Matt Johnson Perimeter Institute/York University

 Most of cosmology is described by General Relativity and Relativistic Hydrodynamics.

- Most of cosmology is described by General Relativity and Relativistic Hydrodynamics.
- The fundamental variables:

 $g_{\mu\nu}(x,t)$   $\rho_i(x,t)$   $u^{\mu}(x,t)$ metric fluid densities fluid velocities

- Most of cosmology is described by General Relativity and Relativistic Hydrodynamics.
- The fundamental variables:

 $g_{\mu\nu}(x,t) \qquad 
ho_i(x,t) \qquad u^{\mu}(x,t)$ metric fluid densities fluid velocities

The laws: Einstein and continuity equations

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} \qquad \nabla_{\mu} T^{\mu\nu} = 0 \qquad T_{\mu\nu} = \sum_{i} \left(\rho_{i} + p_{i}\right) u_{i\mu} u_{i\nu} + p_{i} g_{\mu\nu}$$

- Most of cosmology is described by General Relativity and Relativistic Hydrodynamics.
- The fundamental variables:

 $g_{\mu\nu}(x,t) \qquad 
ho_i(x,t) \qquad u^{\mu}(x,t)$ metric fluid densities fluid velocities

• The laws: Einstein and continuity equations

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} \qquad \nabla_{\mu} T^{\mu\nu} = 0 \qquad T_{\mu\nu} = \sum_{i} \left(\rho_{i} + p_{i}\right) u_{i\mu} u_{i\nu} + p_{i} g_{\mu\nu}$$

Set of coupled PDE's -- need initial conditions!

• For much of the history of the Universe:

$$g_{\mu\nu}(x,t) = \bar{g}_{\mu\nu}(t) + \delta g_{\mu\nu}(x,t)$$
  

$$\rho_i(x,t) = \bar{\rho}_i(t) + \delta \rho_i(x,t) \qquad \delta = \text{ small}$$
  

$$u^{\mu}(x,t) = \bar{u}^{\mu} + \delta u^{\mu}(x,t)$$

• For much of the history of the Universe:

$$g_{\mu\nu}(x,t) = \bar{g}_{\mu\nu}(t) + \delta g_{\mu\nu}(x,t)$$
  

$$\rho_i(x,t) = \bar{\rho}_i(t) + \delta \rho_i(x,t) \qquad \delta = \text{ small}$$
  

$$u^{\mu}(x,t) = \bar{u}^{\mu} + \delta u^{\mu}(x,t)$$

• Why was the universe so nearly homogeneous?

• For much of the history of the Universe:

$$g_{\mu\nu}(x,t) = \bar{g}_{\mu\nu}(t) + \delta g_{\mu\nu}(x,t)$$
  

$$\rho_i(x,t) = \bar{\rho}_i(t) + \delta \rho_i(x,t) \qquad \delta = \text{ small}$$
  

$$u^{\mu}(x,t) = \bar{u}^{\mu} + \delta u^{\mu}(x,t)$$

- Why was the universe so nearly homogeneous?
- For today: this is an extraordinary convenience!

• For much of the history of the Universe:

$$g_{\mu\nu}(x,t) = \bar{g}_{\mu\nu}(t) + \delta g_{\mu\nu}(x,t)$$
  

$$\rho_i(x,t) = \bar{\rho}_i(t) + \delta \rho_i(x,t) \qquad \delta = \text{ small}$$
  

$$u^{\mu}(x,t) = \bar{u}^{\mu} + \delta u^{\mu}(x,t)$$

- Why was the universe so nearly homogeneous?
- For today: this is an extraordinary convenience!



# Initial Conditions

• The homogeneous Universe:

$$p_i = w_i \rho_i$$
  $ar{
ho}_i(t_0)$   
Types of fluids Densities today

# Initial Conditions

• The homogeneous Universe:

$$p_i = w_i \rho_i$$
  
Types of fluids

 $\bar{\rho}_i(t_0)$ 

Densities today

• The Linear Universe:

Characterize statistics of inhomogeneities!

 $\mathcal{P}\left[g_{\mu\nu}\left(x,t=0\right)\right] \quad \mathcal{P}\left[\rho_{i}\left(x,t=0\right)\right] \quad \mathcal{P}\left[u_{i\mu}\left(x,t=0\right)\right]$ 

Assume our Universe is typical.

# 6 Parameter Model of the Universe



Initially small fluctuations collapse to form galaxies, stars, etc.



