precision observations of d, ^4He , η , N_{eff} demand new BBN capabilities
 - resolution of ^7Li problem:
 - observational/stellar astrophysics?
 - ^7Li controversial anomaly: nuclear physics solution?
 - new physics?
- Advance BBN as a tool for precision cosmology
 - incorporate **unitarity** into strong & electroweak interactions ([next slide](#))
 - couple **unitary reaction network (URN)** to full Boltzmann transport code
 - neutrino energy distribution function evolution/transport code
 - fully coupled to nuclear reaction network
 - calculate light primordial element abundance for non-standard BBN
 - active-sterile ν mixing
 - massive particle out-of-equilibrium decays → energetic active SM particles
 - Produce tools/codes for nuc-astro-particle community: test new physics w/BBN
 - existing codes are based on Wagoner's (1969) code

Nuclear reaction network

- Single-process (non-unitary) analysis
 - $\sigma_{\alpha\beta}(E) \pm \delta \sigma_{\alpha\beta}(E)$ from expt
 - fit form (non-res+narrow res) to $\sigma_{\alpha\beta}(E)$
 - compute $\langle\sigma v\rangle(T) \rightarrow$ reactivity \rightarrow network
 - **NB:** norm. systematics can be large
 - ^{17}O case (below)

- Multi-channel (unitary) analysis
 - Construct unitary parametrization
 - R-matrix (Wigner-Eisenbud '47)
 - simultaneous fit of unpolarized/pol'd scatt/reac data → determine T (or S)matrix
 - determines a unitary reaction network (URN) for analyzed compound systems



Wagoner ApJ Suppl '69

Boltzmann eq., cross sections, thermal averages

□ Boltzmann equation

- Toy model, single reaction → $\frac{1}{a^3} \frac{d(n_1 a^3)}{dt} = -\langle \sigma v \rangle \left\{ n_1 n_2 - n_3 n_4 \frac{n_1^{(0)} n_2^{(0)}}{n_3^{(0)} n_4^{(0)}} \right\}$
 - Full code has 144 reactions

- Thermal (Maxwellian) averaged flux(v)*cross section

$$\langle \sigma v \rangle = \left(\frac{8}{\pi \mu} \right)^{1/2} \left(\frac{1}{kT} \right)^{3/2} \int_0^\infty dE E \sigma_{12 \rightarrow 34}(E) e^{-E/kT}$$

□ Energy dependent, angle-integrated cross section is determined from data; Ranking worst → best:

- Guess: sometimes necessary when no data/calc. (e.g. TALYS)
- Parametrize resonance data: undesirable since res/non-res related by unitarity; results in model dependent reaction cross section
- Fit to experimental cross section: can be OK; normalization often problematic; subject to sometimes large systematic uncertainty
- Unitary theory: multichannel R-matrix: sure-fire; downside: need multichannel data

Observables from transition (T) matrix

- Scattering matrix: QM amplitude for (i)initial \rightarrow (f)inal

$$\langle f | S(E) | i \rangle = \delta_{fi} + 2iT_{fi}(E)$$

- All observables \sim T matrix bilinears

- unpolarized differential cross section

$$\frac{d\sigma_{fi}}{d\Omega} = \frac{4\pi}{k^2} \frac{1}{N_{spins,i}} \sum_{spins,f} |T_{fi}|^2$$

- polarization asymmetry

$$P = \frac{\sigma_{\uparrow\uparrow} - \sigma_{\downarrow\uparrow}}{\sigma_{\uparrow\uparrow} + \sigma_{\downarrow\uparrow}}$$

- Diff cross section \rightarrow int'd cross section \rightarrow thermal averaged

$$\sigma(E) = \int d\Omega \frac{d\sigma}{d\Omega} \rightarrow \langle \sigma v \rangle$$

Unitarity: consequences on T matrix

$$\left. \begin{array}{l} \delta_{fi} = \sum_n S_{fn}^\dagger S_{ni} \\ S_{fi} = \delta_{fi} + 2i\rho_f T_{fi} \\ \rho_n = \delta(H_0 - E_n) \end{array} \right\}$$

$$T_{fi} - T_{fi}^\dagger = 2i \sum_n T_{fn}^\dagger \rho_n T_{ni}$$

NB: **unitarity** implies **optical theorem** $\sigma_{\text{tot}} = \frac{4\pi}{k} \text{Im } f(0)$; but **not only** the O.T.

■ Implications of **unitarity constraint on transition matrix**

1. Doesn't uniquely determine T_{ij} ; highly restrictive, however

Elastic: $\text{Im } T_{11}^{-1} = -\rho_1$ (assuming T & P invariance)

Multichannel: $\text{Im } \mathbf{T}^{-1} = -\boldsymbol{\rho}$

2. Unitarity violating transformations

- cannot scale **any** set: $T_{ij} \rightarrow \alpha_{ij} T_{ij} \quad \alpha_{ij} \in \mathbb{R}$

- cannot rotate **any** set: $T_{ij} \rightarrow e^{i\theta_{ij}} T_{ij} \quad \theta_{ij} \in \mathbb{R}$

★ consequence of linear 'LHS' \propto quadratic 'RHS'

Most important feature:
linear \sim quadratic

3. Unitary parametrizations constrain the experimental data itself

★ *normalization*, in particular

★ case studies: ^{17}O & ^9B compound system

Basics of R-matrix (data \rightarrow amplitudes)

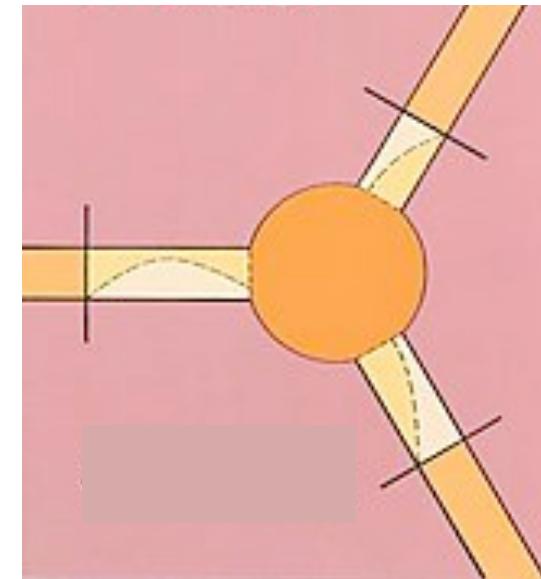
- Assumptions (cf. Lane & Thomas RMP '58)
 - a) Non-relativistic QM (L&T58); LANL-EDA uses rel.
 - b) Two-body channels only ('c'); aux. spectra code
 - c) Conservation of N, Z
 - d) Finite radius a_c beyond $V_{\text{pol}} \approx 0$; sharp boundaries

□ Separated pairs, “channels”

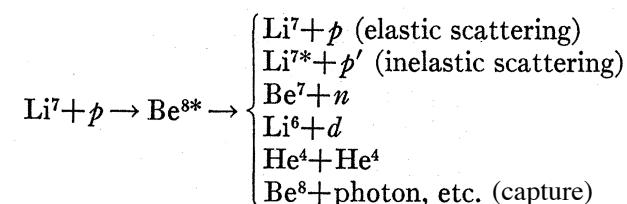
- A nucleons $\rightarrow (A_1, A_2)$
- $c = \{\alpha s_1 m_1 s_2 m_2\} \rightarrow \{\alpha(s_1 s_2) sm_s \ell m_\ell\} \rightarrow \{\alpha(s_1 s_2) s\ell, JM\}$
- Assume $a_c = a_\alpha \rightarrow$ many c have same channel in configuration space

□ Channel surface

- Consider configuration space of $3A$ dimensions
- Set of points: $\cup_c r_{\alpha(c)} = a_{\alpha(c)}$
- Surfaces coincide but assumed to have negl. prob.
- Channels are cylinders normal to channel surf.

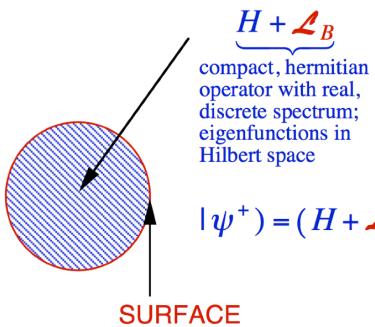


Example: ${}^8\text{Be}$ compound system



R-matrix formalism

INTERIOR (Many-Body) REGION
(Microscopic Calculations)



$$(r_{c'}|\psi_c^+\rangle = -I_{c'}(r_{c'})\delta_{c'c} + O_{c'}(r_{c'})S_{c'c}$$

$$|\psi^+\rangle = (H + \mathcal{L}_B - E)^{-1} \mathcal{L}_B |\psi^+\rangle$$

$$\mathcal{L}_B = \sum_c |c\rangle (d\left(\frac{\partial}{\partial r_c} r_c - B_c\right)),$$

$$(\mathbf{r}_c|c\rangle = \frac{\hbar}{\sqrt{2\mu_c a_c}} \frac{\delta(r_c - a_c)}{r_c} [(\phi_{s_1}^{\mu_1} \otimes \phi_{s_2}^{\mu_2})_s^\mu \otimes Y_l^m(\hat{\mathbf{r}}_c)]_J^M$$

$$R_{c'c} = (c'|(H + \mathcal{L}_B - E)^{-1}|c) = \sum_\lambda \frac{(c'|\lambda)(\lambda|c)}{E_\lambda - E}$$

Bloch operator $\mathcal{L}_B = \sum_c |c\rangle (c| \left[\frac{\partial}{\partial r_c} r_c - B_c \right])$ ensures Hermiticity of Hamiltonian restricted to internal region

Measurements

- R-matrix theory: **unitary**, multichannel parametrization of (not just resonance) data

Interior/Exterior regions

- Interior: strong interactions
- Exterior: Coulomb/non-polarizing interactions
- Channel surface

$$\mathcal{S}_c : r_c = a_c \quad \mathcal{S} = \sum_c \mathcal{S}_c$$

R-matrix elements

- Projections on channel surface functions $(\mathbf{r}_c|c)$ of Green's function

$$G_B = [H + \mathcal{L}_B - E]^{-1}$$

Boundary conditions

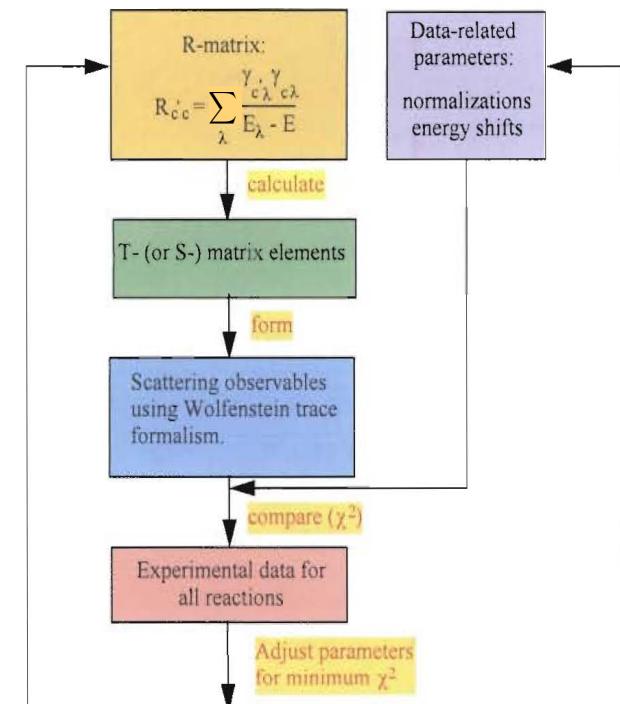
$$B_c = \frac{1}{u_c(a_c)} \frac{du_c}{dr_c} \Big|_{r_c=a_c}$$