## That's it! The rest is details.

# Giving Thanks

- The non-linear Universe
  - GR is highly non-linear inferring the state of the early universe would be like asking for the weather 100 million years ago based on the weather today.
  - No general classification of metrics how to characterize initial conditions?
  - Shock waves, singularities, oh my!

# Giving Thanks

- The linear Universe
  - Simple evolution allows initial conditions to be inferred.
  - Background evolution and growth of structure can be analyzed separately.
  - Simple classification of initial conditions and metric degrees of freedom.
  - Physics on different scales evolves independently (Fourier modes independent).

#### The rest

#### Now for some details....



13.7 Billion Years: the present.9.1 Billion Years: our sun ignites.

100 million years galaxies and first stars form.

380,000 years: neutral atoms form.

1 second: atomic nuclei form.

 $10^{-6}$  seconds: protons and neutrons form. ?Big Bang?

• The metric in a flat, homogeneous, isotropic universe:

$$ds^{2} = -dt^{2} + a^{2}(t)\delta_{ij}dx^{i}dx^{j}$$

$$\begin{tabular}{l} \label{eq:star} & \end{tabular} \\ \end{tabular} \end{tabular} \end{tabular} \end{tabular}$$

• The metric in a flat, homogeneous, isotropic universe:



constant comoving distance = growing physical distance

• Conformal time:

$$\eta = \int \frac{dt}{a(t)} \qquad \qquad ds^2 = a^2(\eta) \left[ -d\eta^2 + \delta_{ij} dx^i dx^j \right]$$

Conformal time:

$$\eta = \int \frac{dt}{a(t)} \qquad \qquad ds^2 = a^2(\eta) \left[ -d\eta^2 + \delta_{ij} dx^i dx^j \right]$$



Conformal time:

$$\eta = \int \frac{dt}{a(t)} \qquad ds^2 = a^2(\eta) \left[ -d\eta^2 + \delta_{ij} dx^i dx^j \right]$$

#### particle horizon



• Equations of motion in a homogeneous universe:

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} \quad \square \searrow \quad H^2 \equiv \left(\frac{\dot{a}}{a}\right) = \frac{8\pi G\rho}{3}$$

• Equations of motion in a homogeneous universe:

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} \qquad \longrightarrow \qquad H^2 \equiv \left(\frac{\dot{a}}{a}\right) = \frac{8\pi G \rho}{3}$$
$$\nabla_{\mu} T^{\mu\nu} = 0 \qquad \implies \qquad \dot{\rho} = -3H\left(\rho + p\right)$$
$$p = w\rho$$

• Equations of motion in a homogeneous universe:

$$G_{\mu\nu} = 8\pi G T_{\mu\nu} \qquad \longrightarrow \qquad H^2 \equiv \left(\frac{\dot{a}}{a}\right) = \frac{8\pi G \rho}{3}$$

$$\nabla_{\mu} T^{\mu\nu} = 0 \qquad \implies \qquad \dot{\rho} = -3H \left(\rho + p\right)$$

$$p = w\rho$$

• Solutions:

$$\rho = \rho_0 a^{-3(1+w)} \qquad a(t) = a_0 t^{\frac{2}{3(1+w)}}$$

different fluids gravitate differently!



• Evolution of the scale factor:



• Energy budget:

$$\left(\frac{H}{H_0}\right)^2 = \sum_i \Omega_i a^{-3(1+w_i)} \qquad \sum_i \Omega_i = 1$$

• Energy budget:

$$\left(\frac{H}{H_0}\right)^2 = \sum_i \Omega_i a^{-3(1+w_i)} \qquad \sum_i \Omega_i = 1$$

$$\int_i^{\text{Dark Matter}} \frac{26.8\%}{4.9\%} \qquad \left(\Omega_r \sim 10^{-4}\right)^2$$

$$\int_i^{\text{Dark Energy}} \frac{68.3\%}{68.3\%}$$





 $a = \frac{1}{1+z}$ • Redshift:  $z = \frac{\lambda_{obs} - \lambda_{em}}{\lambda_{obs}}$  $a = 1, \quad z = 0$  $z_{\rm eq} = 3400,$  $z_* = 1090,$  $z_{\rm re} \sim 11,$  $z_{\rm gal} \sim 11 - 12,$  $z_{\rm surveys} \lesssim 1$ ,  $z_{\Lambda} = .28,$  $z_{\text{Virgo}} = .003$ 