R-matrix implementation in EDA

- EDA = Energy Dependent Analysis
 - Adjust E_λ & $\gamma_{c\lambda}$
- Any number of two-body channels
 - Arbitrary spins, masses, charges (zero mass)
- Scattering observables
 - Wolfenstein trace formalism
- Data
 - Normalization
 - Energy shifts
 - Energy resolution/spread
- Fit (rank-1 var. metric) solution

$$\chi^2_{EDA} = \sum_i \left[\frac{nX_i(\mathbf{p}) - R_i}{\delta R_i} \right]^2 + \left[\frac{nS - 1}{\delta S/S} \right]^2$$

- Covariance determined



^{17}O analysis configuration

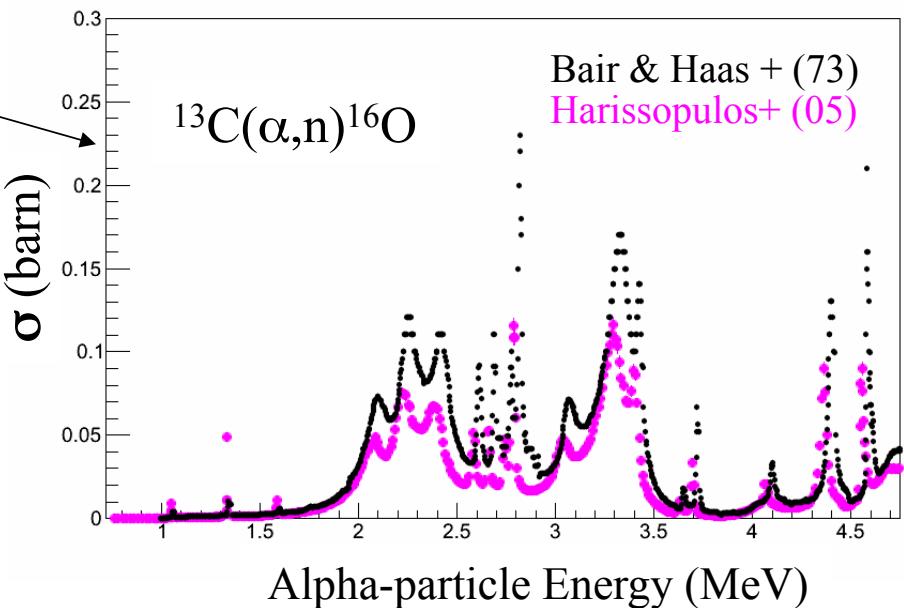
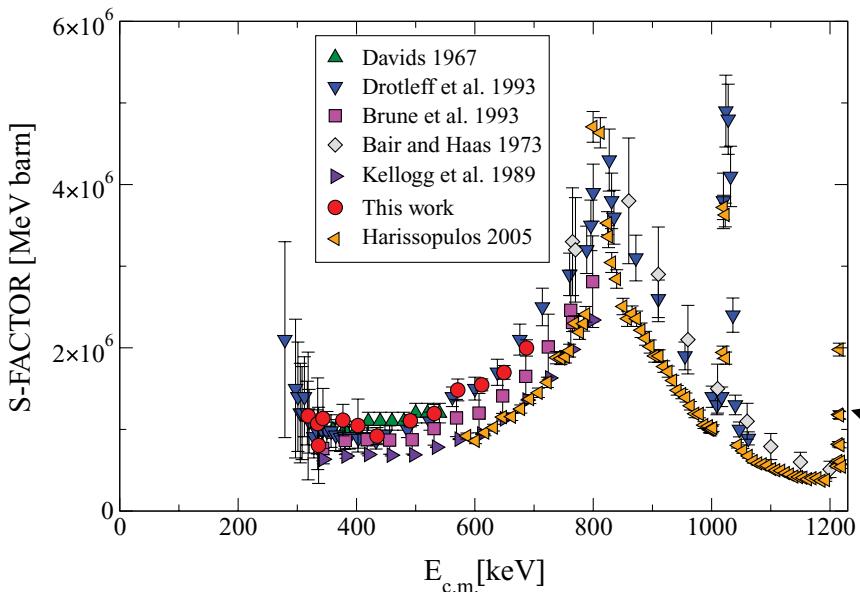
Channel	a_c (fm)	I_{\max}
$n + ^{16}\text{O}$	4.3	4
$\alpha + ^{13}\text{C}$	5.4	5

Reaction	Energies (MeV)	# data points	Data types
$^{16}\text{O}(n,n)^{16}\text{O}$	$E_n = 0 - 7$	2718	$\sigma_T, \sigma(\theta), P_n(\theta)$
$^{16}\text{O}(n,\alpha)^{13}\text{C}$	$E_n = 2.35 - 5$	850	$\sigma_{\text{int}}, \sigma(\theta), A_n(\theta)$
$^{13}\text{C}(\alpha,n)^{16}\text{O}$	$E_\alpha = 0 - 5.4$	874	σ_{int}
$^{13}\text{C}(\alpha,\alpha)^{13}\text{C}$	$E_\alpha = 2 - 5.7$	1296	$\sigma(\theta)$
total		5738	8

^{17}O compound system: experimental status

Recent (Harissopoulos '05) measurement
 $^{13}\text{C}(\alpha, n)^{16}\text{O}$ vs. older (Bair & Haas '73)

Heil et.al. PRC **78** 025803 ('08)



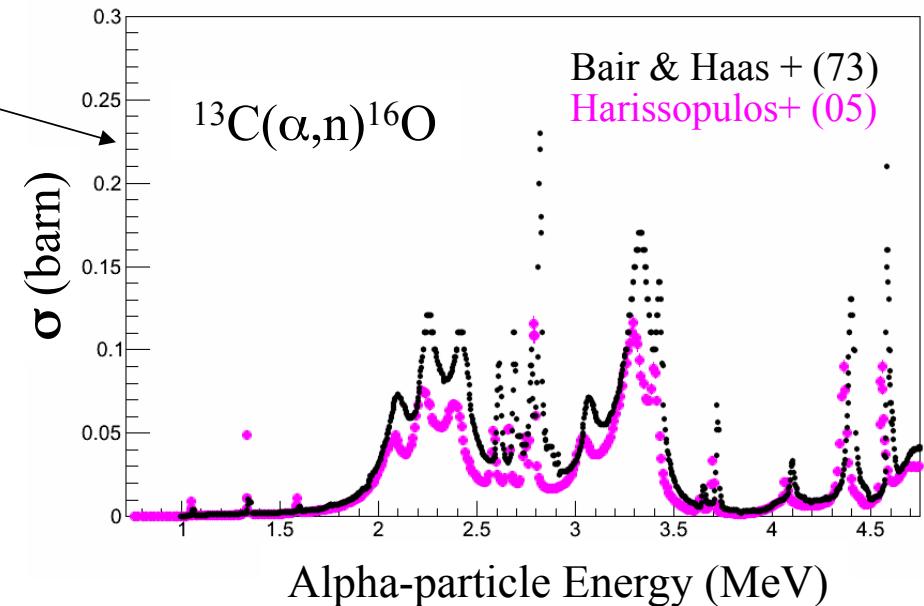
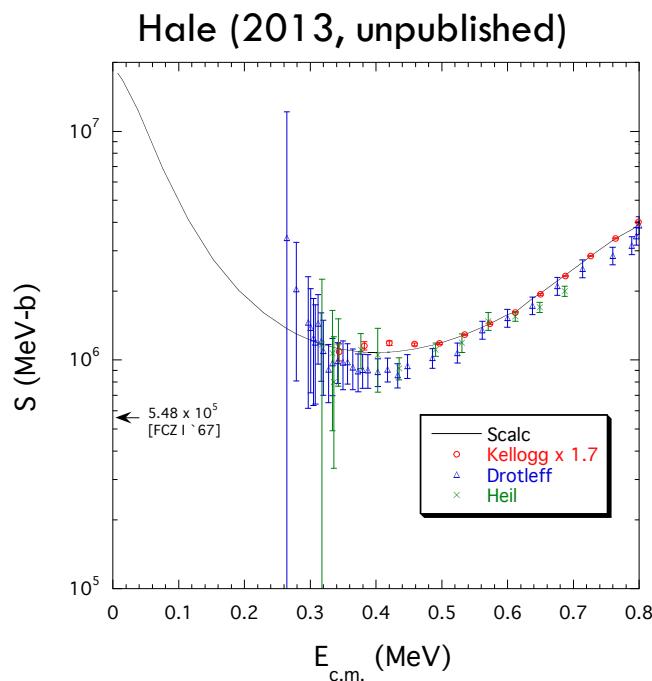
Credit: S. Kunieda

Harissopoulos(05) data 2/3*B&H(73)
Heil(08) data consistent with B&H

Tempting to conclude that B&H73 was right all along!

^{17}O compound system: experimental status

Recent (Harissopoulos '05) measurement
 $^{13}\text{C}(\alpha, n)^{16}\text{O}$ vs. older (Bair & Haas '73)



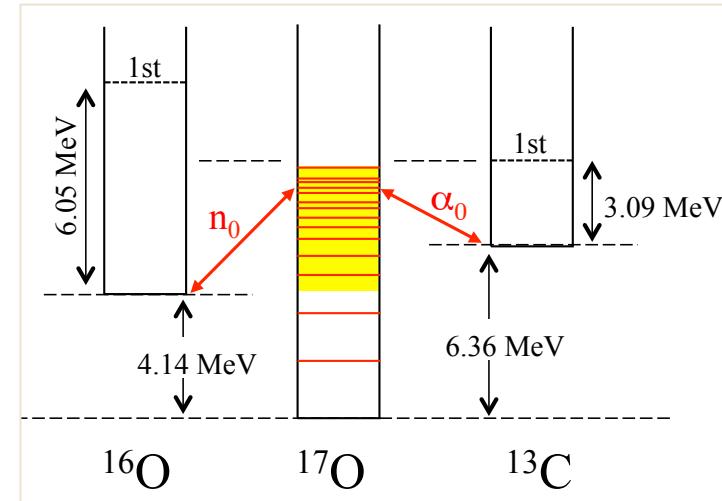
Credit: S. Kunieda

Subthreshold $1/2^+$
deep min in σ_T
 $S(0) \gg S_{\text{FCZ67}}(0)$

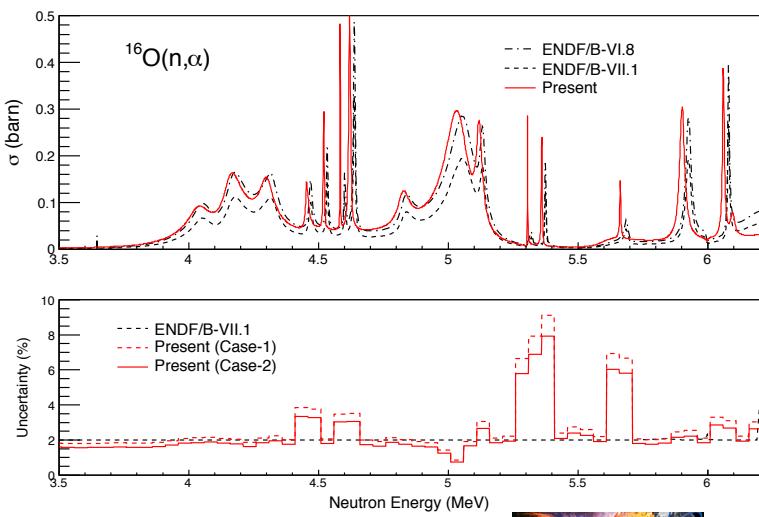
Tempting to conclude that B&H73 was right all along!

R-matrix analyses support B&H73/Heil08

- LANL R-matrix fit to Bair&Haas73
 - two-channel fit: $(^{16}\text{O}, n)$ & $(^{13}\text{C}, \alpha)$
 - $\ell_n = 0, \dots, 4$; $\ell_\alpha = 0, \dots, 5$
 - data included: $\sigma_T(E)$
 - $^{16}\text{O}(n,n)$, $^{16}\text{O}(n,\alpha)$, $^{13}\text{C}(\alpha,n)$
 - σ_{el} , $d\sigma/d\Omega$, A_y
 - χ^2 min: normalizations float
 - Test Hariss05 data
 - remove B&H73/Heil08 data
 - fix Hariss05 norm to unity
 - unable to obtain fit $\chi^2 (< 2.0)$
 - now allow Hariss05 norm to float
 - requires scale factor of ~ 1.5 , consistent with B&H73
- Kunieda/Kawano analysis [2013]
 - cf. LANL R-matrix(EDA)/ENDF/B-VI.8
 - with independent R-matrix code
 - **Right to conclude B&H73 data correct on the basis of unitarity!**



Credit: S. Kunieda



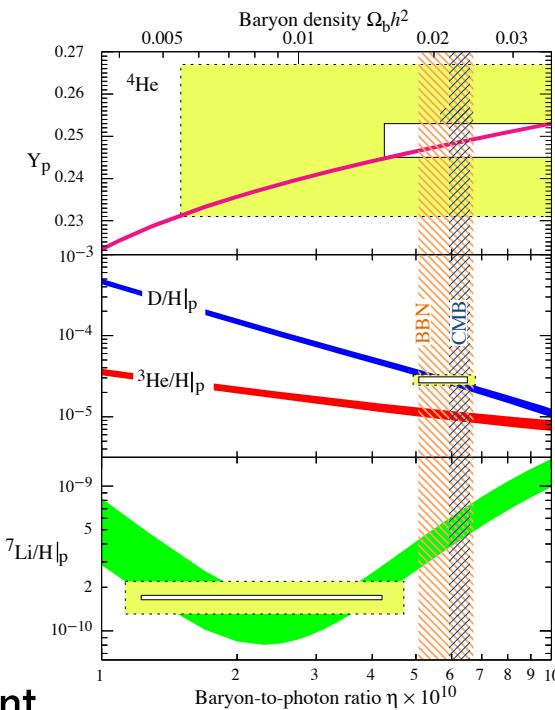
Toward a unitary reaction network for BBN

□ Primordial nucleosynthesis

- Can unitarity play a role in precision BBN?
- D,⁴He abund. agree with theo/expl uncertainties
- At η_{wmap} (CMB) ${}^7\text{Li}/\text{H}|_{\text{BBN}} \sim (2.2\text{--}4.2) * {}^7\text{Li}/\text{H}|_{\text{halo}}$ *
- Discrepancy $\sim 4.5\text{--}5.5 \sigma \rightarrow$ the “Li problem”

□ Resonant destruction ⁷Li

- Prod. mass ⁷Li “well understood”; destruction not
- Cyburt & Pospelov [arXiv:0906.4373](https://arxiv.org/abs/0906.4373); *IJMPE*, 21(2012)
 - ⁷Be(d,p) α α & ⁷Be(d, γ)⁹B resonant enhancement
 - Identify ⁹B $E_{5/2+} \approx 16.7 \text{ MeV} \approx E_{\text{thr}}(\text{d} + {}^7\text{Be}) + 200 \text{ keV}$
 - Near threshold
 - $(E_r, \Gamma_d) \approx (170\text{--}220, 10\text{--}40) \text{ keV}$ solve Li problem
- ‘Large’ widths
 - Conclude “large channel radius” required



NB: both approaches assume validity of TUNL-NDG tables

^9B analysis: included data

- $^6\text{Li} + ^3\text{He}$ elastic *Buzhinski et.al., Izv. Rossiiskoi Akademii Nauk, Ser.Fiz., Vol.43, p.158 (1979)*
 - Differential cross section
 - $1.30 \text{ MeV} < E(^3\text{He}) < 1.97 \text{ MeV}$
- $^6\text{Li} + ^3\text{He} \rightarrow p + ^8\text{Be}^*$ *Elwyn et.al., Phys. Rev. C 22, 1406 (1980)*
 - Integrated cross section
 - Quasi-two-body, excited-state, summed final channel
 - $0.66 \text{ MeV} < E(^3\text{He}) < 5.00 \text{ MeV}$
- $^6\text{Li} + ^3\text{He} \rightarrow d + ^7\text{Be}$ *D.W. Barr & J.S. Gilmore, unpublished (1965)*
 - Integrated cross section
 - $0.42 \text{ MeV} < E(^3\text{He}) < 4.94 \text{ MeV}$
- $^6\text{Li} + ^3\text{He} \rightarrow \gamma + ^9\text{B}$ *Aleksic & Popic, Fizika 10, 273-278 (1978)*
 - Integrated cross section
 - $0.7 \text{ MeV} < E(^3\text{He}) < 0.825 \text{ MeV}$
 - New to ^9B analysis
- New evaluation
 - Separate $^8\text{Be}^*$ states
 - $2^+@200 \text{ keV [16.9 MeV]}$, $1^+@650 \text{ keV [17.6 MeV]}$, $1^+@1.1 \text{ MeV [18.2 MeV]}$
 - $n + ^8\text{B}: E_{\text{thresh}}(^3\text{He}) = 3 \text{ MeV}$
 - Simultaneous analysis with ^9Be mirror system

Data accessed via
EXFOR/CSISRS
database (C4 format)

R-matrix configuration in EDA code

Hadronic channels (in blue, not included)

$A_1 A_2 \pi$	${}^3\text{He} {}^6\text{Li}^+(1)$		$p {}^8\text{Be}^{*+}(2)$		$d {}^7\text{Be}^-(3)$			
ℓ	S	$\frac{3}{2}$	$\frac{1}{2}$	$\frac{5}{2}$	$\frac{3}{2}$	$\frac{5}{2}$	$\frac{3}{2}$	$\frac{1}{2}$
0		${}^4S_{3/2}$	${}^2S_{1/2}$		${}^6S_{5/2}$	${}^4S_{3/2}$		
1		${}^4P_{5/2,3/2,1/2}$	${}^2P_{3/2,1/2}$		${}^6P_{7/2,5/2,3/2}$	${}^4P_{5/2,3/2,1/2}$	${}^6P_{7/2,5/2,3/2}$	${}^4P_{5/2,3/2,1/2}$
2		${}^4D_{7/2,5/2,3/2,1/2}$	${}^2D_{5/2,3/2}$	${}^6D_{9/2,7/2,5/2,3/2,1/2}$	${}^4D_{7/2,5/2,3/2,1/2}$	${}^6D_{9/2,7/2,5/2,3/2,1/2}$	${}^4D_{7/2,5/2,3/2,1/2}$	${}^2D_{5/2,3/2}$
		$E_{\text{thr}}(\text{CM, MeV})$		16.6	16.7	16.5		

Electromagnetic channel:

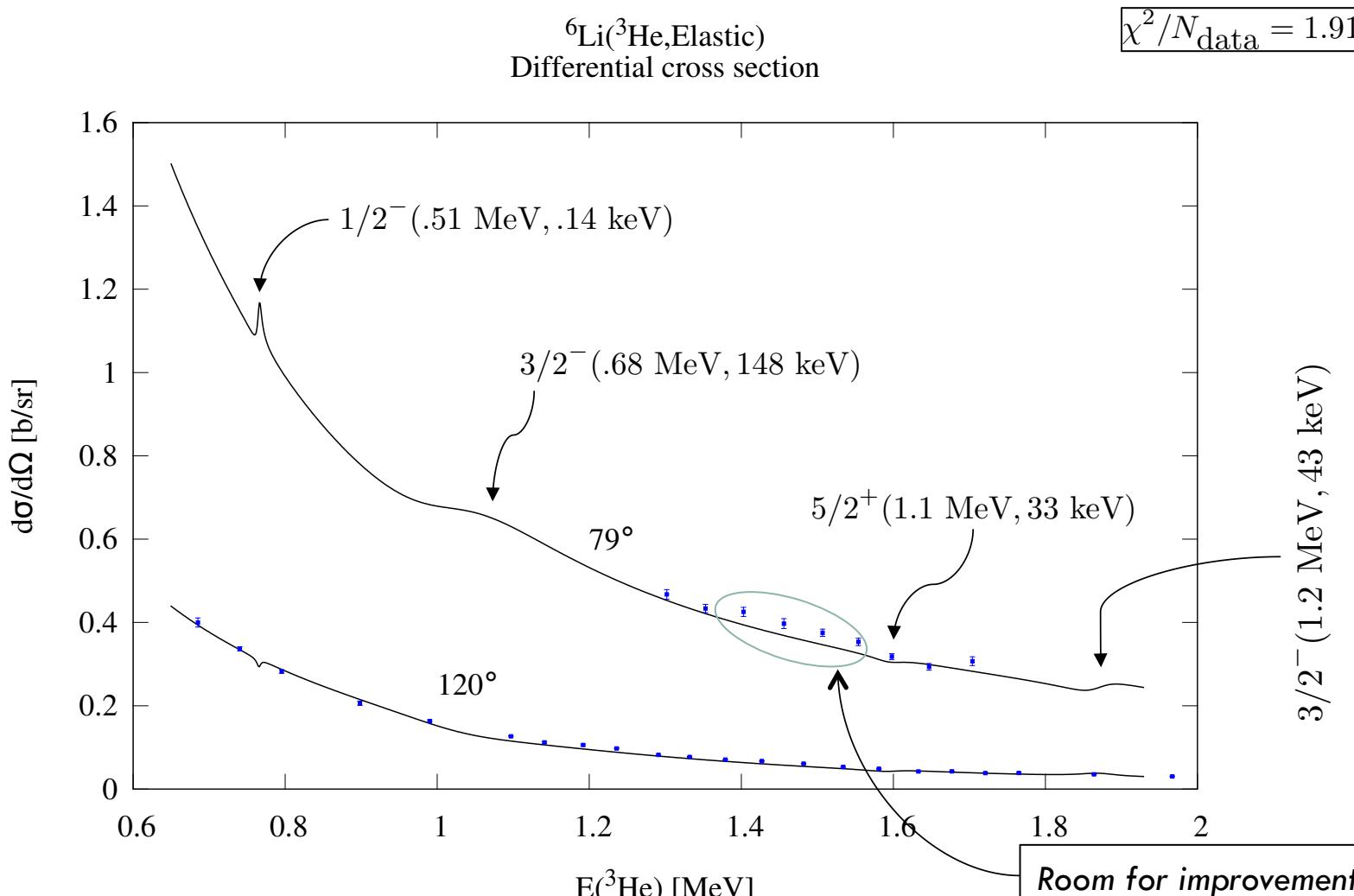
$$\gamma + {}^9B \rightarrow E_1^{3/2}, M_1^{5/2}, M_1^{3/2}, M_1^{1/2}, E_1^{5/2}, E_1^{1/2}$$