• There is structure in the Universe:



• In Fourier space:  

$$\rho(x,t) = \frac{1}{V} \sum \rho(k,t) e^{i\vec{k}\cdot\vec{x}}$$

$$\rho(k,t)$$

$$\bar{\rho}(t)$$

$$k$$

• There is structure in the Universe:



• There is structure in the Universe:


• There is structure in the Universe:



• There is structure in the Universe:



Structure: 
$$rac{\delta 
ho}{ar{
ho}}\gtrsim 1$$

 $\mathcal{O}(10^6)$  galaxies  $\mathcal{O}(10^3)$  clusters  $\mathcal{O}(1)$  superclusters

• There is structure in the Universe:



$$\begin{array}{lll} \mbox{Structure:} & \frac{\delta\rho}{\bar{\rho}}\gtrsim 1 & \mathcal{O}(10^6) \ \mbox{galaxies} \\ & \mathcal{O}(10^3) \ \ \mbox{clusters} \\ & \mathcal{O}(1) \ \ \ \mbox{superclusters} \end{array}$$

- There is structure on all scales which have had a chance to undergo gravitational collapse.
- The largest structures in the Universe define a scale above which the fluctuations in density are linear:

$$k = .1 - .01 {\rm Mpc}^{-1} \qquad \lambda \sim 1\% ~{\rm of}~{\rm observable}~{\rm Universe}$$
 wave number

• There is structure in the Universe:



Luminous: Baryons Photons Semi-Luminous: Neutrinos

• There is structure in the Universe:



Luminous: Baryons Photons Semi-Luminous: Neutrinos Dark: Dark Matter Dark Energy

• For non-relativistic matter in flat space:

$$\begin{aligned} \frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) &= 0 & \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} + \frac{\nabla p}{\rho} + \nabla \phi = 0 \\ \text{continuity} & \text{Euler} \\ \nabla^2 \phi &= 4\pi G\rho \end{aligned}$$

Poisson

• For non-relativistic matter in flat space:

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0 \qquad \frac{\partial \vec{v}}{\partial t} + (\vec{v} \cdot \nabla) \vec{v} + \frac{\nabla p}{\rho} + \nabla \phi = 0$$
continuity
Euler
$$\nabla^2 \phi = 4\pi G\rho$$
Poisson
Linearize:

$$\begin{split} \rho &= \bar{\rho} + \delta \rho \qquad \phi = \bar{\phi} + \delta \phi \qquad \vec{v} = \bar{\vec{v}} + \delta \vec{v} \\ p &= \bar{p} + c_s^2 \delta \rho \\ c_s^2 &= \frac{\partial p}{\partial \rho} = w \end{split}$$

• Linearized equation of motion:

$$\frac{\partial^2 \delta \rho}{\partial t^2} - c_s^2 \nabla^2 \delta \rho - 4\pi G \bar{\rho} \delta \rho = 0$$

• Linearized equation of motion:

$$\frac{\partial^2 \delta \rho}{\partial t^2} - c_s^2 \nabla^2 \delta \rho - 4\pi G \bar{\rho} \delta \rho = 0$$

• Fourier transform:

$$\delta\rho(t,\vec{x}) = \int \frac{d^3k}{(2\pi)^3} \delta\rho(t,k) e^{i\vec{k}\cdot\vec{x}}$$

• Linearized equation of motion:

$$\frac{\partial^2 \delta \rho}{\partial t^2} - c_s^2 \nabla^2 \delta \rho - 4\pi G \bar{\rho} \delta \rho = 0$$

• Fourier transform:

$$\delta\rho(t,\vec{x}) = \int \frac{d^3k}{(2\pi)^3} \delta\rho(t,k) e^{i\vec{k}\cdot\vec{x}} \qquad \frac{\partial^2\delta\rho}{\partial t^2} + \left(c_s^2k^2 - 4\pi G\bar{\rho}\right)\delta\rho = 0$$

• Linearized equation of motion:

$$\frac{\partial^2 \delta \rho}{\partial t^2} - c_s^2 \nabla^2 \delta \rho - 4\pi G \bar{\rho} \delta \rho = 0$$