Full model space:
state number;
channel pair;
LS; J; channel
radius [fm]

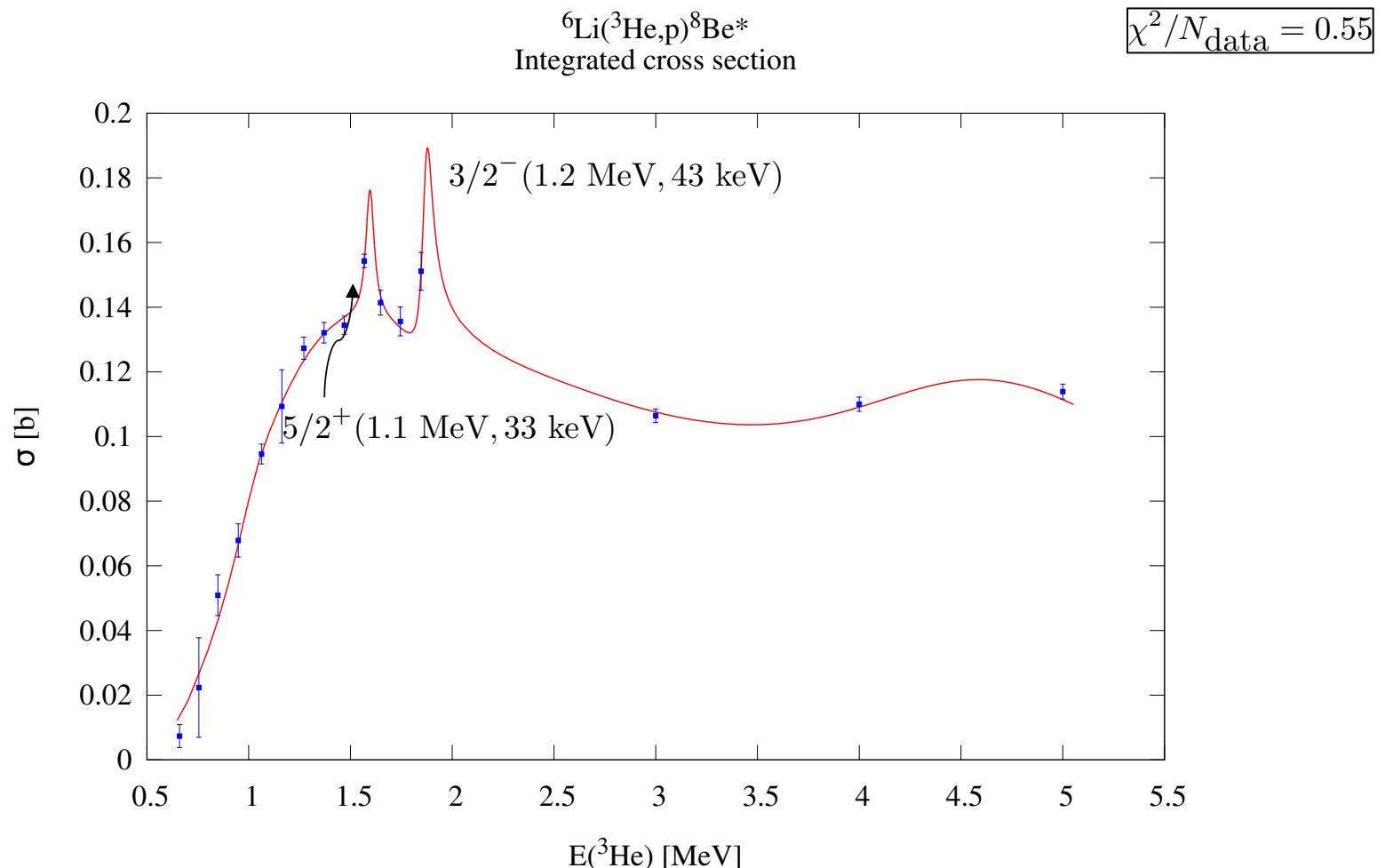
1	1	4s	3/2	7.50000000f	20	1	4p	1/2	7.50000000f
2	1	4d	3/2	7.50000000f	21	1	2p	1/2	7.50000000f
3	1	2d	3/2	7.50000000f	22	2	4p	1/2	5.50000000f
4	2	4s	3/2	5.50000000f	23	3	2s	1/2	7.00000000f
5	3	6p	3/2	7.00000000f	24	4	M1	1/2	50.00000000f
6	3	4p	3/2	7.00000000f	25	1	4d	7/2	7.50000000f
7	3	2p	3/2	7.00000000f	26	3	6p	7/2	7.00000000f
8	4	E1	3/2	50.00000000f	27	1	4d	5/2	7.50000000f
9	1	4p	5/2	7.50000000f	28	1	2d	5/2	7.50000000f
10	2	6p	5/2	5.50000000f	29	2	6s	5/2	5.50000000f
11	2	4p	5/2	5.50000000f	30	3	6p	5/2	7.00000000f
12	3	6s	5/2	7.00000000f	31	3	4p	5/2	7.00000000f
13	4	M1	5/2	50.00000000f	32	4	E1	5/2	50.00000000f
14	1	4p	3/2	7.50000000f	33	1	4d	1/2	7.50000000f
15	1	2p	3/2	7.50000000f	34	1	2s	1/2	7.50000000f
16	2	6p	3/2	5.50000000f	35	3	4p	1/2	7.00000000f
17	2	4p	3/2	5.50000000f	36	3	2p	1/2	7.00000000f
18	3	4s	3/2	7.00000000f	37	4	E1	1/2	50.00000000f
19	4	M1	3/2	50.00000000f	38	2	6p	7/2	5.50000000f



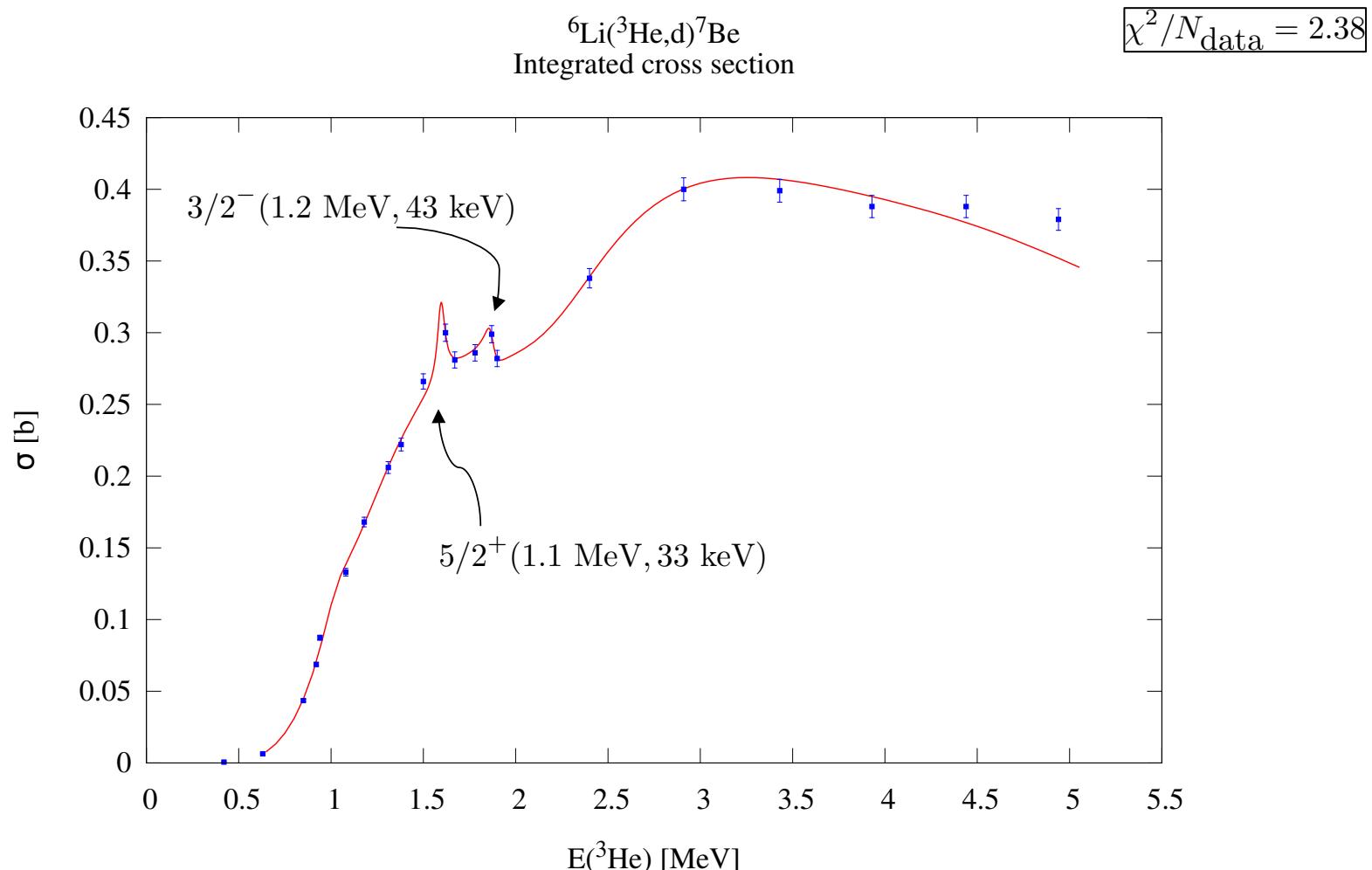
Observable fit: ${}^3\text{He} + {}^6\text{Li}$ elastic DCS