• Fourier transform:

$$\delta\rho(t,\vec{x}) = \int \frac{d^3k}{(2\pi)^3} \delta\rho(t,k) e^{i\vec{k}\cdot\vec{x}} \qquad \frac{\partial^2\delta\rho}{\partial t^2} + \left(c_s^2k^2 - 4\pi G\bar{\rho}\right)\delta\rho = 0$$

$$\delta \rho(t,k) = A \exp(i\omega(k)t) + B \exp(-i\omega(k)t)$$

$$\omega(k) = \sqrt{k^2 c_s^2 - 4\pi G\bar{\rho}}$$

$$\omega(k) = \sqrt{k^2 c_s^2 - 4\pi G\bar{\rho}}$$

$$\omega(k) = \sqrt{k^2 c_s^2 - 4\pi G\bar{\rho}}$$

• Jeans scale: competition between pressure and gravity

$$\lambda_J = \frac{2\pi}{k_J} = c_s \left(\frac{\pi}{G\bar{\rho}}\right)^{1/2}$$

$$\omega(k) = \sqrt{k^2 c_s^2 - 4\pi G\bar{\rho}}$$

Jeans scale: competition between pressure and gravity



• In an expanding universe waves are stretched:



$$\int_{\Delta x} \int_{\Delta x} \int_{\Delta s} \int_{$$

• In an expanding universe waves are stretched:



- Only gravitationally bound (non-linear) structures separate from the Hubble flow.
- Expansion inhibits collapse.

 $t_{\rm coll} \sim (4\pi G \bar{\rho})^{-1/2} \sim H^{-1}$  Expansion will be relevant!

 $t_{\rm coll} \sim (4\pi G \bar{\rho})^{-1/2} \sim H^{-1}$  Expansion will be relevant!

• Including expansion (on small scales):

$$\left[\frac{d^2}{dt^2} + 2H\frac{d}{dt} + \left(c_s^2\frac{k_{\rm com}^2}{a^2} - 4\pi G\bar{\rho}\right)\right]\frac{\delta\rho(t,k_{\rm com})}{\bar{\rho}(t)} = 0$$

 $t_{\rm coll} \sim (4\pi G \bar{\rho})^{-1/2} \sim H^{-1}$  Expansion will be relevant!

Including expansion (on small scales):

$$\left[\frac{d^2}{dt^2} + 2H\frac{d}{dt} + \left(c_s^2\frac{k_{\rm com}^2}{a^2} - 4\pi G\bar{\rho}\right)\right]\frac{\delta\rho(t,k_{\rm com})}{\bar{\rho}(t)} = 0$$



 $t_{\rm coll} \sim (4\pi G \bar{\rho})^{-1/2} \sim H^{-1}$  Expansion will be relevant!

Including expansion (on small scales):

$$\left[\frac{d^2}{dt^2} + 2H\frac{d}{dt} + \left(c_s^2\frac{k_{\rm com}^2}{a^2} - 4\pi G\bar{\rho}\right)\right]\frac{\delta\rho(t,k_{\rm com})}{\bar{\rho}(t)} = 0$$

 $c_s^2 = 0 \quad \begin{array}{c|c} \mbox{radiation} & \mbox{matter} & \mbox{dark energy} \\ \hline \frac{\delta\rho}{\bar{\rho}} \propto \log(a) & \mbox{d}{\bar{\rho}} \propto a & \mbox{d}{\bar{\rho}} \propto \cosh t. \end{array}$ 

To have structure, need:

$$\frac{\delta\rho}{\bar{\rho}}(t_{\rm eq})\gtrsim 3\times 10^{-4}$$

1

- An important scale: comoving horizon
- Horizon crossing: k = aH



- An important scale: comoving horizon
- Horizon crossing: k = aH



conformal time  $\,\eta\,$ 

1

- An important scale: comoving horizon
- Horizon crossing: k = aH



conformal time  $\,\eta\,$ 

1

1

- An important scale: comoving horizon
- Horizon crossing: k = aH



- An important scale: comoving horizon
- Horizon crossing: k = aH



conformal time  $\,\eta\,$ 

1

- An important scale: comoving horizon
- Horizon crossing: k = aH



conformal time  $\eta$ 

1

- An important scale: comoving horizon
- Horizon crossing: k = aH



conformal time  $\,\eta\,$ 

1

- An important scale: comoving horizon
- Horizon crossing: k = aH



conformal time  $\,\eta\,$ 

1

#### **Relativistic Perturbations**

• The full model:

$$\delta g_{\mu\nu}(x,t) \quad \delta \rho_i(x,t) \quad \delta u^{\mu}(x,t)$$
$$\nabla_{\mu} \delta T^{\mu\nu} = 0 \qquad \delta G_{\mu\nu} = 8\pi G \delta T_{\mu\nu}$$

#### **Relativistic Perturbations**

• The full model:

$$\delta g_{\mu\nu}(x,t) \quad \delta \rho_i(x,t) \quad \delta u^{\mu}(x,t)$$

$$\nabla_{\mu} \delta T^{\mu\nu} = 0 \qquad \delta G_{\mu\nu} = 8\pi G \delta T_{\mu\nu}$$
• Metric fluctuations



#### **Relativistic Perturbations**

• The full model:

$$\delta g_{\mu\nu}(x,t) \quad \delta \rho_i(x,t) \quad \delta u^{\mu}(x,t)$$

$$\nabla_{\mu} \delta T^{\mu\nu} = 0 \qquad \delta G_{\mu\nu} = 8\pi G \delta T_{\mu\nu}$$
fluctuations
$$- 4 \text{ scalar} \longleftrightarrow \delta \rho_i(x,t)$$

• Metric fluctuations  

$$\delta g_{\mu\nu}(x,t) \xrightarrow{\qquad 4 \text{ scalar } \leftrightarrow \delta \rho_i(x,t)} 4 \text{ vector } \text{ decay, not major player} 2 \text{ tensor } \text{ gravity waves}$$

• Gauge choice: only 6 DOF.

$$ds^{2} = a(\eta)^{2} \left[ -(1+2\Psi) \, d\eta^{2} + (1+2\Phi) \, \delta_{ij} dx^{i} dx^{j} \right]$$

"Conformal Newtonian Gauge"



```
photons tightly
coupled to baryons
imperfect fluid with
   c_s^2 \simeq 1/2
dark matter
perturbations begin
to grow
```

equality



equality recombination				
	photons tightly coupled to baryons imperfect fluid with $c_s^2 \simeq 1/2$	neutral atoms form photons and baryons decouple: CMB is released!		
	dark matter perturbations begin to grow	baryons begin to collapse into dark matter halos		

equality recomb		ination reioni	zation
	photons tightly coupled to baryons imperfect fluid with $c_s^2 \simeq 1/2$	neutral atoms form photons and baryons decouple: CMB is released!	
	dark matter perturbations begin to grow	baryons begin to collapse into dark matter halos	
#### Important events

equ	ality recomb	ination reioni	zation
	photons tightly coupled to baryons imperfect fluid with $c_s^2\simeq 1/2$	neutral atoms form photons and baryons decouple: CMB is released!	first non-linear structure stars reionize the Universe
	dark matter perturbations begin to grow	baryons begin to collapse into dark matter halos	hierarchical structure formation some CMB photons re- scatter

## Baryon Acoustic Oscillations

• Before recombination, photons and baryons are coupled:

$$\left[\frac{d^2}{dt^2} + 2H\frac{d}{dt} + \left(c_s^2\frac{k_{\rm com}^2}{a^2} - 4\pi G\bar{\rho}\right)\right]\frac{\delta\rho(t,k_{\rm com})}{\bar{\rho}(t)} = 0$$
  
effective pressure from coupling

## Baryon Acoustic Oscillations

• Before recombination, photons and baryons are coupled:

$$\left[\frac{d^2}{dt^2} + 2H\frac{d}{dt} + \left(c_s^2\frac{k_{\rm com}^2}{a^2} - 4\pi G\bar{\rho}\right)\right]\frac{\delta\rho(t,k_{\rm com})}{\bar{\rho}(t)} = 0$$
  
effective pressure from coupling

• Sound horizon:

$$r_s = \int_0^\eta d\eta \ c_s(\eta)$$

$$\frac{\delta\rho}{\bar{\rho}} \propto \cos\left[kr_s(\eta)\right]$$

frequency of oscillation for sound waves

Standard ruler!

## Baryon Acoustic Oscillations

• Before recombination, photons and baryons are coupled:



• The Universe is filled with a gas of photons.

$$dN(t, \vec{x}) = f(t, \vec{x}, p) \frac{d^3 x d^3 p}{(2\pi)^3}$$

$$f = \frac{1}{\exp\left[\frac{p}{T(t)}\right] - 1}$$
  
Bose-Einstein

distribution

• The Universe is filled with a gas of photons.

$$dN(t, \vec{x}) = f(t, \vec{x}, p) \frac{d^3 x d^3 p}{(2\pi)^3}$$

$$f = \frac{1}{\exp\left[\frac{p}{T(t)}\right] - 1}$$
  
Bose-Einstein  
distribution

T(t) -- all photons stretched equally

- The Universe's most perfect blackbody.
- The temperature today is 2.73 K.