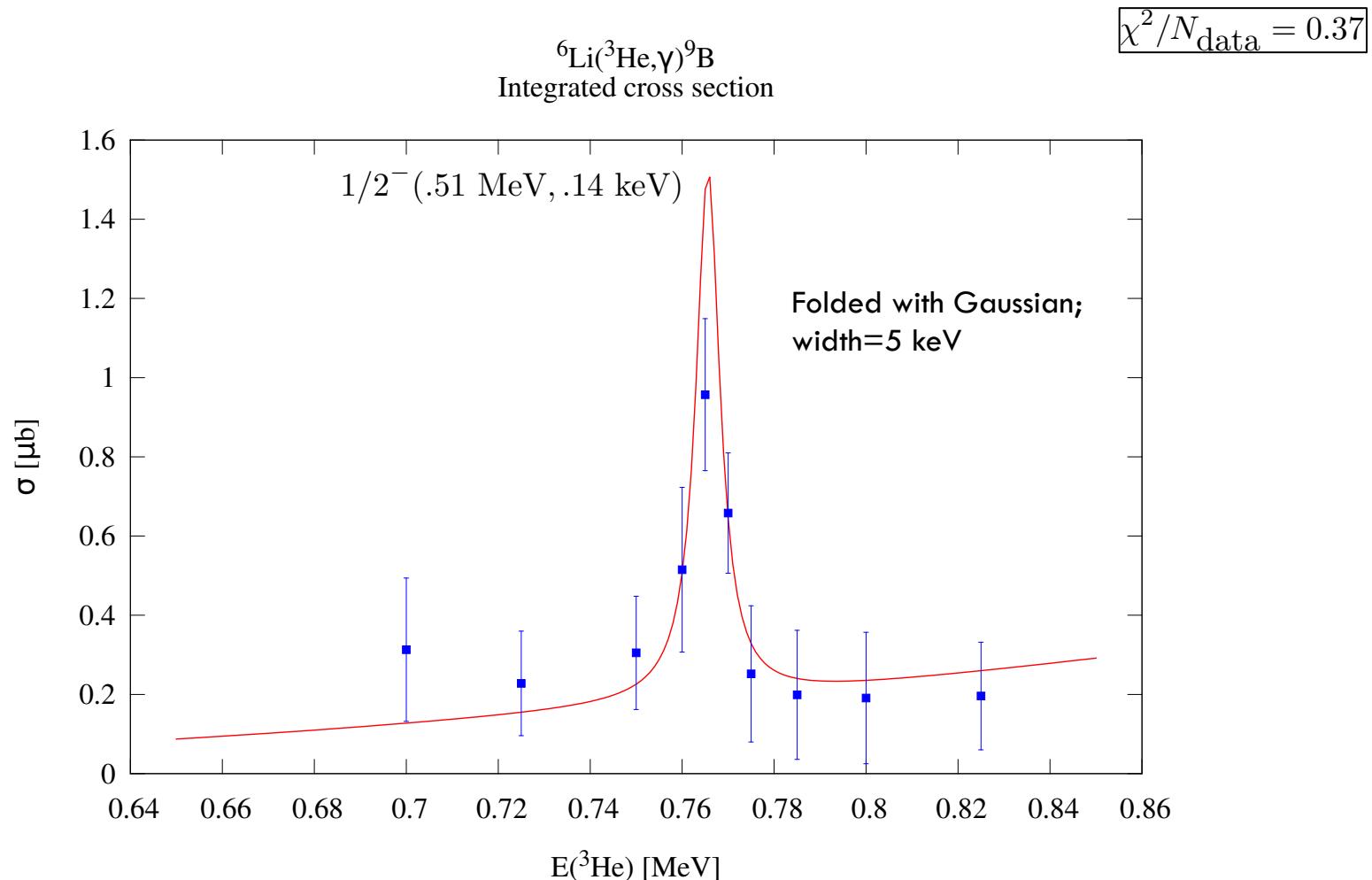
Observable fit: ${}^6\text{Li}({}^3\text{He},\text{p}){}^8\text{Be}^*$ integrated x-sec



Observable fit: ${}^6\text{Li}({}^3\text{He},\text{d}){}^7\text{Be}$ integrated x-sec



Observable fit: ${}^6\text{Li}({}^3\text{He}, \gamma) {}^9\text{B}$ integrated x-sec



⁹B analysis result: resonance structure

Ex(MeV)	Jpi	Gamma(keV)	Er(MeV)	ImEr(MeV)	E(³ He)	Strength
16.46539	1/2-	768.46	- .1369	-0.3842	-0.2054	0.06 weak
17.11317	1/2-	0.14	0.5109	-0.6771E-04	0.7664	1.00 strong
17.20115	5/2-	871.63	0.5989	-0.4358	0.8984	0.40 weak
17.28086	3/2-	147.78	0.6785	-0.0739	1.0178	0.77 strong
17.66538	5/2+	33.33	1.0631	-0.0167	1.5947	0.98 strong
17.83619	7/2+	2036.21	1.2339	-1.0181	1.8509	0.15 weak
17.84773	3/2-	42.52	1.2454	-0.0213	1.8681	0.97 strong
18.04821	3/2+	767.11	1.4459	-0.3836	2.1689	0.54 weak
18.42292	1/2+	5446.32	1.8206	-2.7232	2.7309	0.03 weak
18.67716	1/2-	10278.41	2.0749	-5.1392	3.1124	0.15 weak
19.60923	3/2-	1478.22	3.0069	-0.7391	4.5104	0.52 weak

TUNL-NDG/ENSDF
parameters

**NB: no strong resonance seen
~100 keV of ³He+⁶Li threshold**

E_x^a (MeV ± keV)	$J^\pi; T$	$\Gamma_{c.m.}$ (keV)	Decay
16.024 ± 25	$T = (\frac{1}{2})$	180 ± 16	
16.71 ± 100^h	$(\frac{5}{2}^+); (\frac{1}{2})$		
17.076 ± 4	$\frac{1}{2}^-; \frac{3}{2}$	22 ± 5	$(\gamma, {}^3\text{He})$
17.190 ± 25		120 ± 40	$p, d, {}^3\text{He}$
$17.54 \pm 100^{h,i}$	$(\frac{7}{2}^+); (\frac{1}{2})$		
17.637 ± 10^i		71 ± 8	$p, d, {}^3\text{He}, \alpha$

Summary

- Provided overview of current work in the LANL light nuclear reaction program
- Emphasize the utility of multichannel, unitary parametrization of light nuc data
 - ^{17}O norm issue: are Bair & Haas '73 data conclusive?
 - ^9B resonance spectrum:
 - no resonances in ^9B that reside within ~ 200 (~ 100) keV of the $d+^7\text{Be}$ ($^3\text{He}+^6\text{Li}$) threshold with ‘large’ widths 10—40 keV
 - Appears to rule out scenarios considered by Cyburt & Pospelov (2009) that low-lying, robust resonance in ^9B could explain the “Li problem”

End Lecture II

BSMs scenarios

- New particles: WIMPs, Axion, SUSY, ...
- GR modifications: new propagating DsOF; scalar-tensor
- Modifications of Cosmological SM: non-zero ν chem. pot.; non-equil. phenomena
- Variation of fundamental couplings
- Cosmic variance
- Neutrino sector
 - solar, atmospheric & reactor neutrinos oscillation experiment prove at least two neutrinos have mass
 - “sterile neutrinos”: mass \rightarrow neutrinos have left- & right-hand spin states
 - only left-hand neutrinos interact in SM
 - Massless neutrinos (recall)
 - have only one spin state

Neutrino Mass: what we know and don't know

We know the **mass-squared differences**:

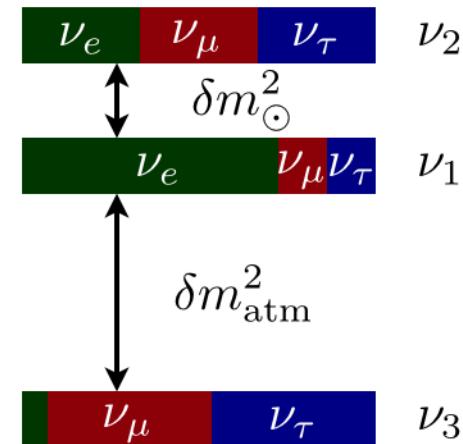
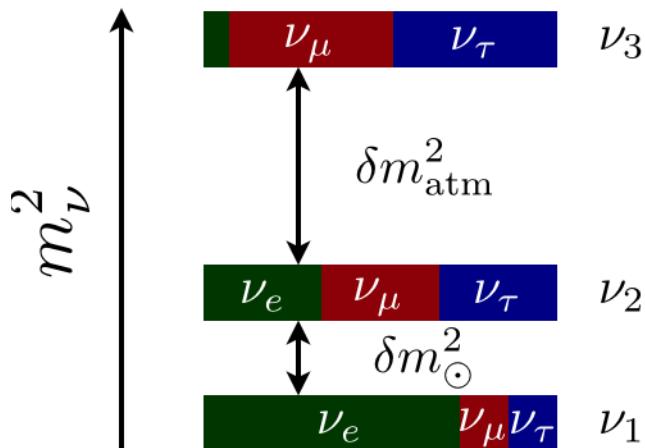
$$\left\{ \begin{array}{l} \delta m_{\odot}^2 \approx 7.6 \times 10^{-5} \text{ eV}^2 \\ \delta m_{\text{atm}}^2 \approx 2.4 \times 10^{-3} \text{ eV}^2 \end{array} \right.$$

$$\text{e.g., } \delta m_{21}^2 \equiv m_2^2 - m_1^2$$

We **do not** know the **absolute masses** or the **mass hierarchy**:

normal mass hierarchy

inverted mass hierarchy



Neutrino mass mixing 101

- Take-away message from experiments: “neutrinos have mass”

- neutrino flavor eigenstates

$$|\nu_e\rangle, |\nu_\mu\rangle, |\nu_\tau\rangle$$

- interact via left-hand (L) components

$$\bar{\psi}_e \gamma_\mu \frac{1}{2}(1 - \gamma_5) \psi_{\nu_e} = \bar{\psi}_{e,L} \gamma_\mu \psi_{e,L}$$

- Mass term, however, mixes L & R:

$$\bar{\psi}_e \psi_e = \bar{\psi}_{e,R} \psi_{e,L} + \bar{\psi}_{e,L} \psi_{e,R}$$

$$\begin{pmatrix} |\nu_e\rangle \\ |\nu_\mu\rangle \\ |\nu_\tau\rangle \end{pmatrix} = U_m \begin{pmatrix} |\nu_1\rangle \\ |\nu_2\rangle \\ |\nu_3\rangle \end{pmatrix} \quad U_m = U_{23} U_{13} U_{12} M$$

$$U_{23} \equiv \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos \theta_{23} & \sin \theta_{23} \\ 0 & -\sin \theta_{23} & \cos \theta_{23} \end{pmatrix}$$

$$U_{13} \equiv \begin{pmatrix} \cos \theta_{13} & 0 & e^{i\delta} \sin \theta_{13} \\ 0 & 1 & 0 \\ -e^{-i\delta} \sin \theta_{13} & 0 & \cos \theta_{13} \end{pmatrix}$$