• Perturbations are characterized by:

$$f = \frac{1}{\exp\left[\frac{p}{T(t)(1+\Theta(\vec{x},t,\hat{p}))}\right] - 1} \qquad \Theta = \frac{\delta T}{T}$$

• Perturbations are characterized by:

Observers at each position see an

anisotropic distribution of photons

$$f = \frac{1}{\exp\left[\frac{p}{T(t)(1+\Theta(\vec{x},t,\hat{p}))}\right] - 1}$$

 $\Theta = \frac{\delta T}{T}$ 

$$dN(t_{\rm now}, \vec{x}_{\rm here}) = \int f(t_{\rm now}, \vec{x}, p) \delta(\vec{x}_{\rm here}) \frac{d^3 x d^3 p}{(2\pi)^3}$$

Full set of coupled  $\left(\frac{\delta\rho}{-}\right)$ ,  $v_{\rm dm}$ ,  $\left(\frac{\delta\rho}{-}\right)$ ,  $v_{\rm b}$ ,  $\Theta$ ,  $\Phi$ ,  $\Psi$ the tempera anisotrop

ature 
$$(\bar{\rho})_{dm}^{ature}, (\bar{\rho})_{b}^{ature}, (\bar{\rho})_{b}^{ature}$$

Full set of coupled variables go into finding  $\left( \frac{\delta \rho}{\bar{\rho}} \right)$  the temperature anisotropy

$$\left(\frac{\rho}{\bar{\rho}}\right)_{\rm dm}, v_{\rm dm}, \left(\frac{\delta\rho}{\bar{\rho}}\right)_{\rm b}, v_{\rm b}, \Theta, \Phi, \Psi$$

• Sachs-Wolfe -- valid on largest scales

$$\left(\frac{\Delta T}{T}\right)_{\text{fin}} = \left(\frac{\Delta T}{T}\right)_{\text{init}} - \Phi_{\text{init}} = -\frac{\Phi_{\text{init}}}{3}$$

$$\begin{array}{c} \text{Intrinsic} & \text{gravitational} \\ \text{temperature} & \text{redshift} \\ \text{variations} \end{array}$$

Full set of coupled variables go into finding the temperature anisotropy  $\left(\frac{\delta \mu}{\bar{\rho}}\right)$ 

$$\left(\frac{\overline{\rho}}{\overline{\rho}}\right)_{\mathrm{dm}}, v_{\mathrm{dm}}, \left(\frac{\delta\rho}{\overline{\rho}}\right)_{\mathrm{b}}, v_{\mathrm{b}}, \Theta, \Phi, \Psi$$

• Sachs-Wolfe -- valid on largest scales

$$\begin{pmatrix} \Delta T \\ T \end{pmatrix}_{\text{fin}} = \begin{pmatrix} \Delta T \\ T \end{pmatrix}_{\text{init}} - \Phi_{\text{init}} = -\frac{\Phi_{\text{init}}}{3}$$
Intrinsic gravitational temperature redshift variations

Integrated Sachs-Wolfe -- time dependence of potentials.

• In linear theory, can sum up the contribution from each fourier mode separately:



• In linear theory, can sum up the contribution from each fourier mode separately:



• In linear theory, can sum up the contribution from each fourier mode separately:



• Convenient to perform spherical harmonic transform:

$$\Theta(t, \vec{x}, \hat{p}) = \sum_{\ell} \sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m}(\hat{p})$$

• Convenient to perform spherical harmonic transform:

$$\Theta(t, \vec{x}, \hat{p}) = \sum_{\ell} \sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m}(\hat{p})$$

• The transfer function:  $\Phi_{\text{init}}(k) \rightarrow a_{\ell m}$ 

$$a_{\ell m} = \int \frac{d^3 k}{(2\pi)^3} \Delta_{\ell}(k) \Phi_{\text{init}}(k) Y_{\ell m}(\hat{k})$$

projection, evolution, ISW, etc.