$\theta_{12}, \theta_{23}, \theta_{13}, \delta$

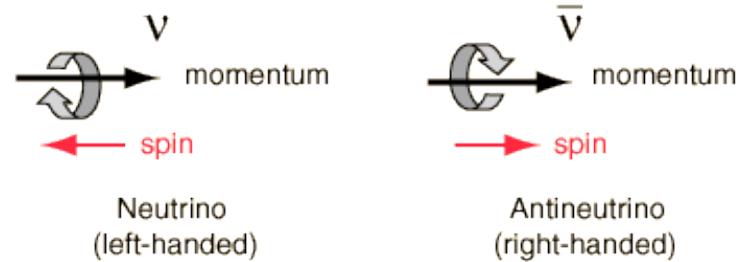
$$U_{12} \equiv \begin{pmatrix} \cos \theta_{12} & \sin \theta_{12} & 0 \\ -\sin \theta_{12} & \cos \theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

$$\begin{aligned} \theta_{12} &\approx 0.59^{+0.02}_{-0.015} \\ \theta_{23} &\approx 0.785^{+0.124}_{-0.124} \approx \frac{\pi}{4} \\ \theta_{13} &\approx 0.154^{+0.065}_{-0.065} \end{aligned}$$

$\delta = CP$ violating phase =?

Sterile* neutrinos

- What are they?
 - Related to right-handed components



- Wherefore?
 - Mass → right-handed neutrinos → must exist by Lorentz invariance
 - but may have mass modified by interactions
 - Non-interacting(?!): only example of particles that interact solely via GR
 - Interactions → necessarily beyond SM physics

- What (if anything) do they do?
 - perhaps they mix with active (e, μ, τ) neutrinos?
 - then they're not really "sterile"

$$|\nu_e\rangle = \cos\theta|\nu_1\rangle + \sin\theta|\nu_2\rangle$$
$$|\nu_s\rangle = -\sin\theta|\nu_1\rangle + \cos\theta|\nu_2\rangle$$

- Why would we want (need?) them?
 - leptogenesis; baryogenesis
 - BBN & N_{eff}

Hints for light sterile neutrinos?

- mini-BooNE
 - neutrino oscillation experiment $\nu_e \rightarrow \nu_s \rightarrow \nu_\mu$
 - appearance with $\delta m^2 \sim 1 \text{ eV}^2$
 - result inconsistent with flavor oscillation alone
- ~~Neutrino reactor anomaly~~
 - ~~3 σ deficit neutrinos detected in short-baseline (<100m) reactor ν experiments~~
 $\bar{\nu}_e$ deficit from $\bar{\nu}_e \rightarrow \bar{\nu}_s$ (???) – a disappearance experiment
 - A. Hayes et al. (2013) find “large corrections”
- Extra radiation at photon-decoupling (Neff) ??
 - CMB observations (PolarBear, ACT, SPT, Planck, CMBPol,...)
 - ‘extra’ RED could reconcile H_0 and σ_8 inferred from CMB and astronomical observation

Dark radiation

- γ -decoupling (last scattering) $T \sim 0.2$ eV ($z \sim 1000$)
- N_{eff} : “effective number of neutrino degrees of freedom”
 - A misnomer; it refers to any/all relativistic particles at decoupling
 - ‘Baby’ formula: $\rho_{\text{rad}} = 2 \left[1 + \frac{7}{8} \left(\frac{4}{11} \right)^{4/3} N_{\text{eff}} \right] \frac{\pi^2}{30} T_{\gamma}^4$
 - We’ve done this better...
- CSM+SMPP \rightarrow predicts $N_{\text{eff}} = 3.046$ [Dicus et. al. '83; Dolgov, Hansen, Semikoz '97, '99; Gnedin² '98,...]
 - annihilation of neutrinos-antineutrinos at weak decoupling
 - QED corrections
- Measurements
 - WMAP9 (2012): 3.26(35); Planck (2013): 3.30(50); ACT(2013): 2.79(56); SPT-SZ (2012): 3.71(35)
- Sterile neutrinos can affect the physics of dark radiation

CMB as a probe of steriles: caveats

- Sterile neutrinos can decay *out-of-equilibrium*
 - “dilution”: steriles are “sub-weakly” interacting
 - non-thermal energy spectra/number densities
- Care must be applied when
 - computing N_{eff} : non-equilibrium effects; relativistic vs. non-relativistic kinematics
 - determining N_{eff} and Y_p (mass fraction ${}^4\text{He}$)
 - current Planck collab. procedure is inconsistent w.r.t. N_{eff} and Y_p
 - in preparation: “**Neutrino physics in the era of precision cosmology**”
- neutron/proton ratio (and therefore ${}^4\text{He}$)
 - competing weak reaction rates determine $Y_p({}^4\text{He})$
 - **very sensitive to neutrino energy spectra**

$$\begin{aligned}\nu_e + n &\leftrightarrow p + e^- \\ \bar{\nu}_e + p &\leftrightarrow n + e^+ \\ n &\leftrightarrow p + e^- + \bar{\nu}_e\end{aligned}$$

Dilution physics (I)

- Consider the presence of ν_s
 - heavy (~ 100 MeV), unstable (~ 10 s)
- Thermal effects
 - Assume interaction of steriles sufficiently strong at $T \sim$ few GeV to maintain thermal equilibrium with e , ν , γ , ...
 - Further, the sterile decouples at $T \sim$ few MeV
 - assume relativistic kinematics throughout
 - proper entropy is conserved: $s \propto a^3 = \text{constant}$ (FLRW)
 - sterile neutrino temperature distribution cooled or “diluted”
$$\frac{T_{\nu_s}(a_{wdc})}{T_\gamma(a_{wdc})} = \left(\frac{g_*(a_{wdc})}{g_*(a_{\nu_s dc})} \right)^{1/3} = \left(\frac{10.75}{61.75} \right)^{1/3} \approx \frac{1}{1.8}$$
 - number density comparable to photons (since lifetime chosen 10's secs)
 - $n(\nu_s) \sim 0.1 n(\gamma)$
- NB: ν_s is out-of-equilibrium with $e\mu\nu\gamma$

Dilution physics (II)

□ Heavy particle decay during/after weak decoupling

□ Interactions

Exothermic

$$\nu_s \rightarrow 3\nu_i$$

$$\nu_s \rightarrow \nu_i + \gamma$$

Endothermic

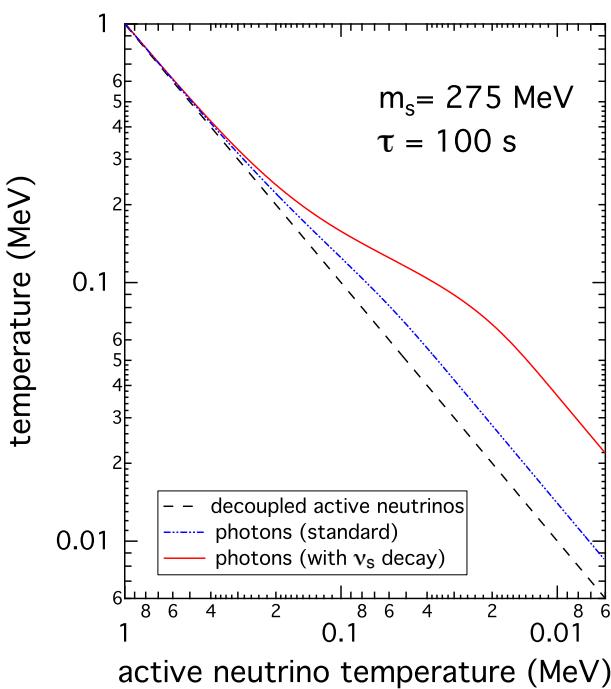
$$\nu_s \rightarrow \nu_i + e^- + e^+$$

$$\nu_s \rightarrow \nu + \mu^+ + \mu^-$$

$$\nu_s \rightarrow \nu + \pi^0$$

$$\nu_s \rightarrow \pi^\pm + e^\mp$$

$$\nu_s \rightarrow \pi^\pm + \mu^\mp$$



□ Entropy production

□ due to out-of-equilibrium decay

□ plasma cools slower than decoupled actives

□ Dilution

□ decoupled actives diluted down

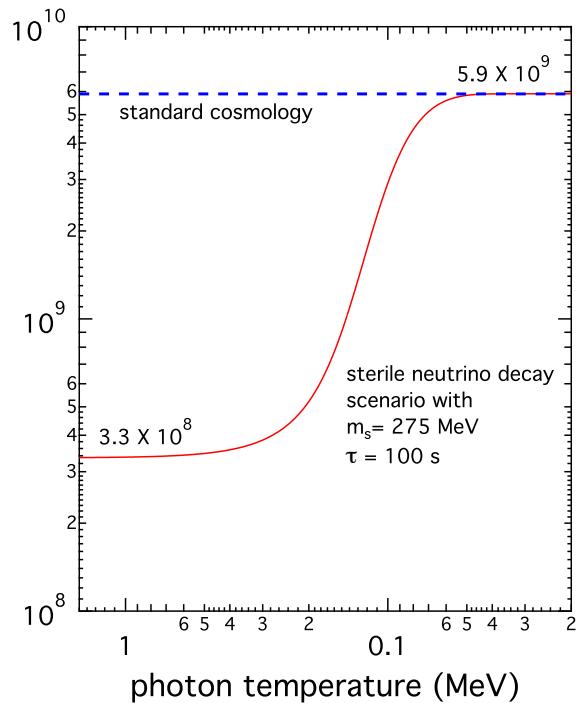
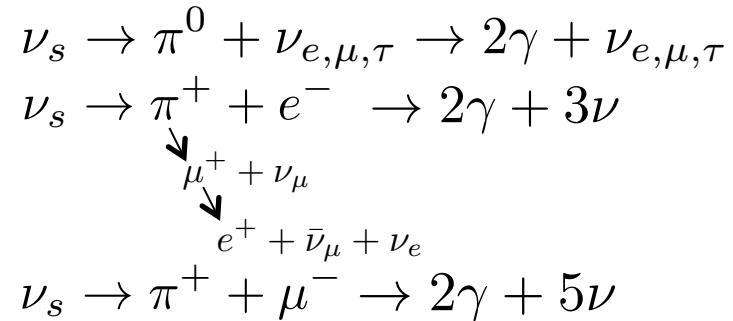
□ Two effects

■ coupling to plasma → reduction in N_{eff}

■ coupling to actives → increase N_{eff}

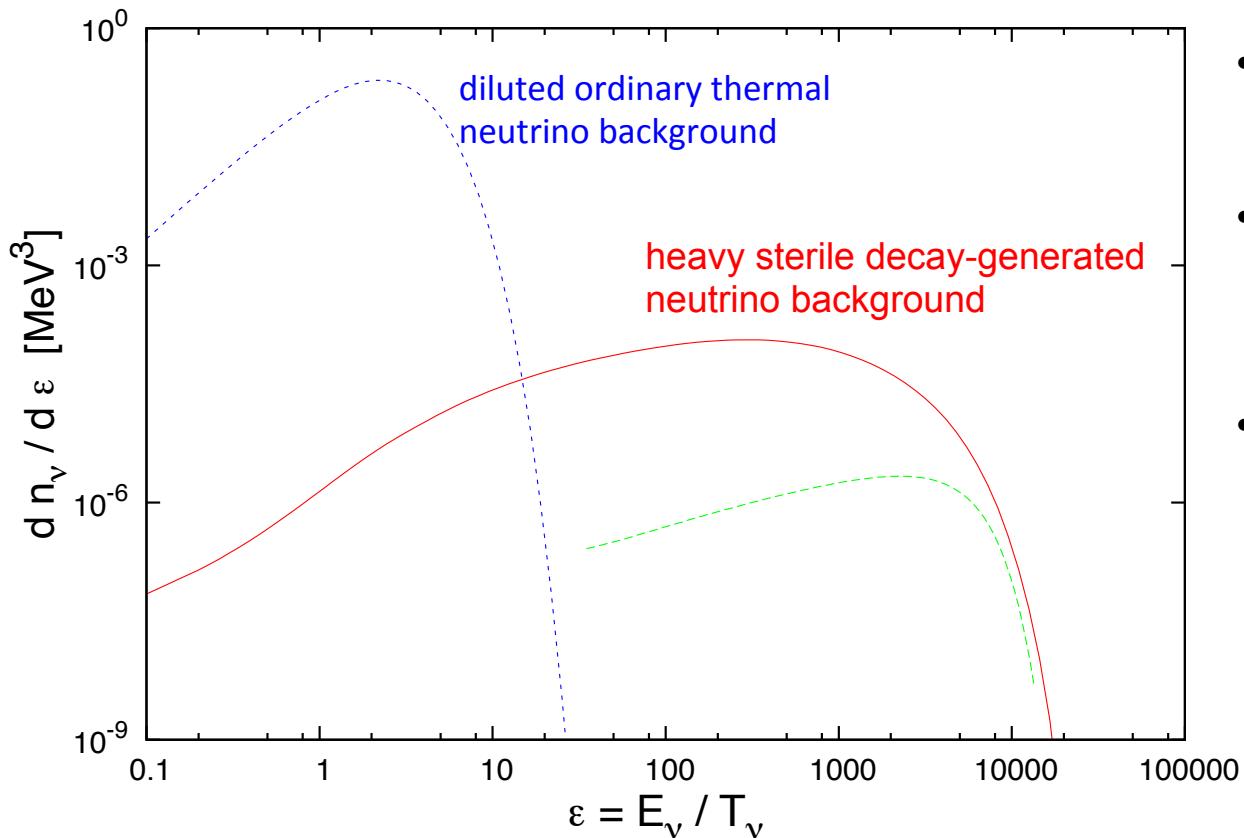
Dilution physics (III)

- Photons thermalize
 - sterile neutrino decay ($m_s <$ few GeV)
- But active neutrinos may not
 - energy/decay-epoch dependent



- Heavy sterile neutrino decay
 - dilution of background ($C\nu B$)
 - generation of radiation energy density: N_{eff}
 - prodigious entropy production

Non-equilibrium distribution of CνB



- Heavy sterile dilutes the normal background neutrino spectrum
- decay-generated spectrum $\sim 10^3$ times more energetic than standard; never non-relativistic
- can't detect neutrino rest mass cosmologically

The Big Question: what effect on BBN? Y_p

Code capabilities & design

- Capabilities
 - Boltzmann equation solver: two classes of Boltzmann equations
 - Nucleosynthesis: Unitary Reaction Network for BBN (previous slides)
 - Neutrino energy transport: **new capability – never before achieved**

$$\begin{aligned}\frac{Df_1}{Dt} = & \int \frac{s}{2E_1} \frac{d^3p_2}{(2\pi)^3(2E_2)} \frac{d^3p_3}{(2\pi)^3(2E_3)} \frac{d^3p_4}{(2\pi)^3(2E_4)} \\ & \times \langle |\mathcal{M}|^2 \rangle (2\pi)^4 \delta^4(P_1 + P_2 - P_3 - P_4) F(p_1, p_2, p_3, p_4) \\ \frac{Df_1}{Dt} = & \frac{\kappa}{32(2\pi)^3} \int_0^\infty p_1 p_2^3 dp_2 \int_{-1}^1 dx \frac{(1-x)^2}{\sqrt{p_1^2 + p_2^2 + 2p_1 p_2 x}} \int_{E_{\min}}^{E_{\max}} dp_3 F(p_1, p_2, p_3, p_1 + p_2 - p_3).\end{aligned}$$

- Various reactions result in seven evaluations of this **triple integral**
- Achieved short turn-around time by **parallelization**

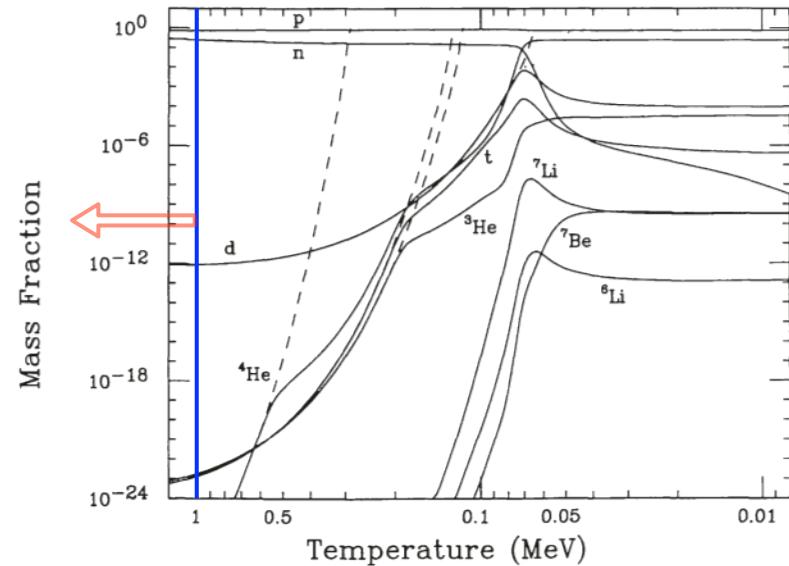
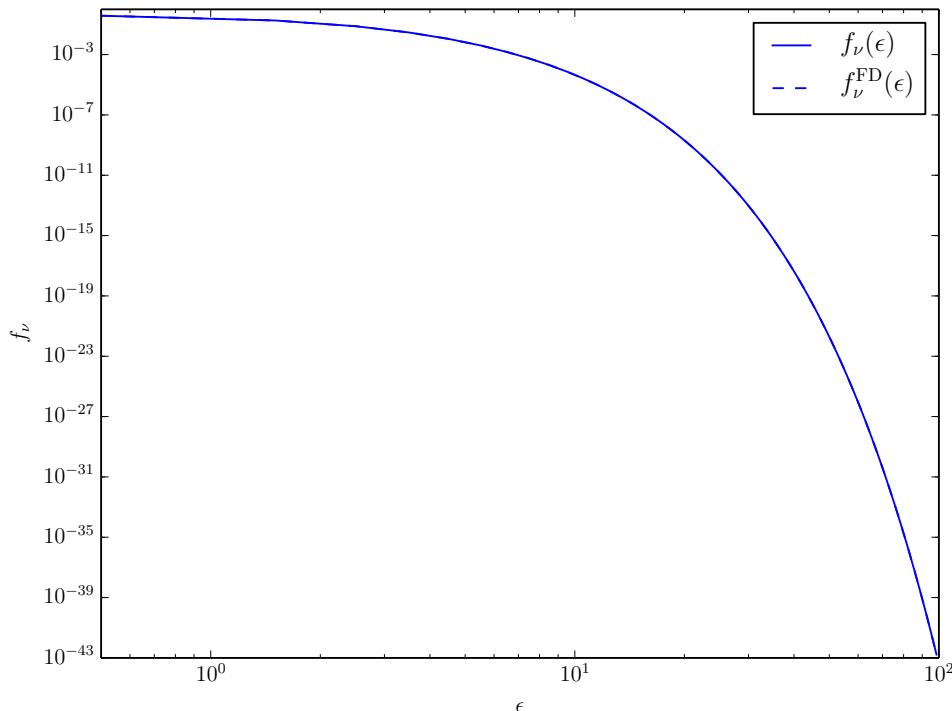
- Design
 - Modular code design for adaptability for **public code release**
 - Allow insertion of “physics packages” to test BSM (not just sterile ν 's)

Code testing/preliminary results

- Evolve assuming equilibrium from 30 MeV → 3 MeV
- Then turn-on only elastic ν -lepton scattering