• Convenient to perform spherical harmonic transform:

$$\Theta(t, \vec{x}, \hat{p}) = \sum_{\ell} \sum_{m=-\ell}^{\ell} a_{\ell m} Y_{\ell m}(\hat{p})$$

• The transfer function:  $\Phi_{\text{init}}(k) \rightarrow a_{\ell m}$ 

$$a_{\ell m} = \int \frac{d^3 k}{(2\pi)^3} \Delta_{\ell}(k) \Phi_{\text{init}}(k) Y_{\ell m}(\hat{k})$$

projection, evolution, ISW, etc.

• Computed numerically: CAMB, CMBFast, etc.

## The Power Spectrum

• Fluctuations are characterized statistically:





$$\langle \Phi(k)\Phi(k')\rangle = \delta^3(k-k')P(k)$$

$$\langle a_{\ell m} a_{\ell' m'} \rangle = \delta_{\ell \ell'} \delta_{m m'} C_{\ell}$$

gaussian:  $\langle \Phi(k) \Phi(k') \Phi(k'') \rangle = 0$ 

$$C_{\ell} = \frac{2}{\pi} \int dk \ k^2 \Delta_{\ell}^2(k) P(k)$$

Relates statistics of primordial fluctuations to the statistics of fluctuations in the CMB

#### The Power Spectrum

$$P(k) = Ak^{n_s - 1} \longrightarrow C_\ell$$



#### Comments

- In a statistically homogeneous and isotropic universe with gaussian fluctuations, the power spectrum is all the information there is.
- There is power on scales of order the size of the observable universe -- superhorizon fluctuations.
- The structure of the acoustic peaks is determined by the contents of the universe as well as the initial conditions.

# Other cool things in CMB

• Lensing of the CMB: information on intervening structure



- Polarization of the CMB: primordial gravitational waves
- Sunyaev-Zeldovich effect: shadows of galaxy clusters in the CMB
- Combined mass and number of neutrinos.

#### 6 Parameter Model of the Universe

#### $\Lambda CDM$

	Planck		Planck+lensing		Planck+WP	
Parameter	Best fit	68% limits	Best fit	68% limits	Best fit	68% limits
$\overline{\Omega_{\mathrm{b}}h^2}$	0.022068	$0.02207 \pm 0.00033$	0.022242	$0.02217 \pm 0.00033$	0.022032	$0.02205 \pm 0.00028$
$\Omega_{\rm c} h^2$	0.12029	$0.1196 \pm 0.0031$	0.11805	$0.1186 \pm 0.0031$	0.12038	$0.1199 \pm 0.0027$
100θ <sub>MC</sub>	1.04122	$1.04132 \pm 0.00068$	1.04150	$1.04141 \pm 0.00067$	1.04119	$1.04131 \pm 0.00063$
au	0.0925	$0.097 \pm 0.038$	0.0949	$0.089 \pm 0.032$	0.0925	$0.089^{+0.012}_{-0.014}$
$n_{\rm s}$	0.9624	$0.9616 \pm 0.0094$	0.9675	$0.9635 \pm 0.0094$	0.9619	$0.9603 \pm 0.0073$
$\ln(10^{10}A_{\rm s})$	3.098	$3.103 \pm 0.072$	3.098	$3.085 \pm 0.057$	3.0980	$3.089^{+0.024}_{-0.027}$

## 6 Parameter Model of the Universe

#### $\Lambda CDM$

	Planck		Planck+lensing		Planck+WP		
Parameter	Best fit	68% limits	Best fit	68% limits	Best fit	68% limits	
$\Omega_{ m b}h^2$	0.022068	$0.02207 \pm 0.00033$	0.022242	$0.02217 \pm 0.00033$	0.022032	0.02205 = 0.00028	
$\Omega_{\rm c} h^2$	0.12029	0.1196 ± 0.0031	0.11805	$0.1186 \pm 0.0031$	0.12038	0.1199 = 0.0027	
$100\theta_{\rm MC}$	1.04122	$1.04132 \pm 0.00068$	1.04150	$1.04141 \pm 0.00067$	1.04119	1.04131 = 0.00063	
τ	0.0925	$0.097 \pm 0.038$	0.0949	$0.089 \pm 0.032$	0.0925	$0.089^{+0.012}_{-0.014}$	
$n_{\rm s}$	0.9624	0.9616 ± 0.0094	0.9675	$0.9635 \pm 0.0094$	0.9619	0.9603 = 0.0073	
$\ln(10^{10}A_{\rm s})$	3.098	$3.103 \pm 0.072$	3.098	$3.085 \pm 0.057$	3.0980	$3.089^{+0.024}_{-0.027}$	