$$T_\nu = 2.892E+01 \text{ MeV}$$



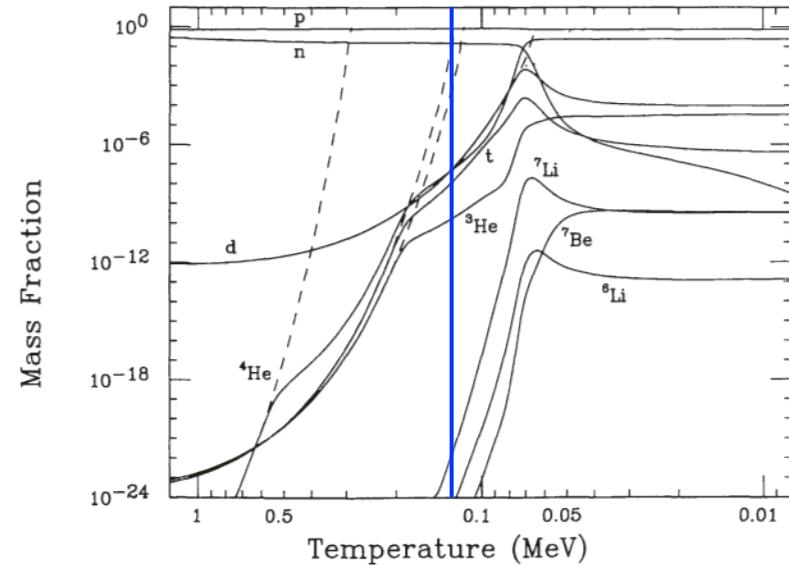
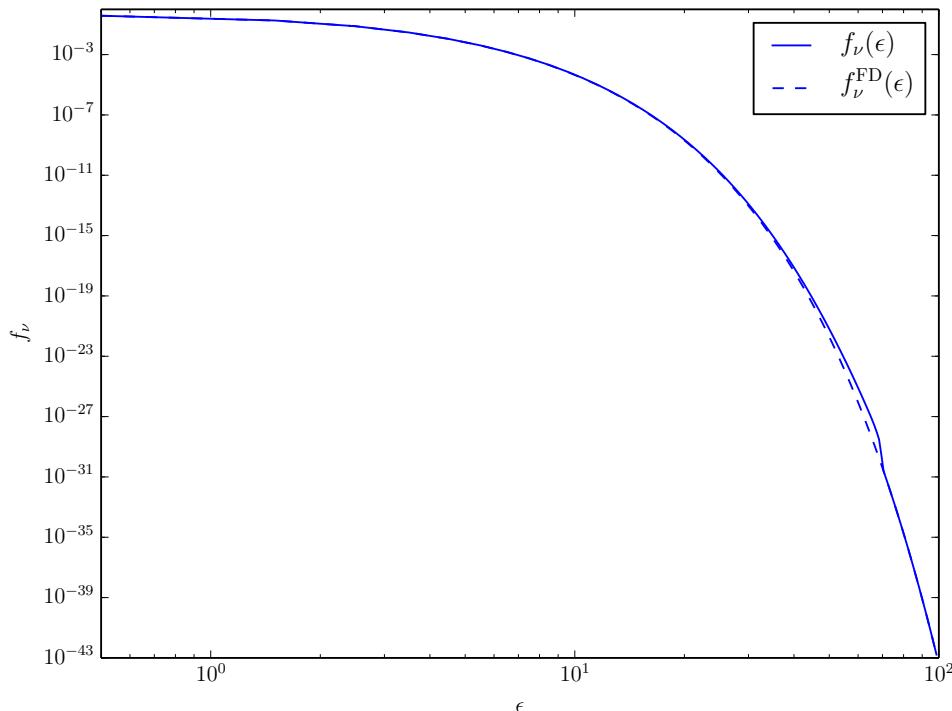
- since the ν & anti- ν are cooler than the e^\pm anticipate upscattering

Code testing/preliminary results

- Evolve assuming equilibrium from 30 MeV → 3 MeV
- Then turn-on only elastic ν -lepton scattering

$$\nu_i + e^\pm \rightarrow \nu_i + e^\pm \quad i = e, \mu, \tau$$

$T_\nu = 1.134\text{E-}01$ MeV



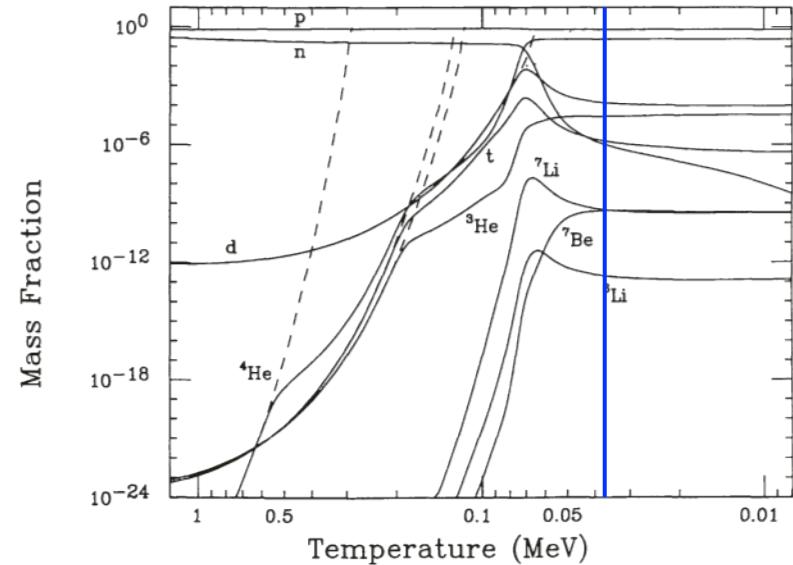
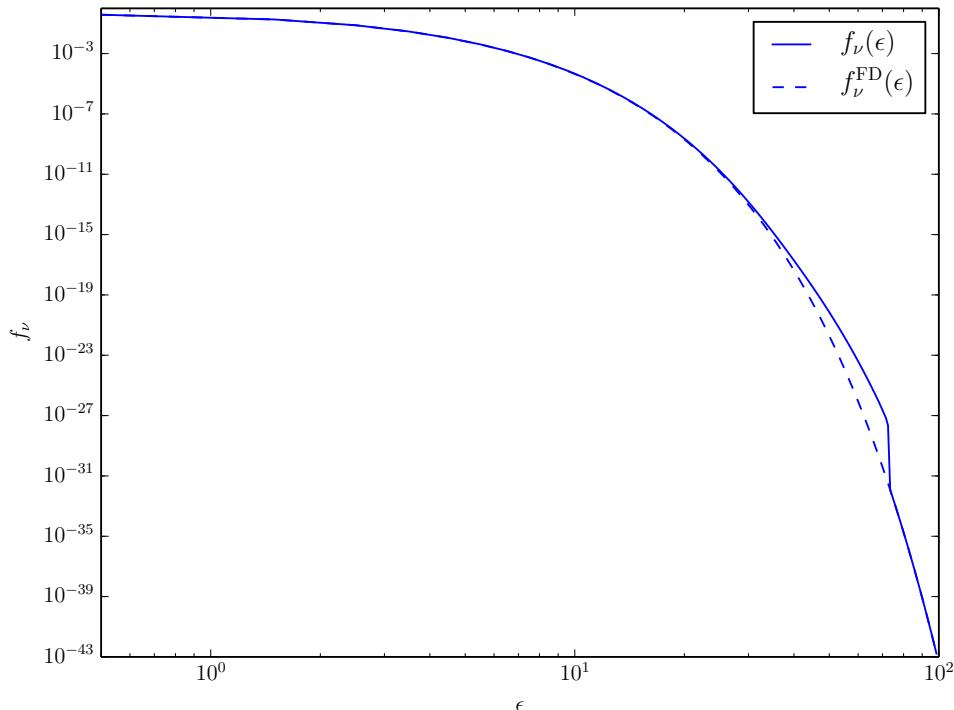
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Code testing/preliminary results

- Evolve assuming equilibrium from 30 MeV → 3 MeV
- Then turn-on only elastic ν -lepton scattering

$$\nu_i + e^\pm \rightarrow \nu_i + e^\pm \quad i = e, \mu, \tau$$

$T_\nu = 3.875\text{E-}02$ MeV



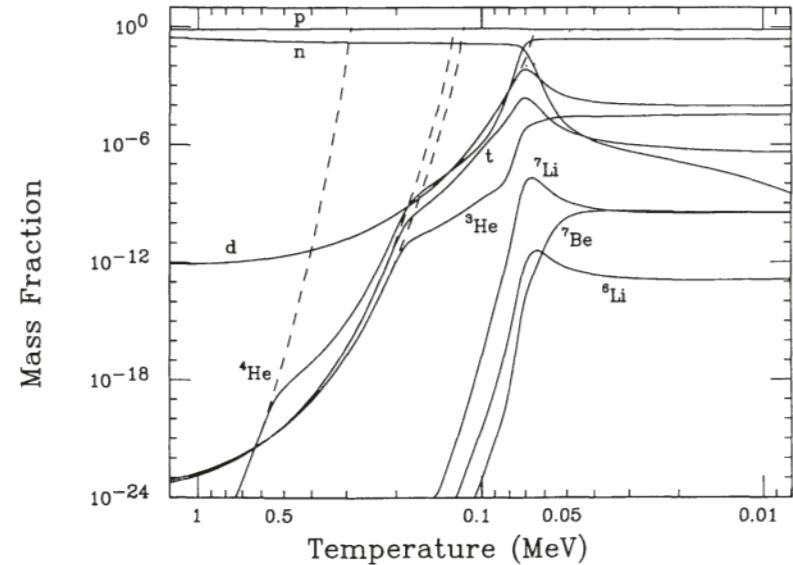
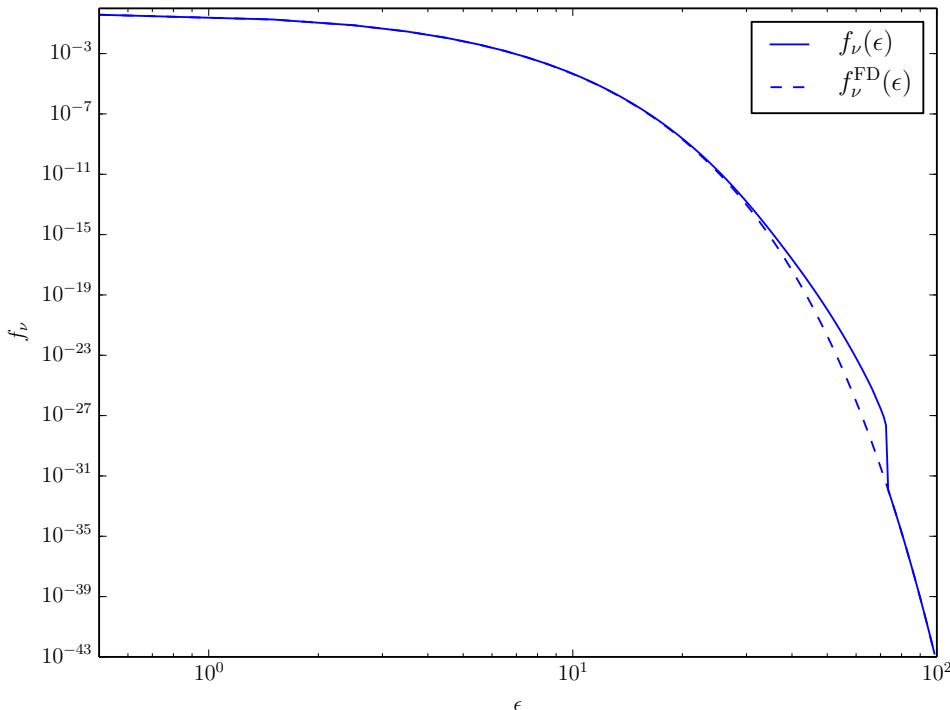
- since the ν & anti- ν are cooler than the e^\pm anticipate upscattering

Code testing/preliminary results

- Evolve assuming equilibrium from 30 MeV → 3 MeV
- Then turn-on only elastic ν -lepton scattering



$$T_\nu = 1.886 \text{E-}03 \text{ MeV}$$



- since the ν & anti- ν are cooler than the e^\pm anticipate upscattering
- **INTERESTING:** because “ ν decoup. complete by e^+e^- annihilation”

Elastic scattering

Initial transport temperature [keV]	N_{eff}
20	3.0055
40	3.0055
100	3.005666
200	3.005936
400	3.006555
1000	3.008414
3000	3.013428

 e^{\pm} annihilation

Initial transport temperature [keV]	N_{eff}
20	3.005584
40	3.005590
100	3.005682
200	3.005985
400	3.006604
1000	3.008309
3000	3.xxxxxx

These preliminary/test results give a nice demonstration that the fundamentals of the neutrino energy transport are working.